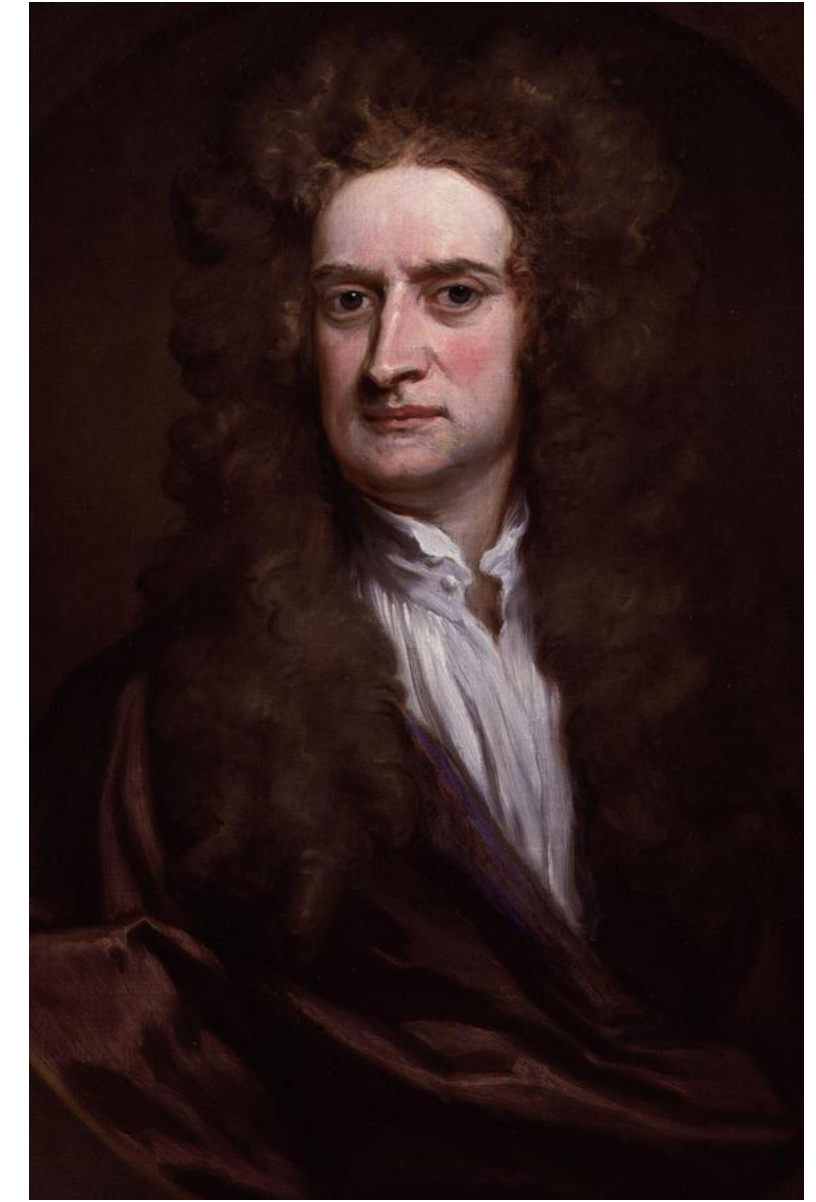
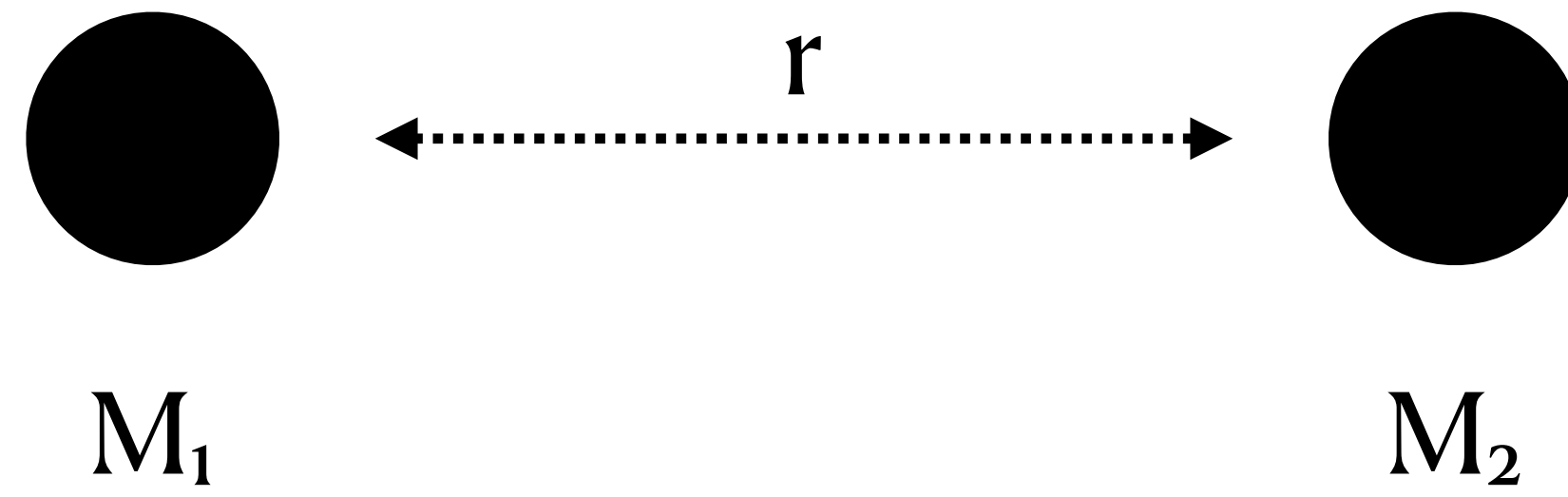


Testing Gravity and Quantum Mechanics

Surjeet Rajendran

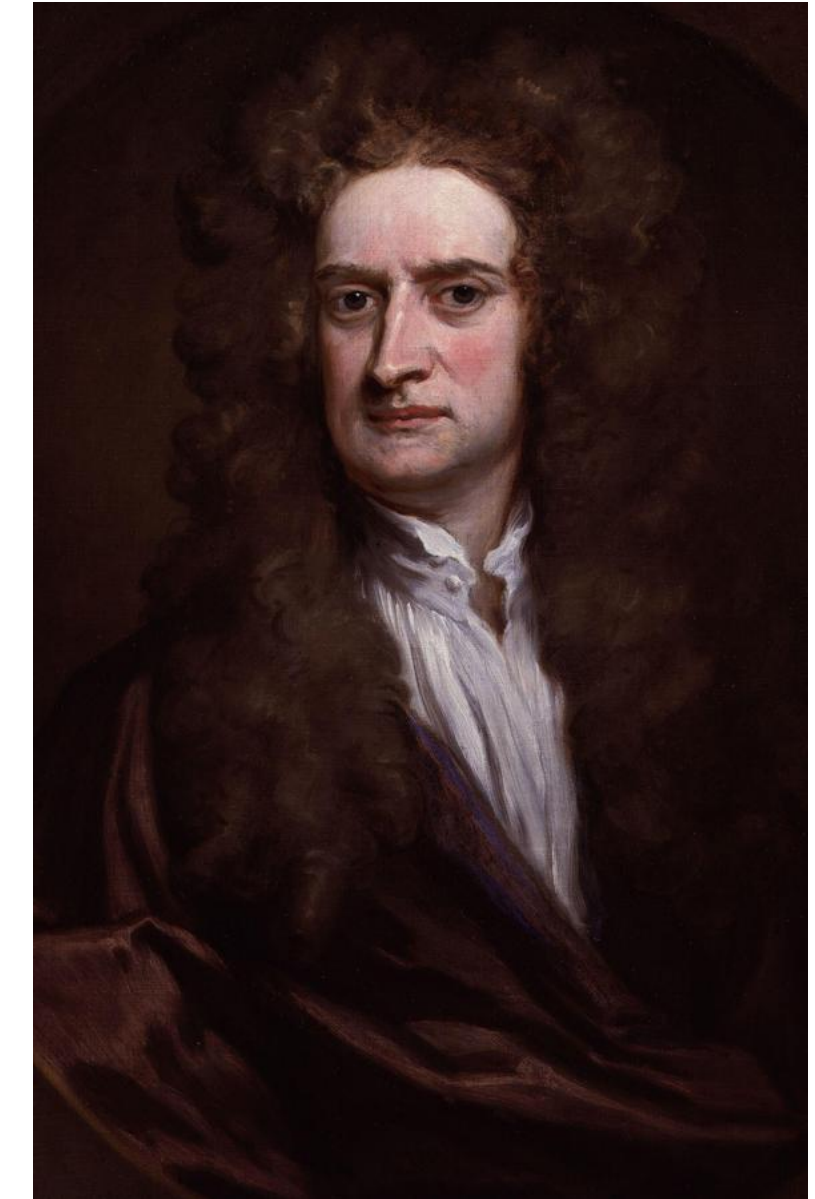
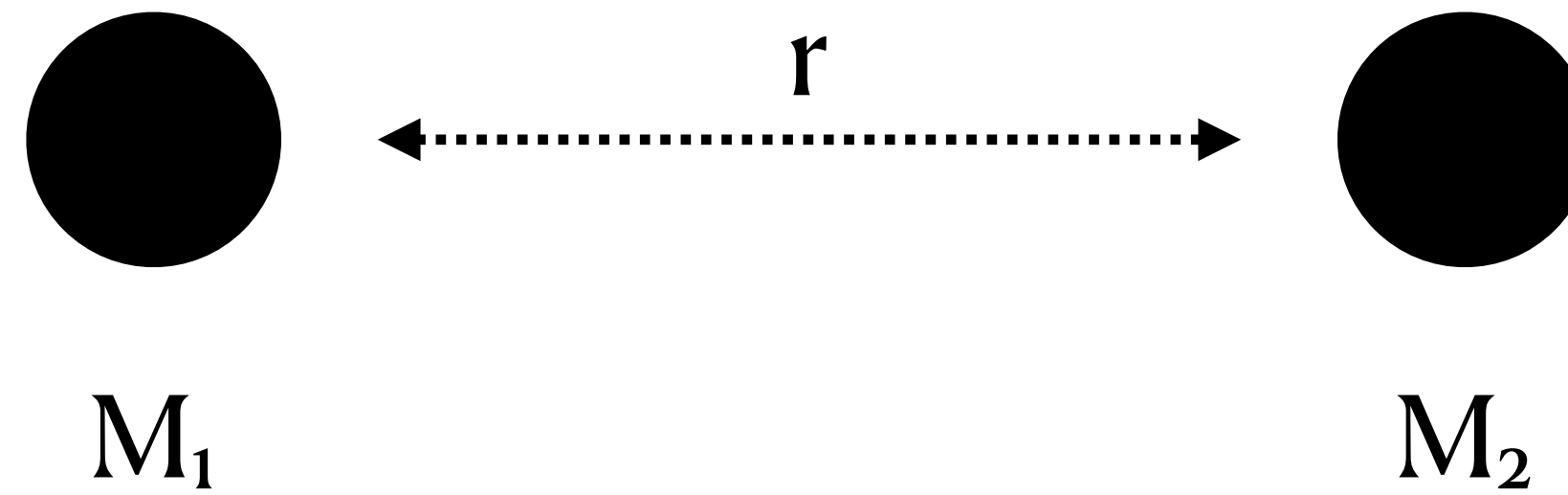
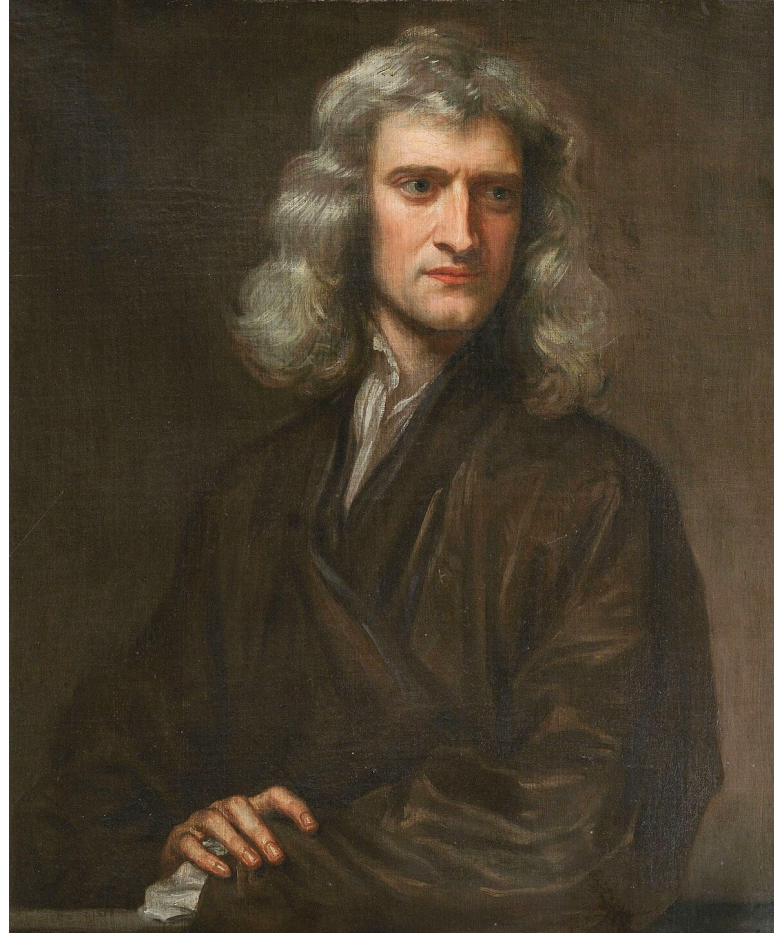
Gravity

Testing Gravity



$$F = \frac{GM_1M_2}{r^2}$$

Testing Gravity

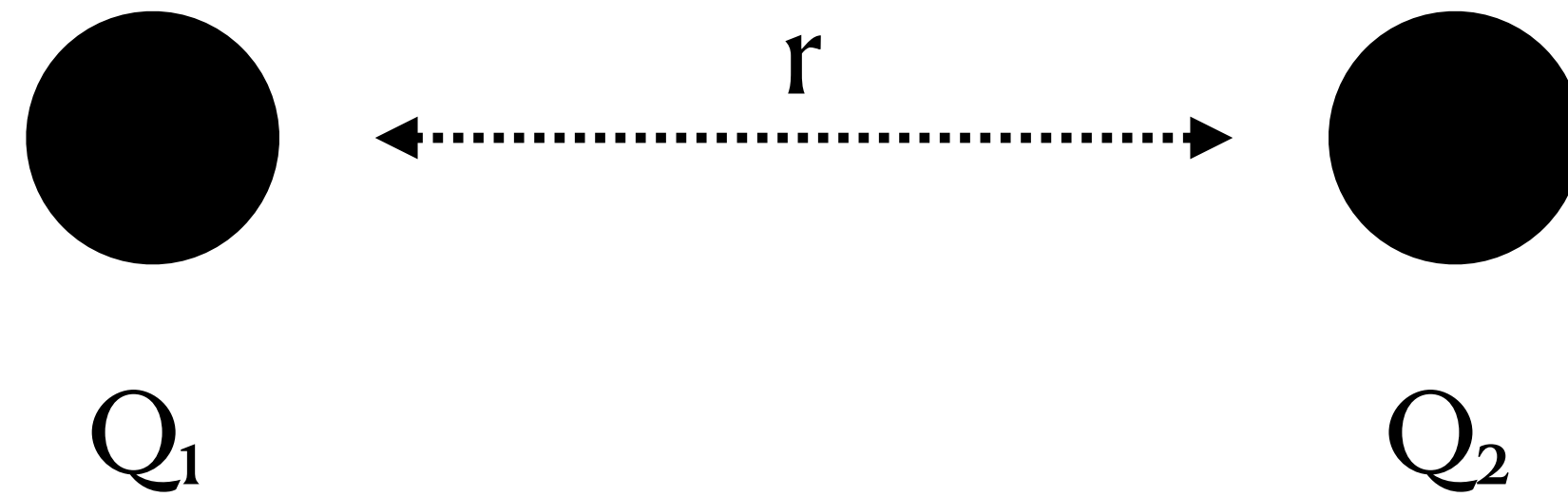


$$F = \frac{GM_1M_2}{r^2}$$

What happens at small r ?

What happens at large M ?

Electromagnetism

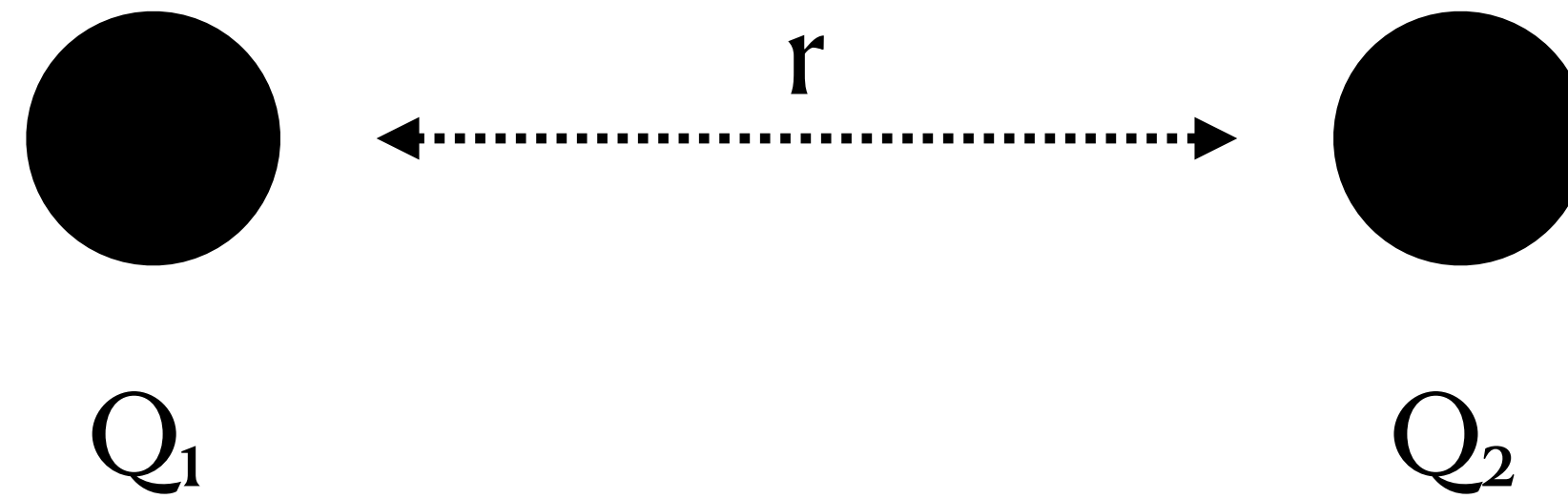


$$F_E = \frac{Q_1 Q_2}{r^2}$$

What happens at small r ? Large Q ?



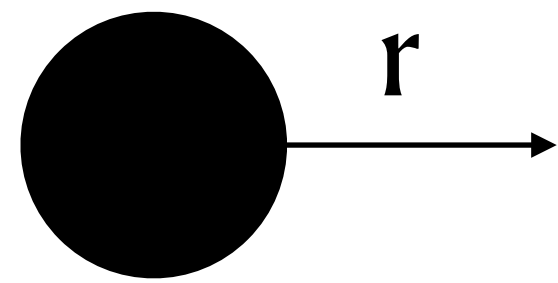
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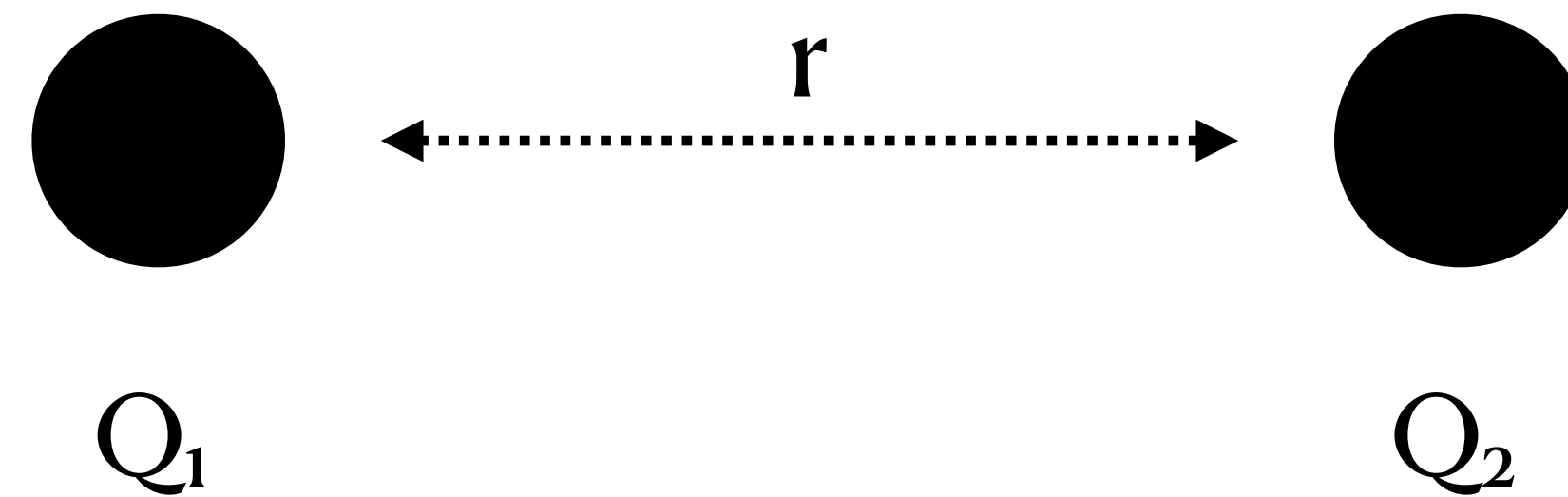
Small r



$$U = \int_{r_i}^{r_f} d^3r \left(\frac{Q_1}{r^2} \right)^2 = \int_{r_i}^{r_f} dr \frac{Q_1^2}{r^2} = Q_1^2 \left(\frac{1}{r_i} - \frac{1}{r_f} \right)$$



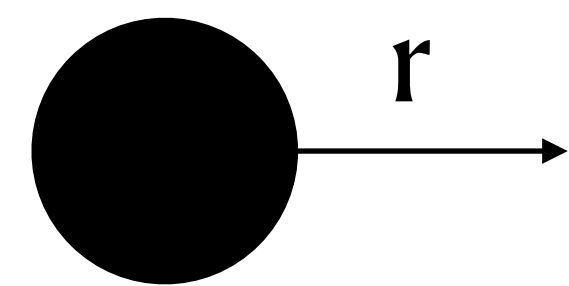
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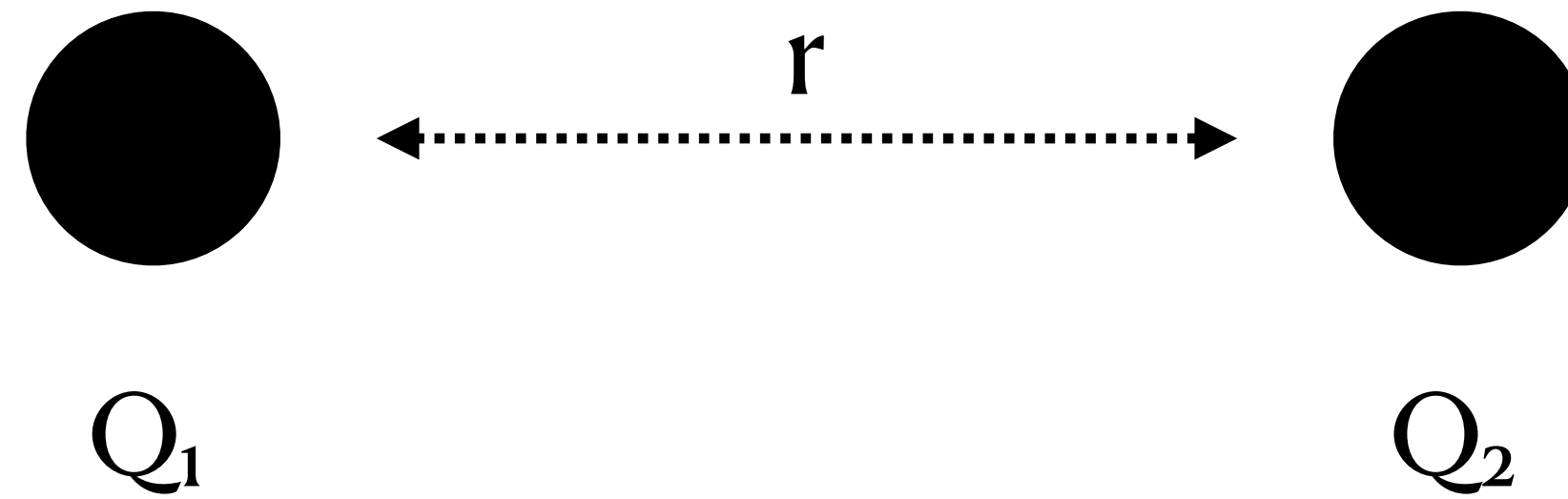
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$$U \rightarrow \infty \text{ as } r_i \rightarrow 0$$

Finite Electron mass - theory wrong at short distance!

Classical theory replaced by quantum mechanics

Electromagnetism



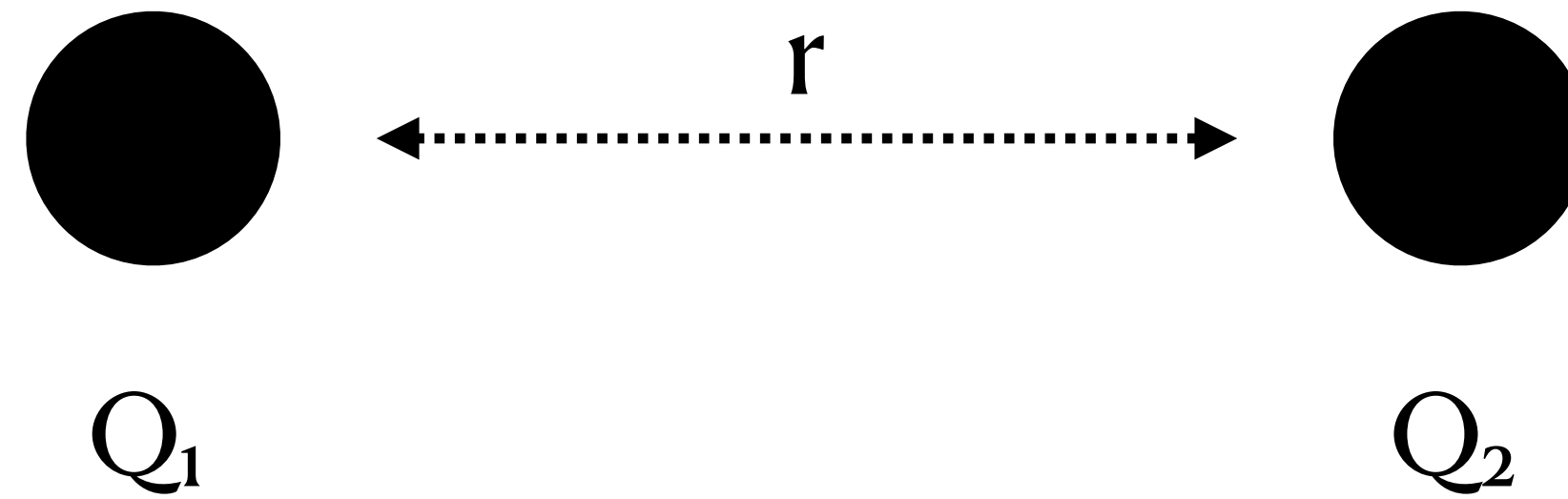
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Large Q

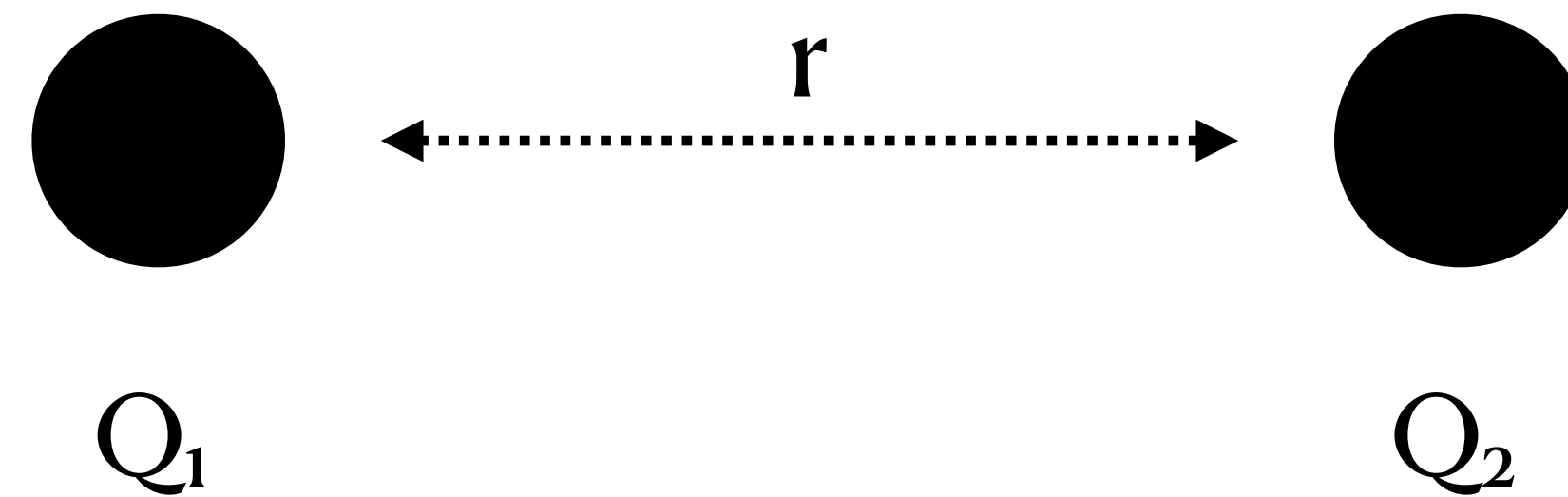
$$U = \int_{r_i}^{r_f} d^3r \left(\frac{Q_1}{r^2} \right)^2$$

$$U > 2m_e$$





Electromagnetism



$$F_E = \frac{Q_1 Q_2}{r^2}$$

What happens at small r ? Large Q ?

Large Q

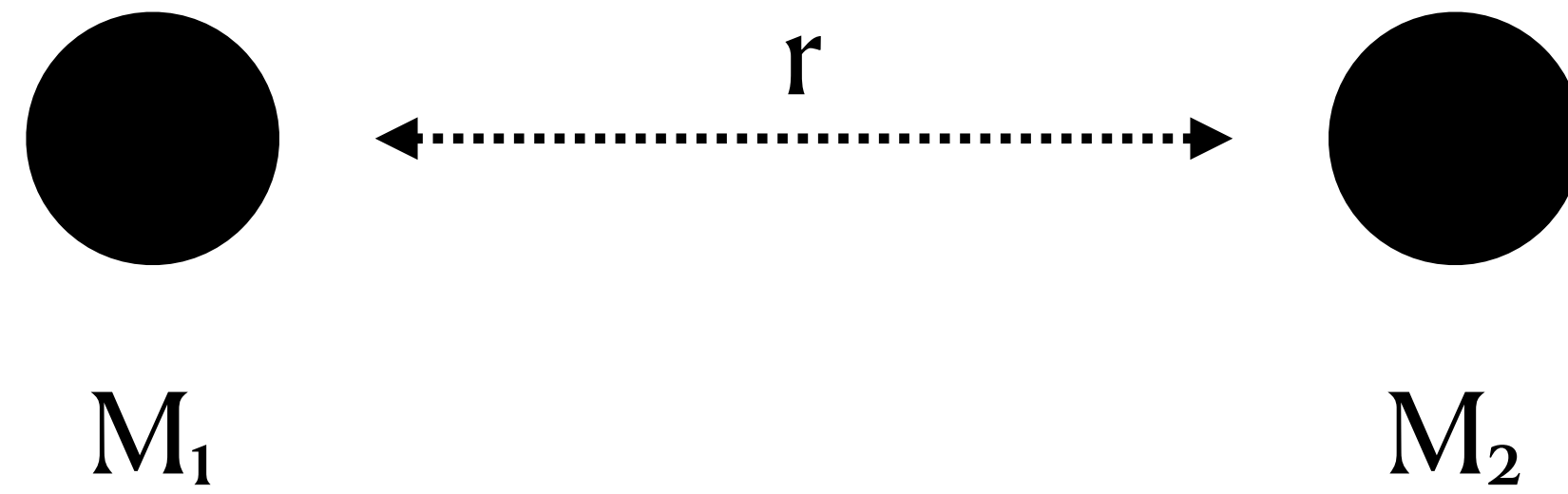
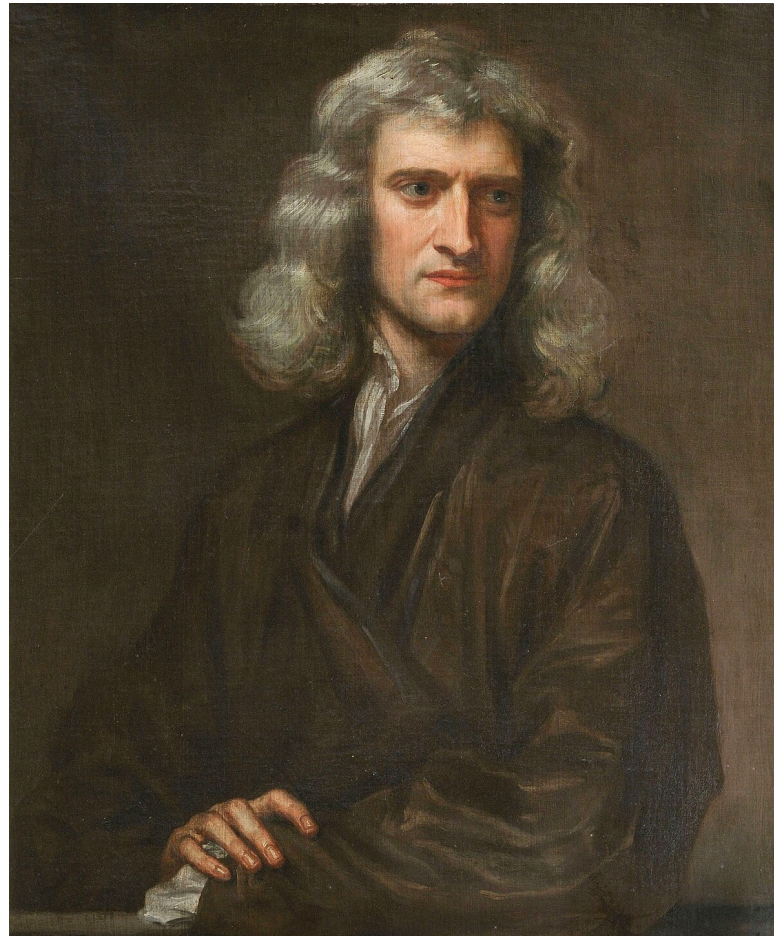
$$U = \int_{r_i}^{r_f} d^3r \left(\frac{Q_1}{r^2} \right)^2$$

$$U > 2m_e$$

Pair produce electrons + positrons - neutralize field

Classical theory replaced by quantum mechanics

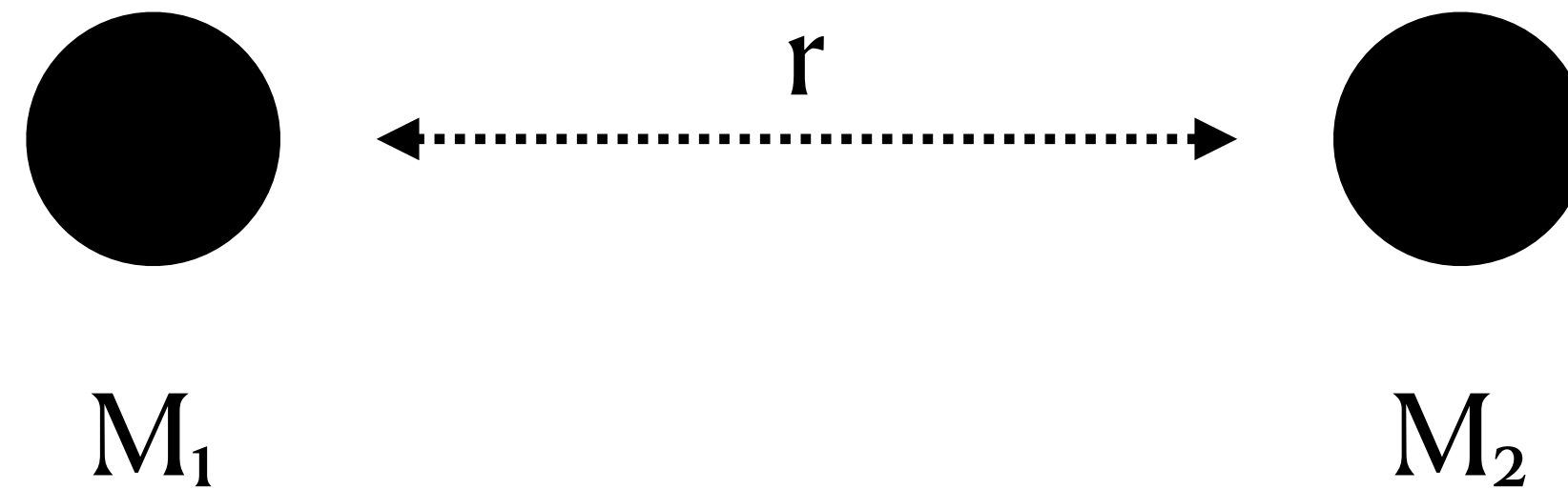
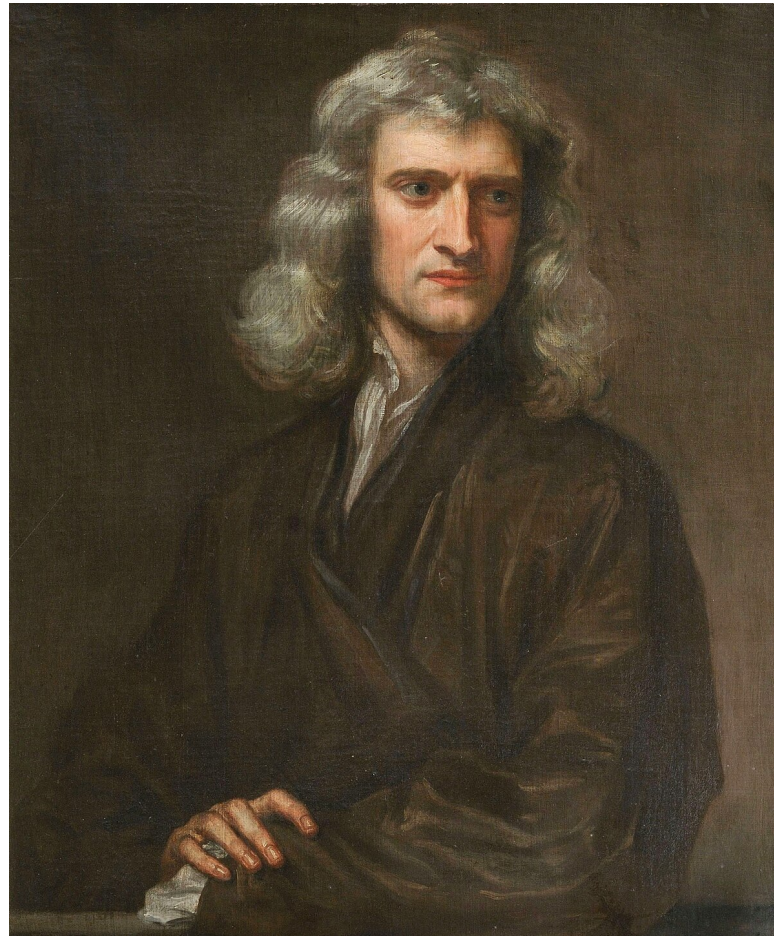
Testing Gravity



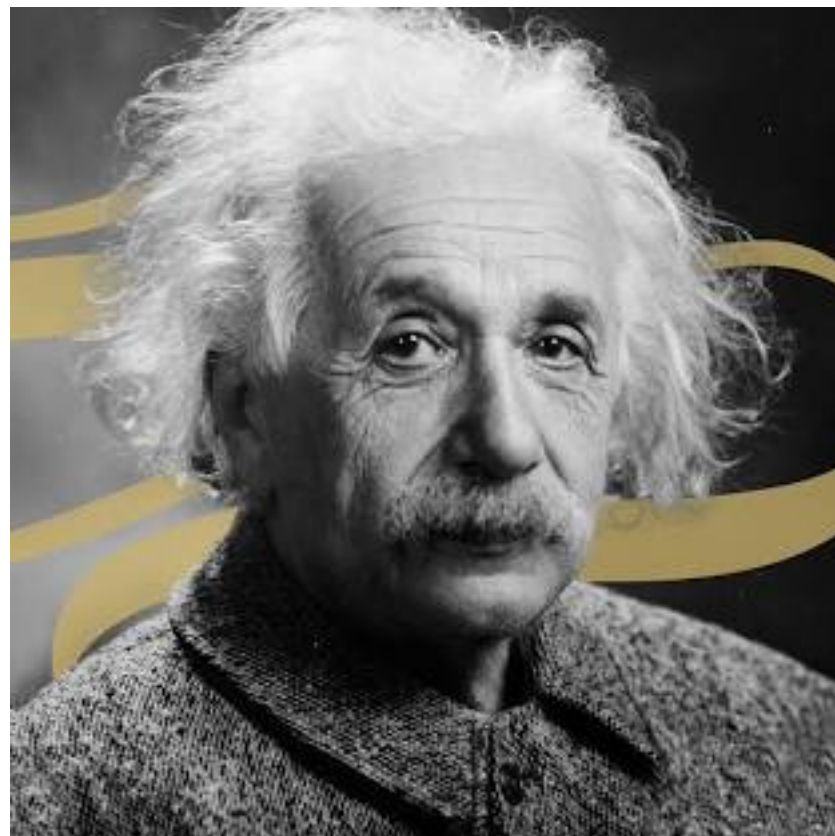
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Large M

Testing Gravity



$$F = \frac{GM_1M_2}{r^2}$$

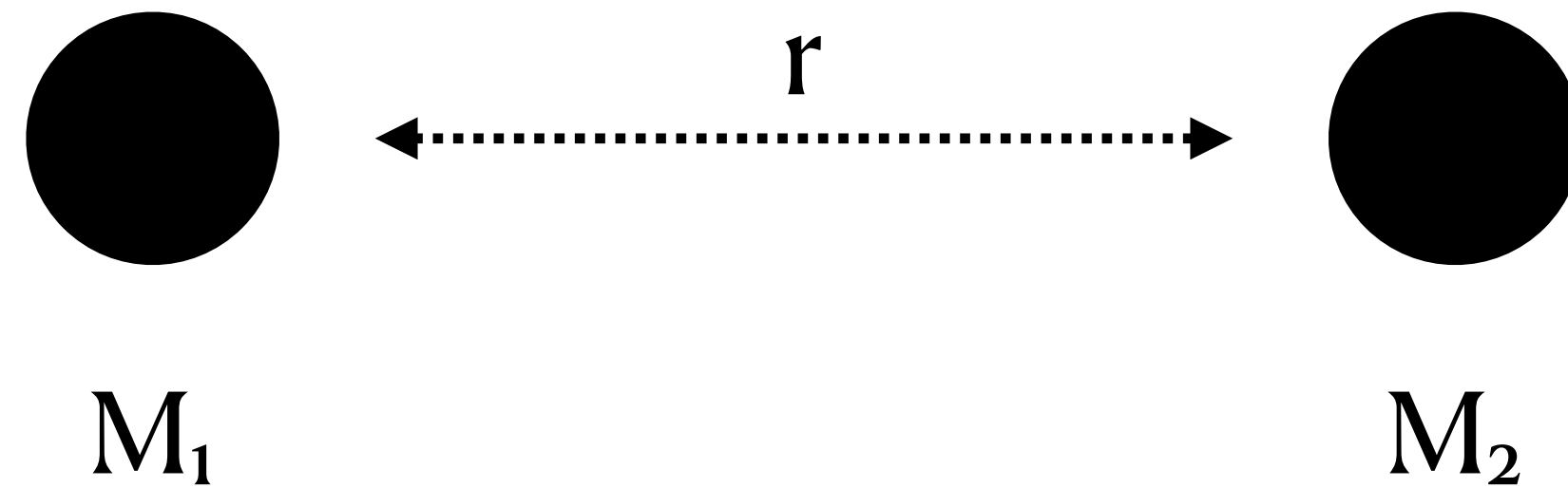
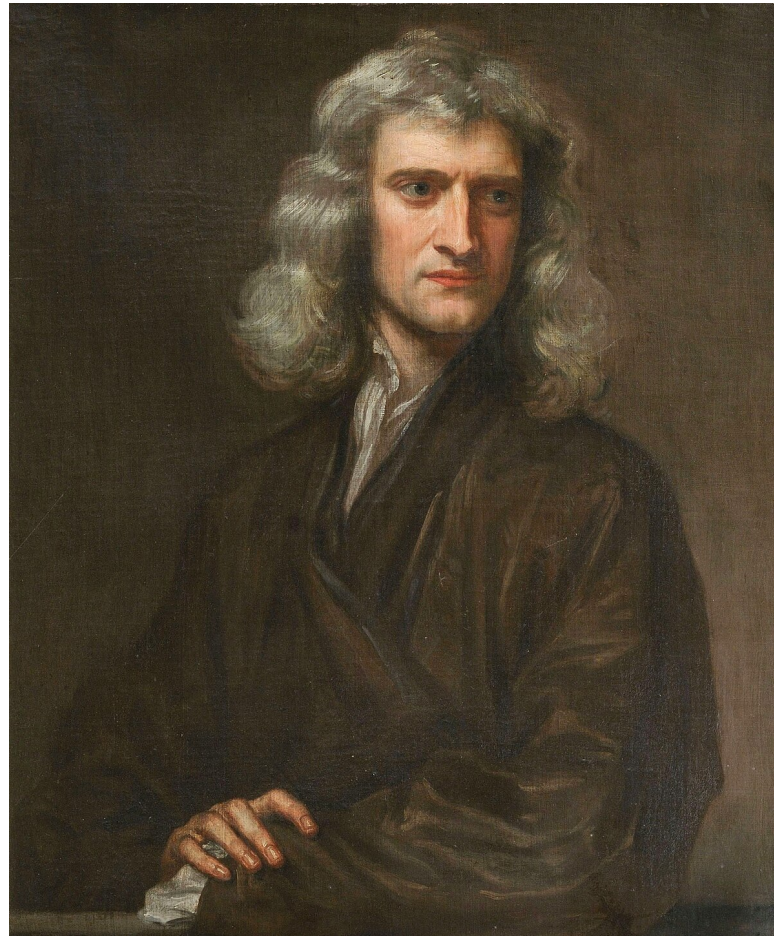


Large M

Gravity produces more gravity -
interacts with itself. Non-Linear

$$G_{\mu\nu} = T_{\mu\nu}$$

Testing Gravity



$$F = \frac{GM_1M_2}{r^2}$$

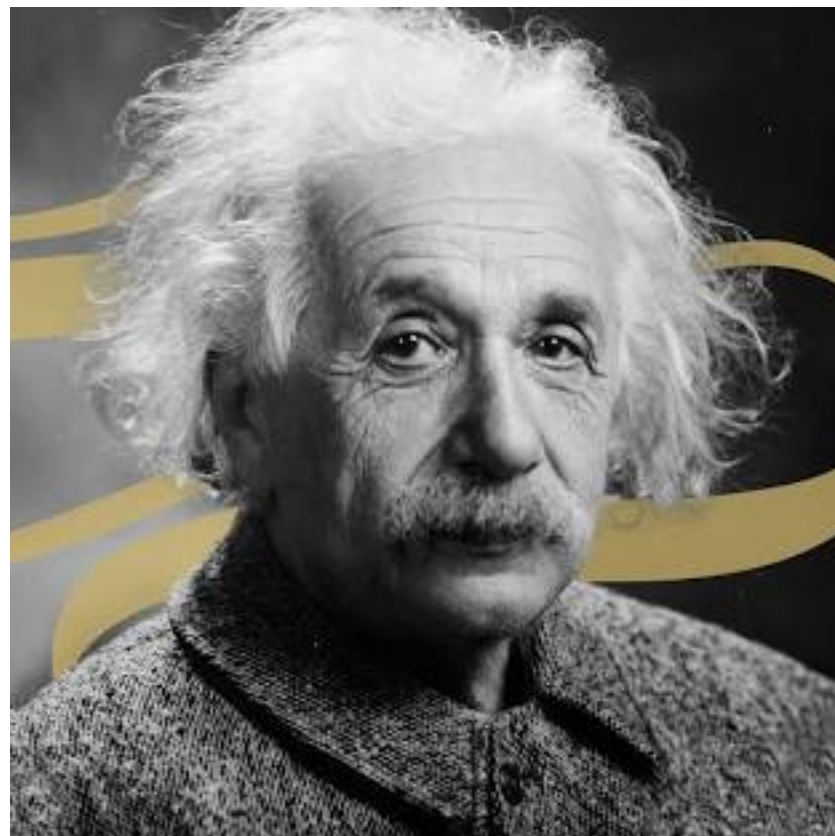
Large M

Gravity produces more gravity -
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Even bigger M

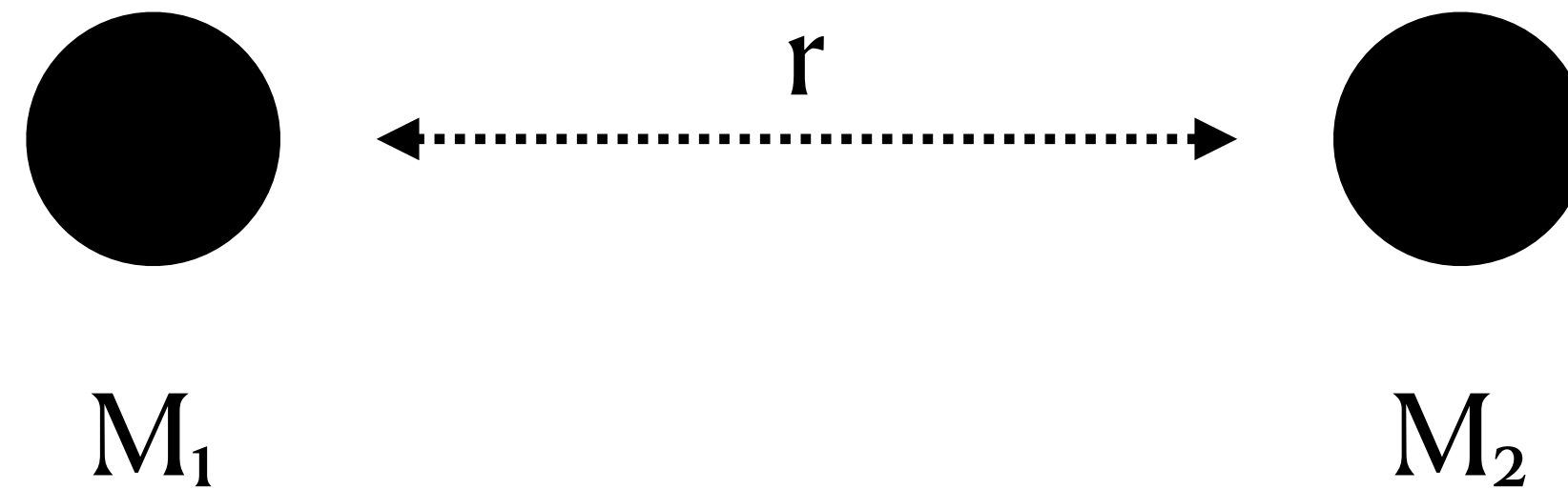
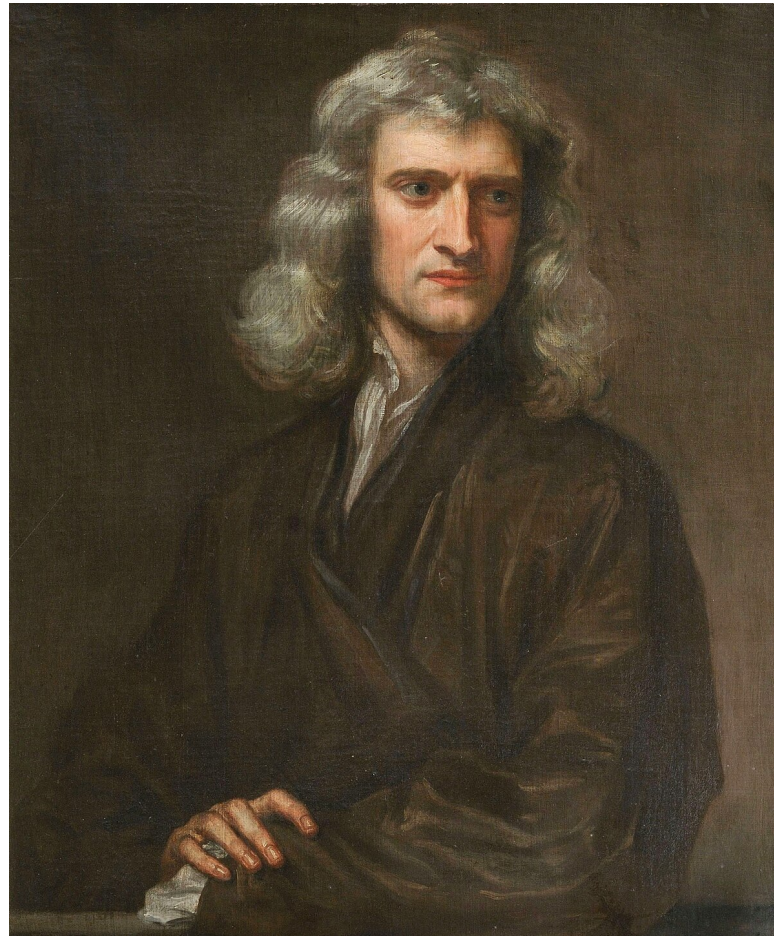
No idea what happens!



Black Hole



Testing Gravity



$$F = \frac{GM_1M_2}{r^2}$$

Large M

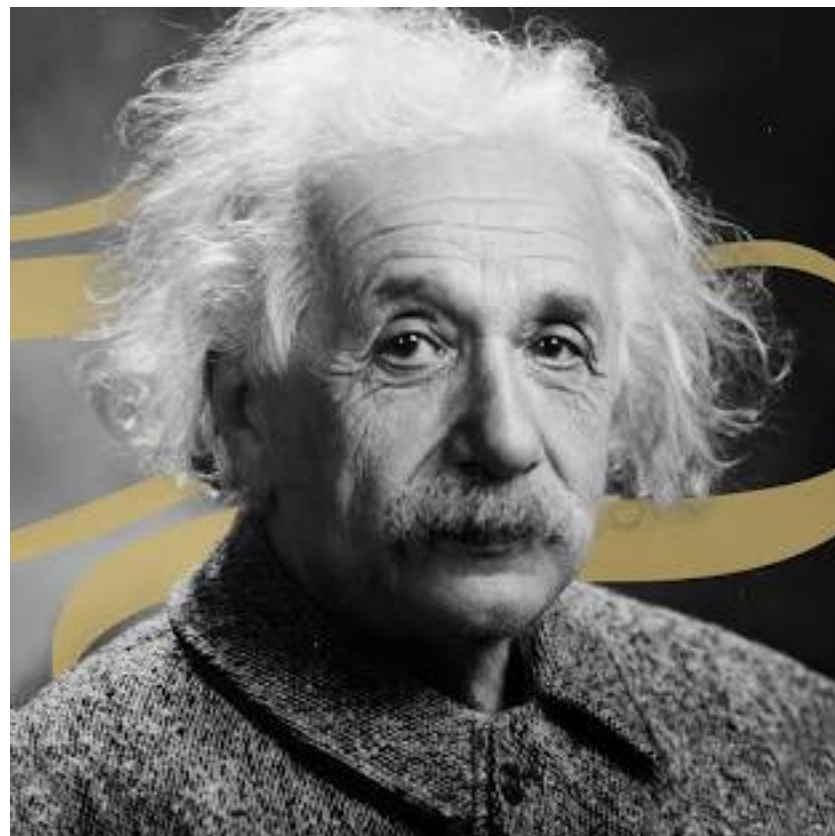
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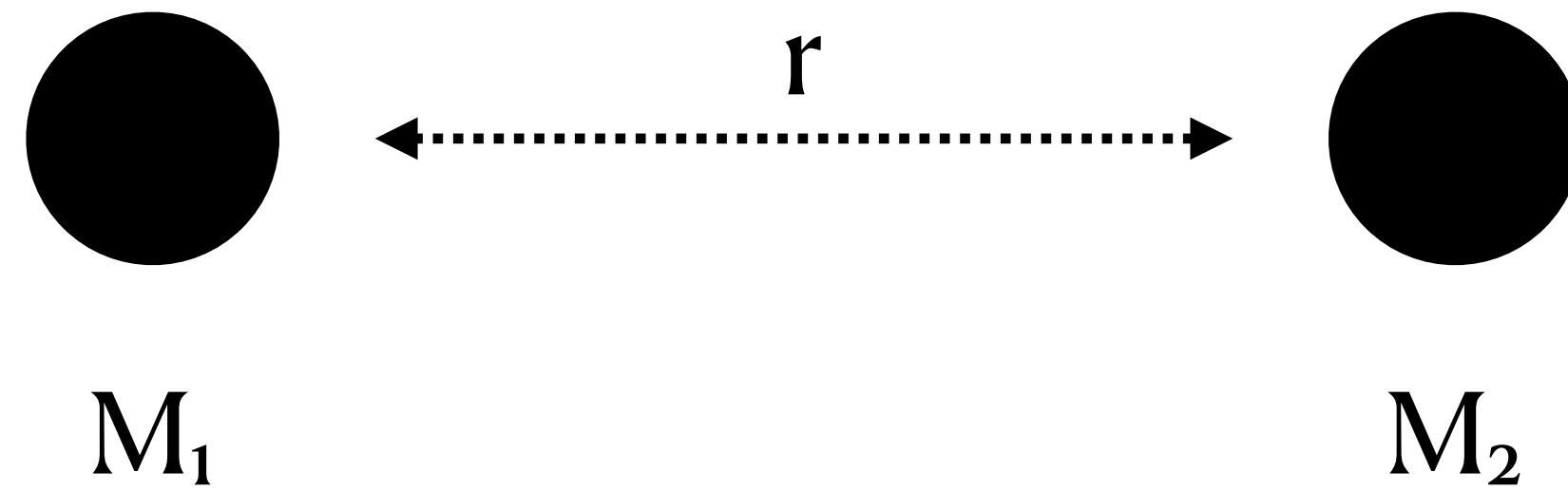
Need new theory - “quantum gravity”



Black Hole



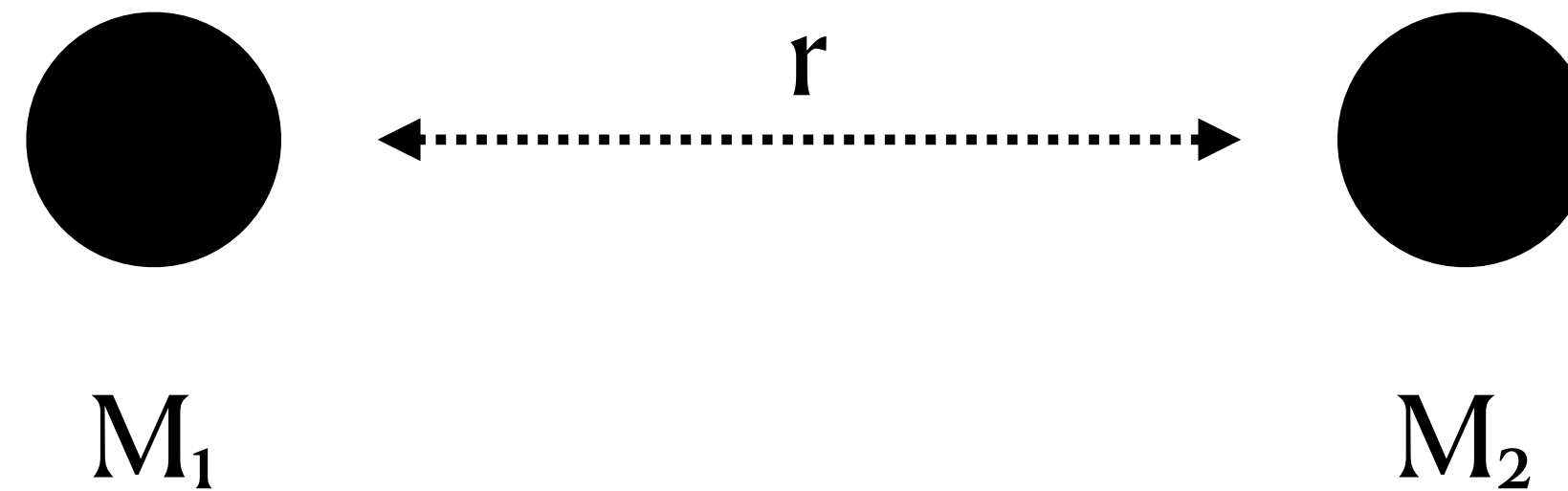
Testing Gravity



Small r ?

$$U = \int_{r_i}^{r_f} d^3r \left(\frac{GM}{r^2} \right)^2 \rightarrow \infty \text{ as } r_i \rightarrow 0$$

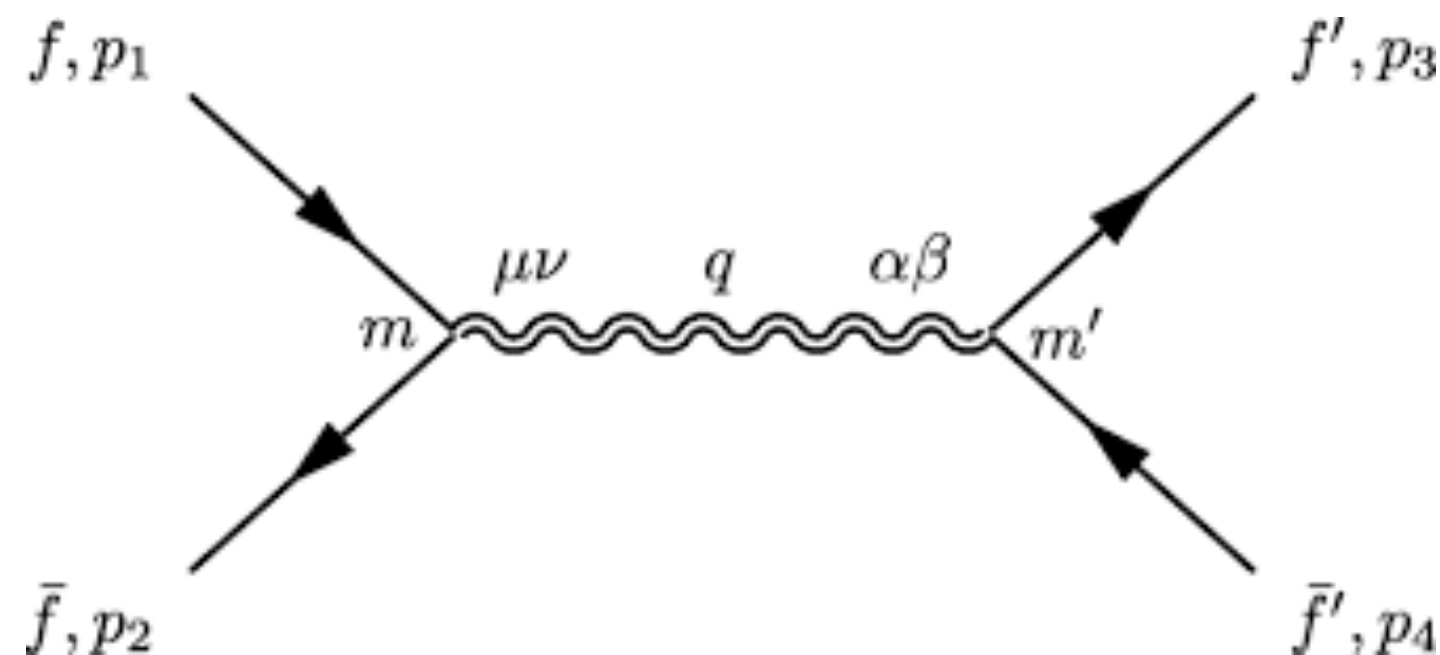
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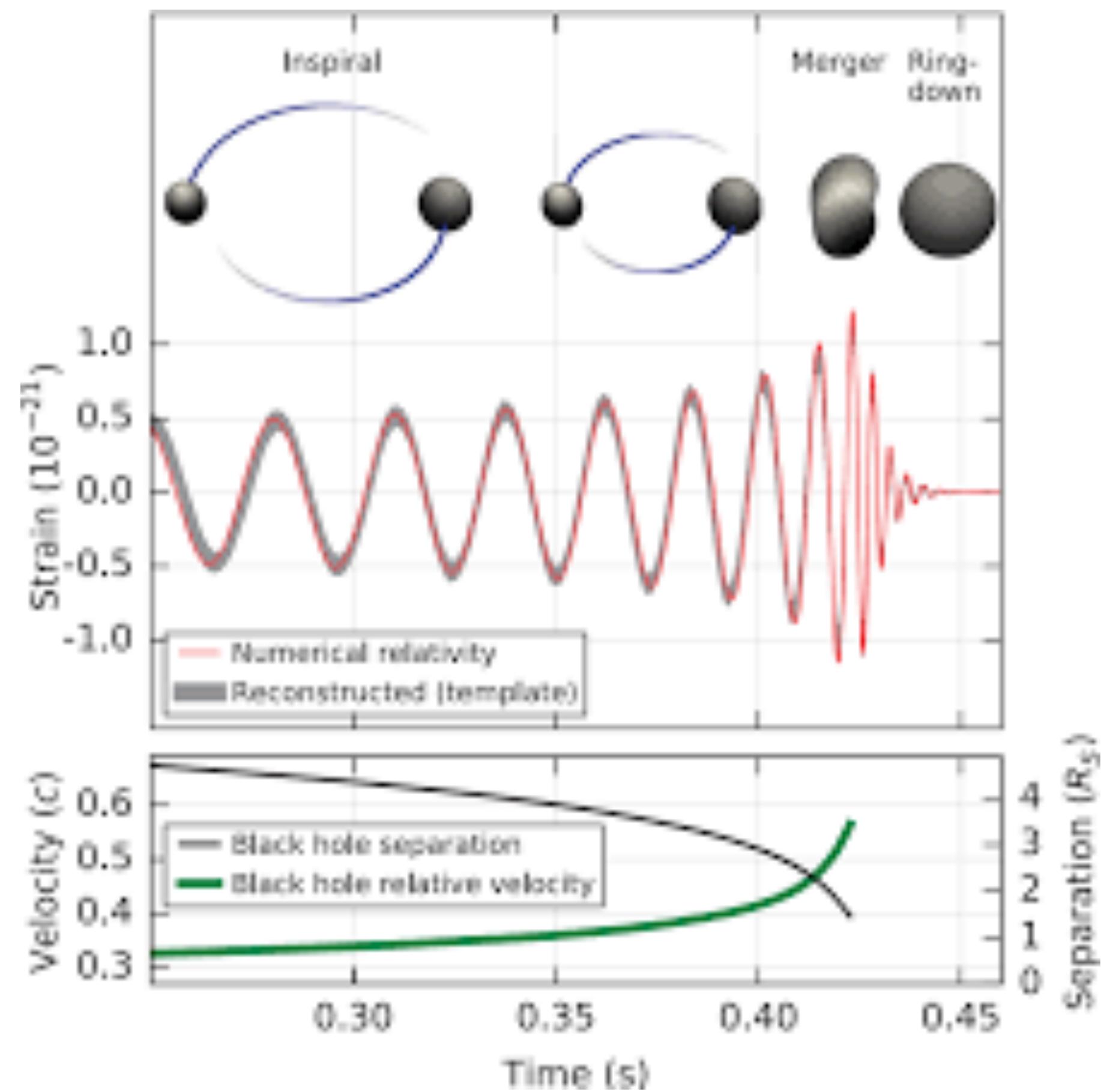


$$\propto GE^2 \rightarrow \infty \text{ as } E \rightarrow \infty$$

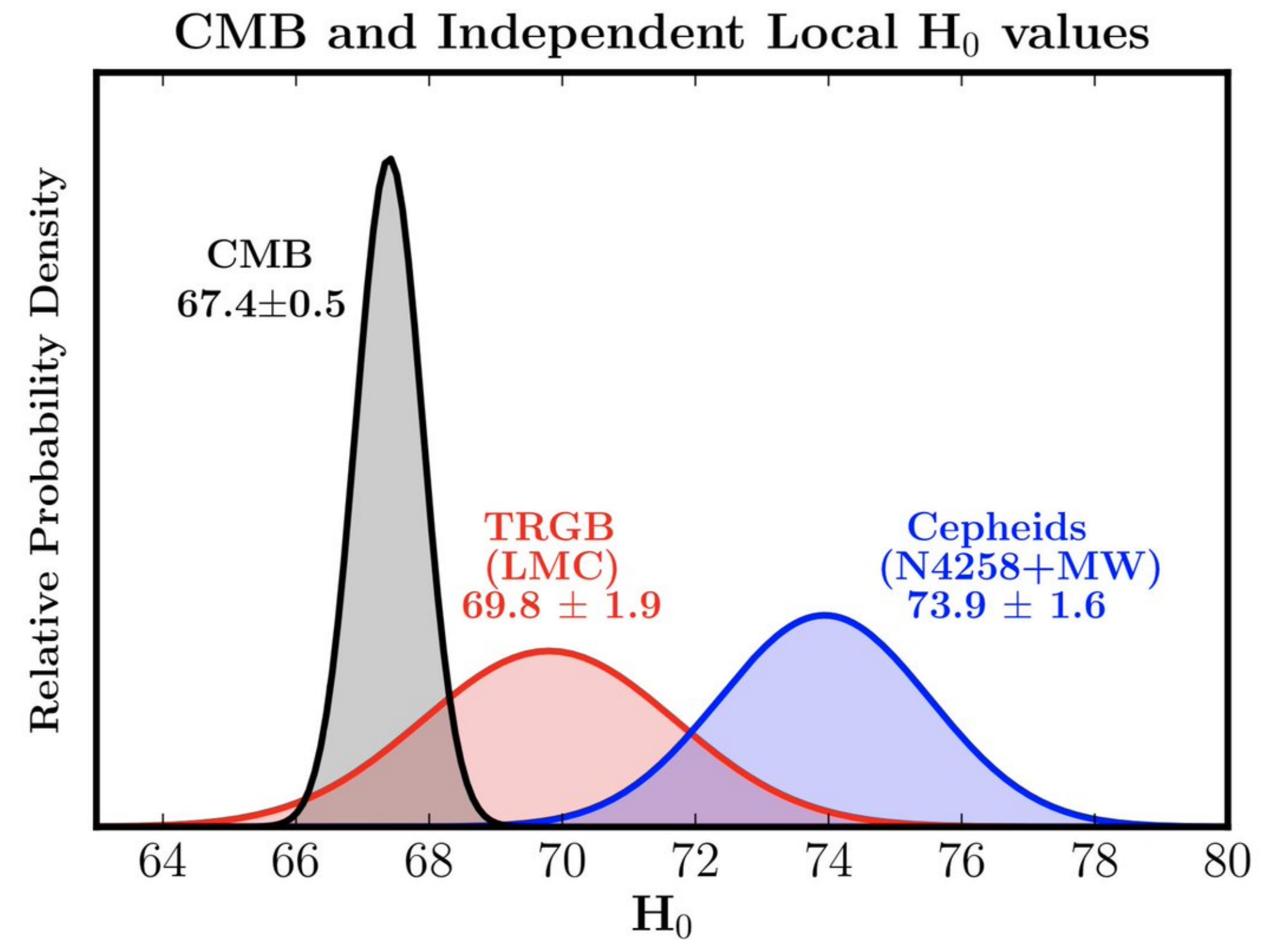
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Testing General Relativity

$$G_{\mu\nu} = T_{\mu\nu}$$



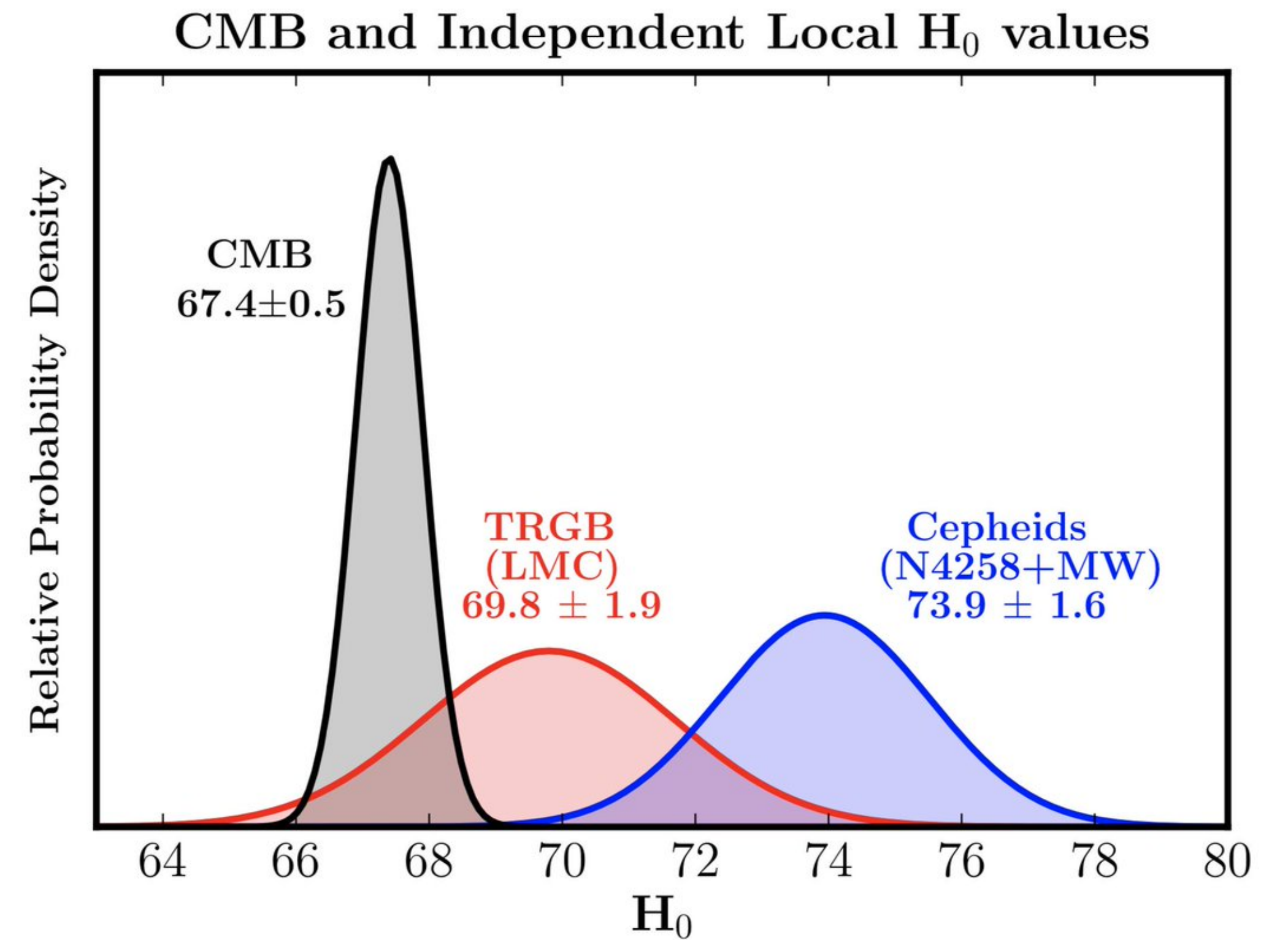
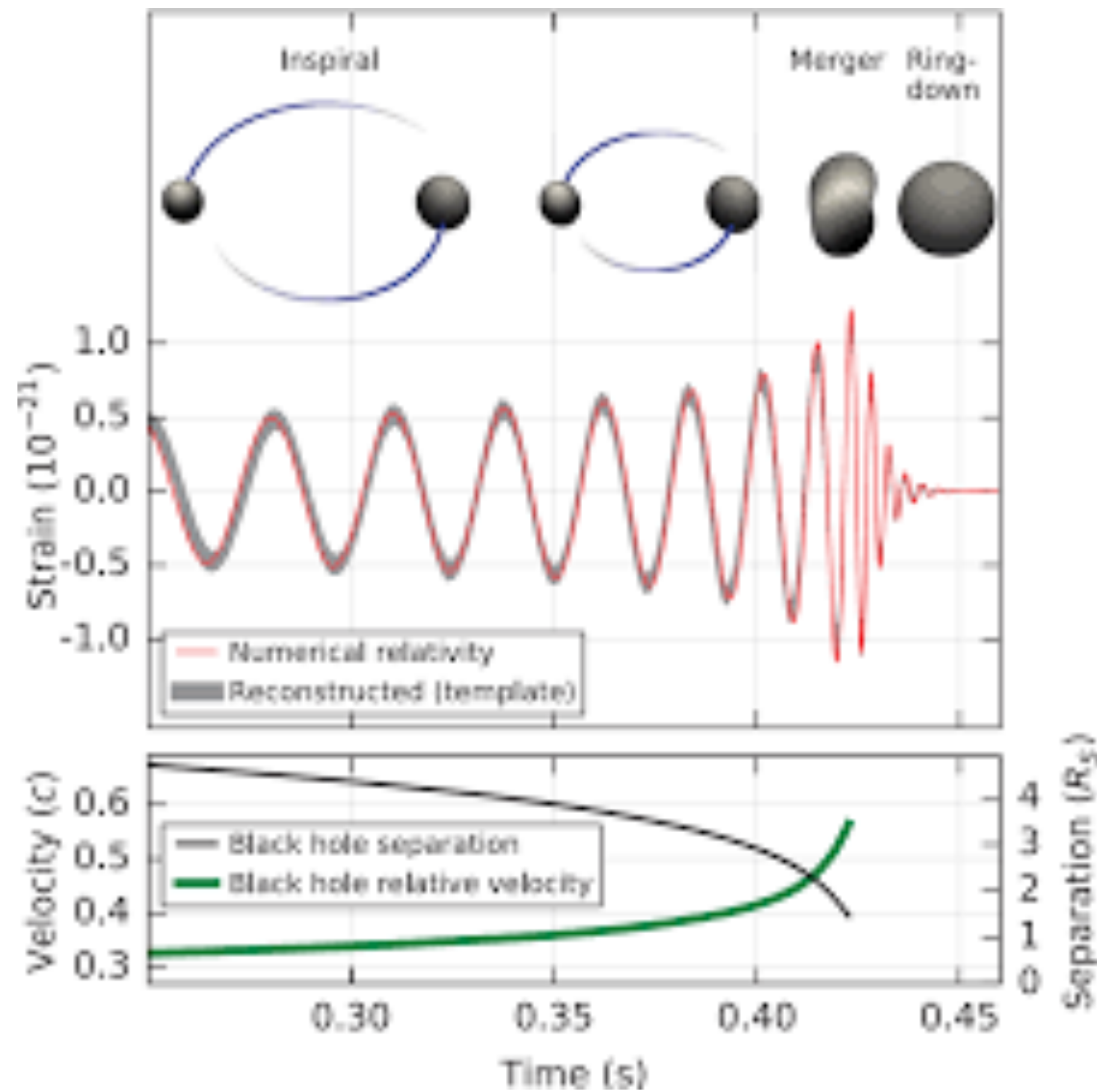
LIGO



Cosmology

Testing General Relativity

$$G_{\mu\nu} = T_{\mu\nu}$$

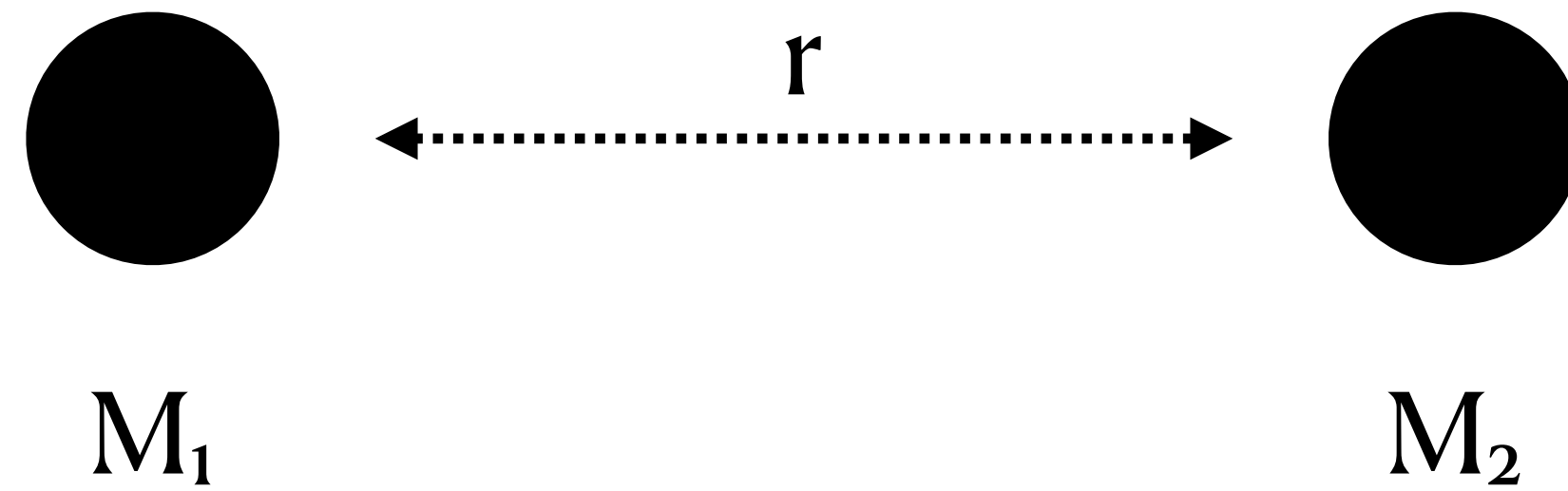


Cosmology

LIGO

Tests of non-linear nature of gravity

Short Distance Tests of Gravity

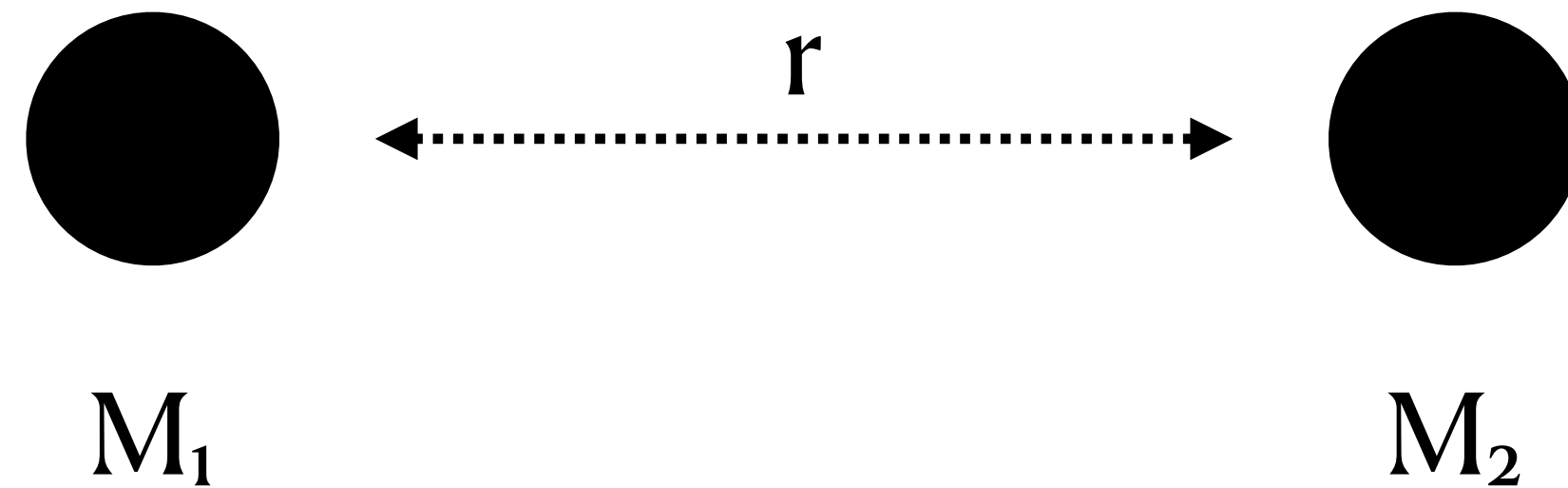


Small r ?

$$F = \frac{GM_1 M_2}{r^2}$$

How well do we know that this law is correct experimentally?

Short Distance Tests of Gravity



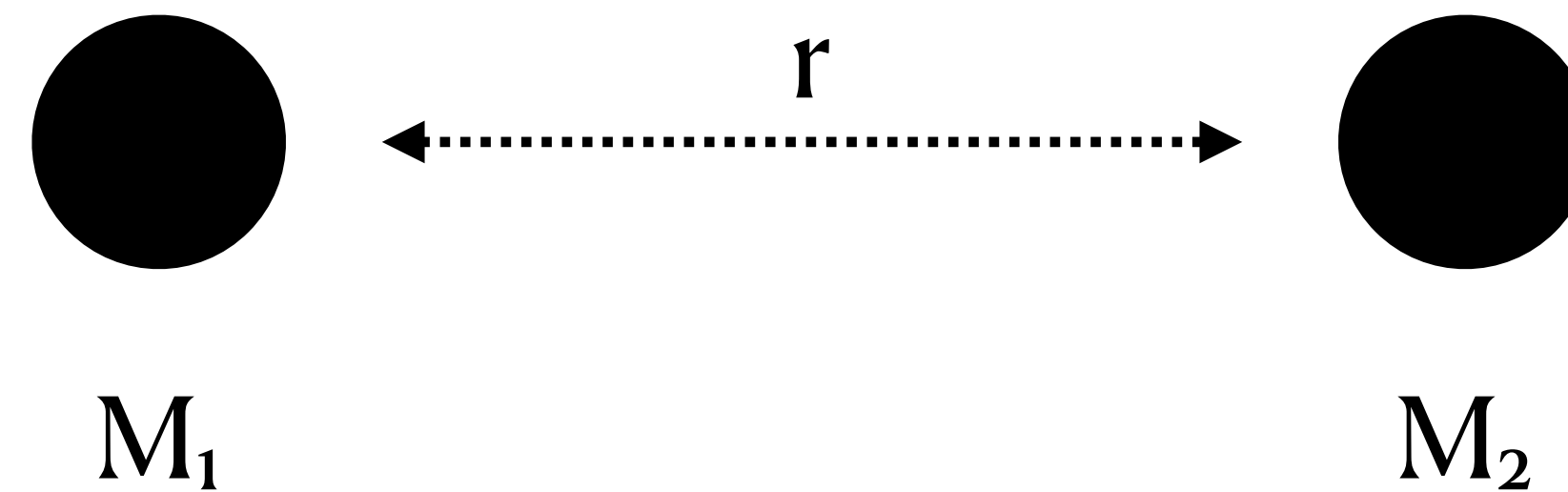
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Far from quantum gravity regime

Short Distance Tests of Gravity



Small r ?

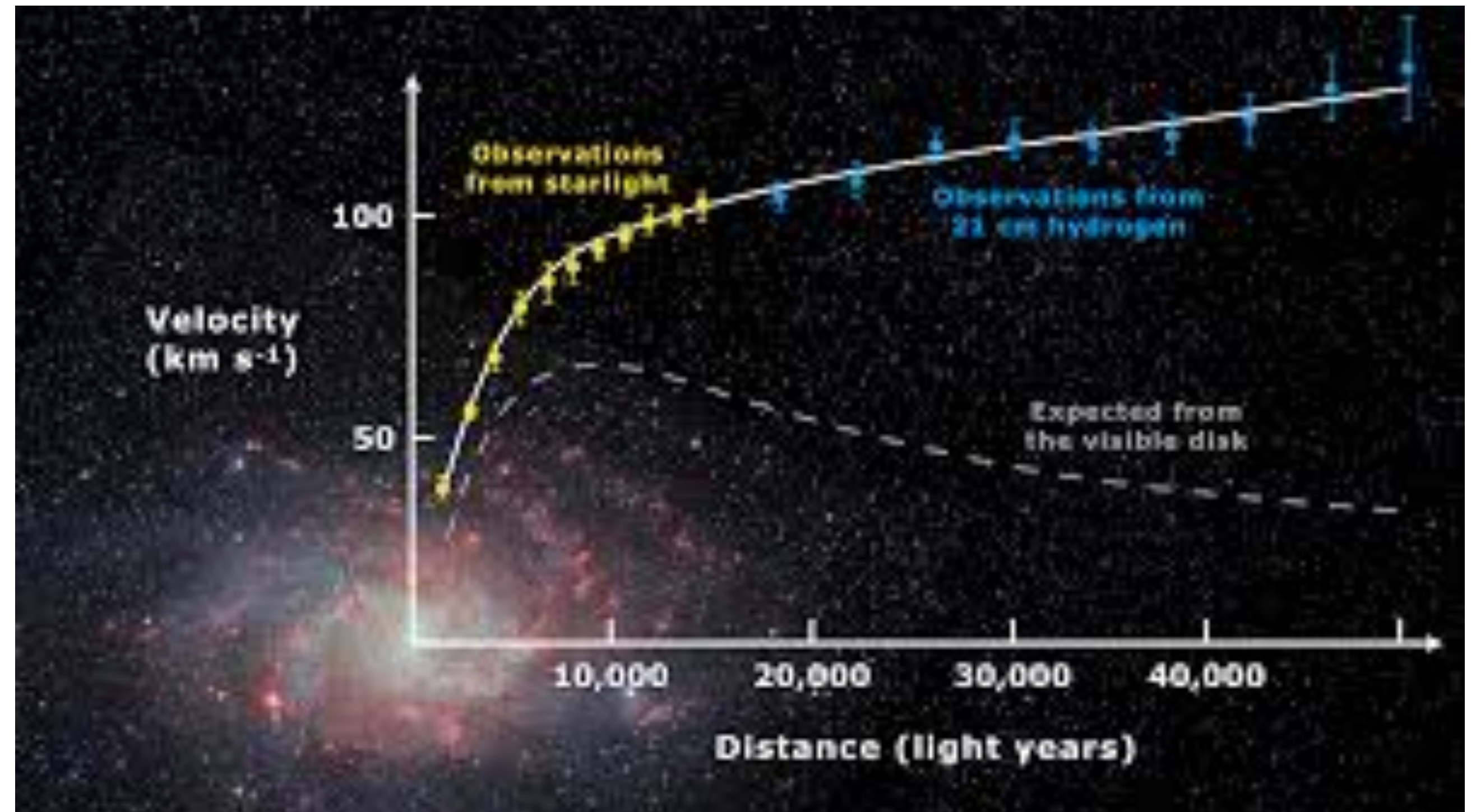
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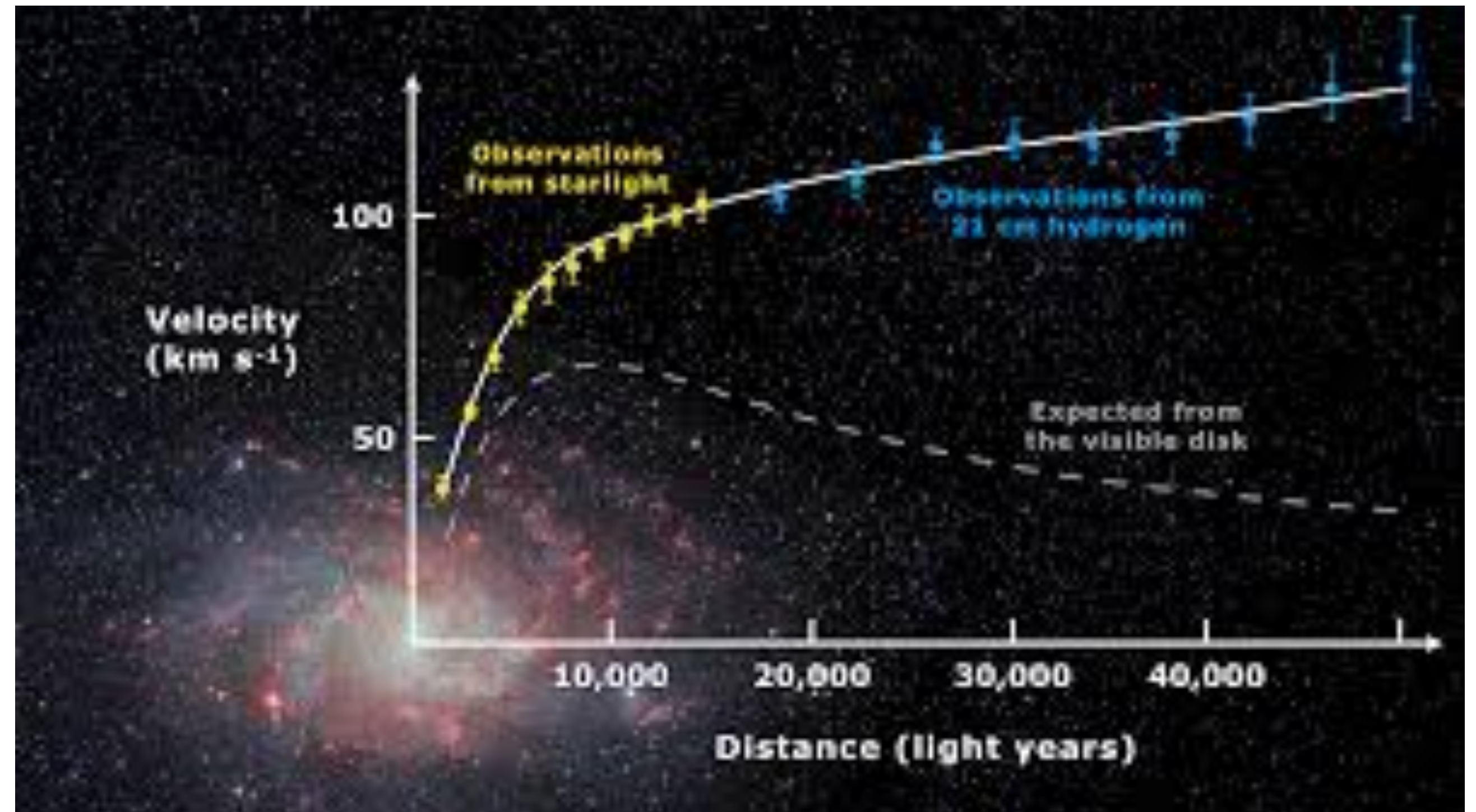
Far from quantum gravity regime

What can we find?

Testing Gravity

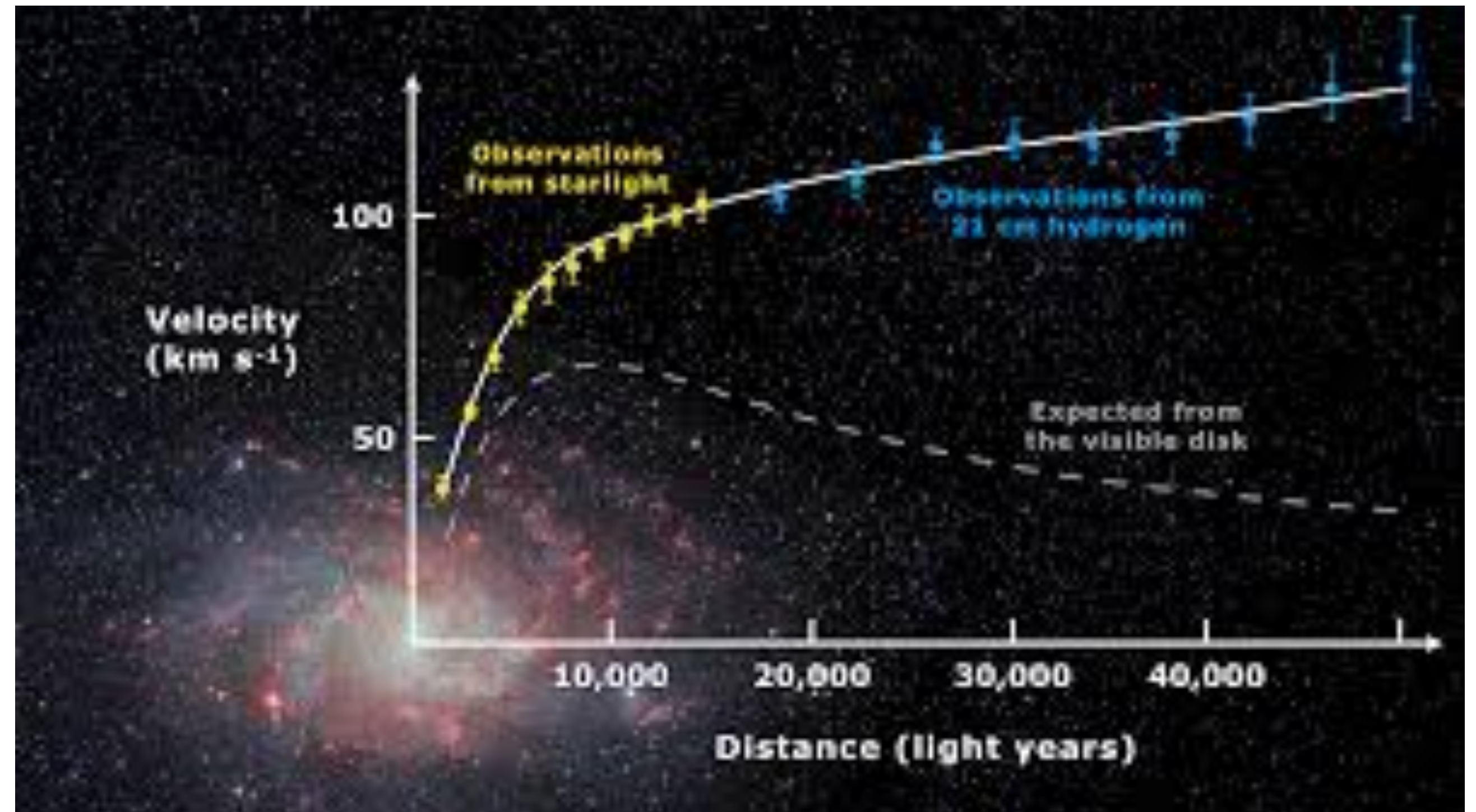


Testing Gravity



No reason to expect deviations from gravity at long distance - but found dark matter!

Testing Gravity



No reason to expect deviations from gravity at long distance - but found dark matter!

What can we hope to find at short distance?

New Short Distance Forces

Light bosonic particles motivated by BSM Physics
(e.g. radions, moduli, relaxions)

$$\mathcal{L} \supset y_q \phi \bar{q} q + \frac{\phi}{f_\gamma} F_{\mu\nu}^2 + \frac{\phi}{f_g} G_{\mu\nu}^2 + \frac{\tilde{h}_{\mu\nu}}{f_T} F^\mu{}_\sigma F^{\nu\sigma} + g \phi h^2 + \frac{m_\phi^2}{2} \phi^2$$

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How do we find them?

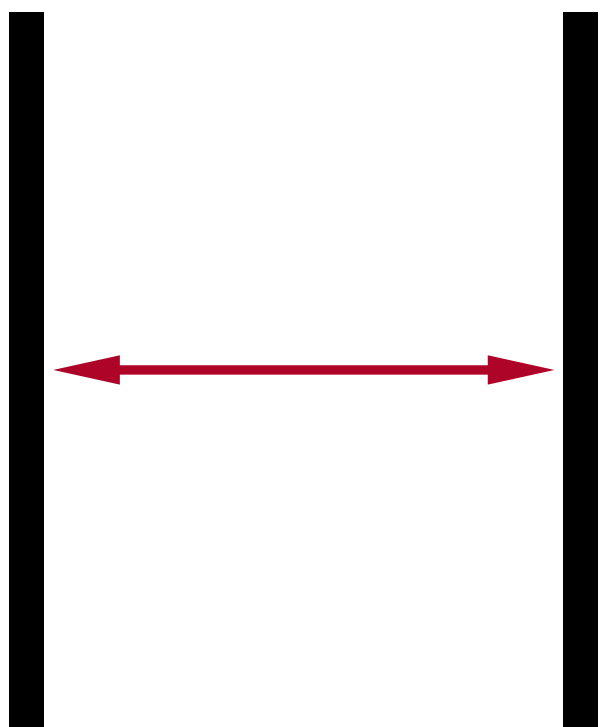
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How do we find them?

Take two objects, measure anomalous forces between them



$$F = \alpha \frac{G m_p^2}{r^2} e^{-\frac{r}{\lambda}}$$

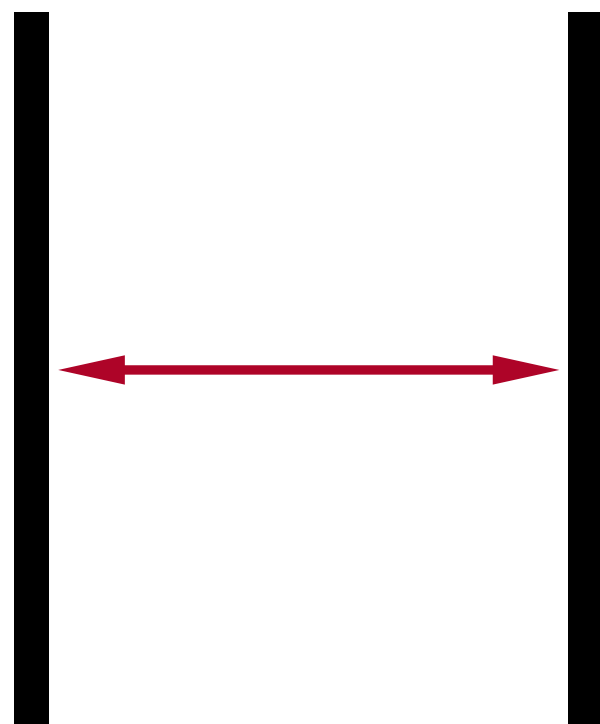
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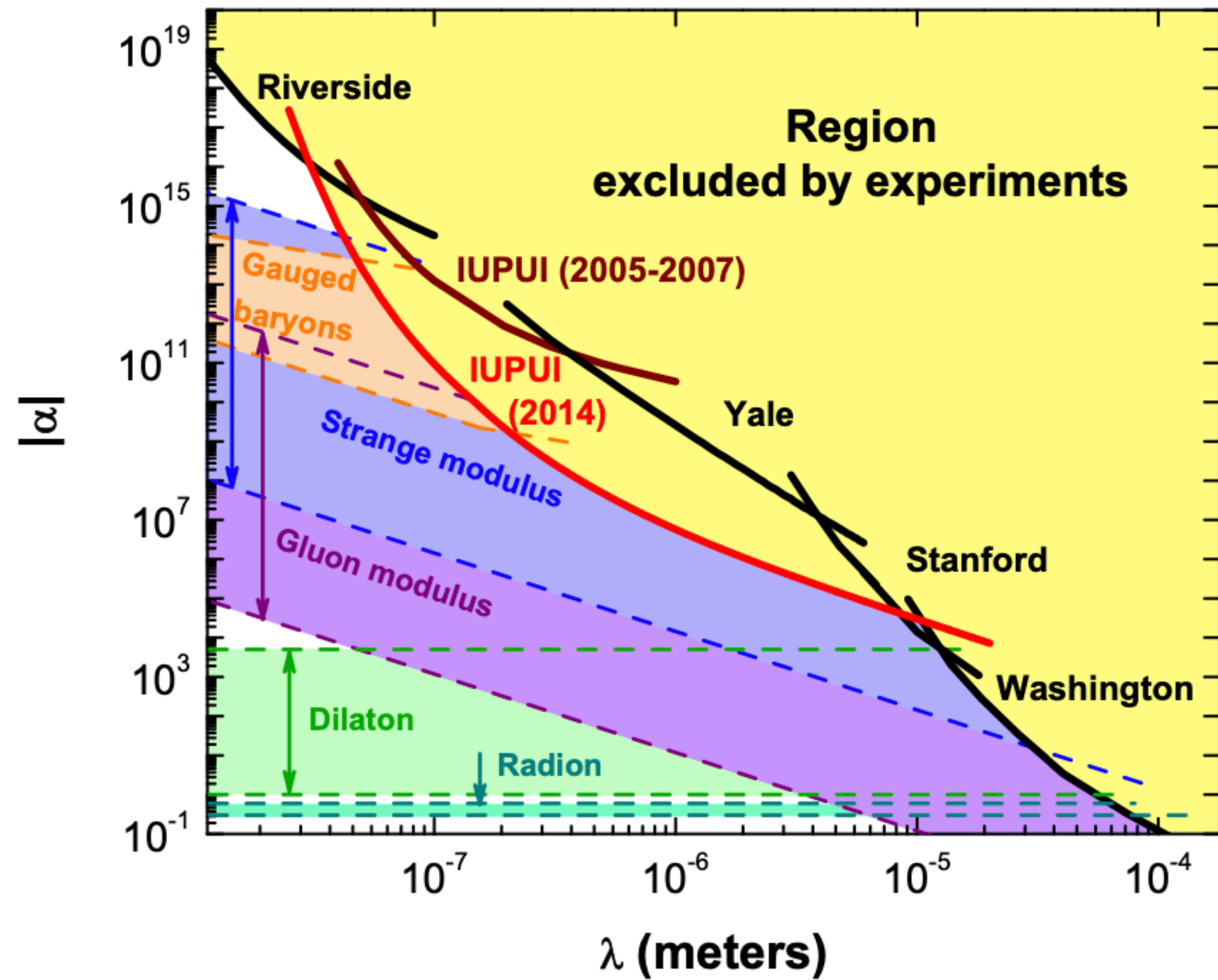
Take two objects, measure anomalous forces between them



$$F = \alpha \frac{G m_p^2}{r^2} e^{-\frac{r}{\lambda}}$$

Measure Relative Acceleration

Where are we?

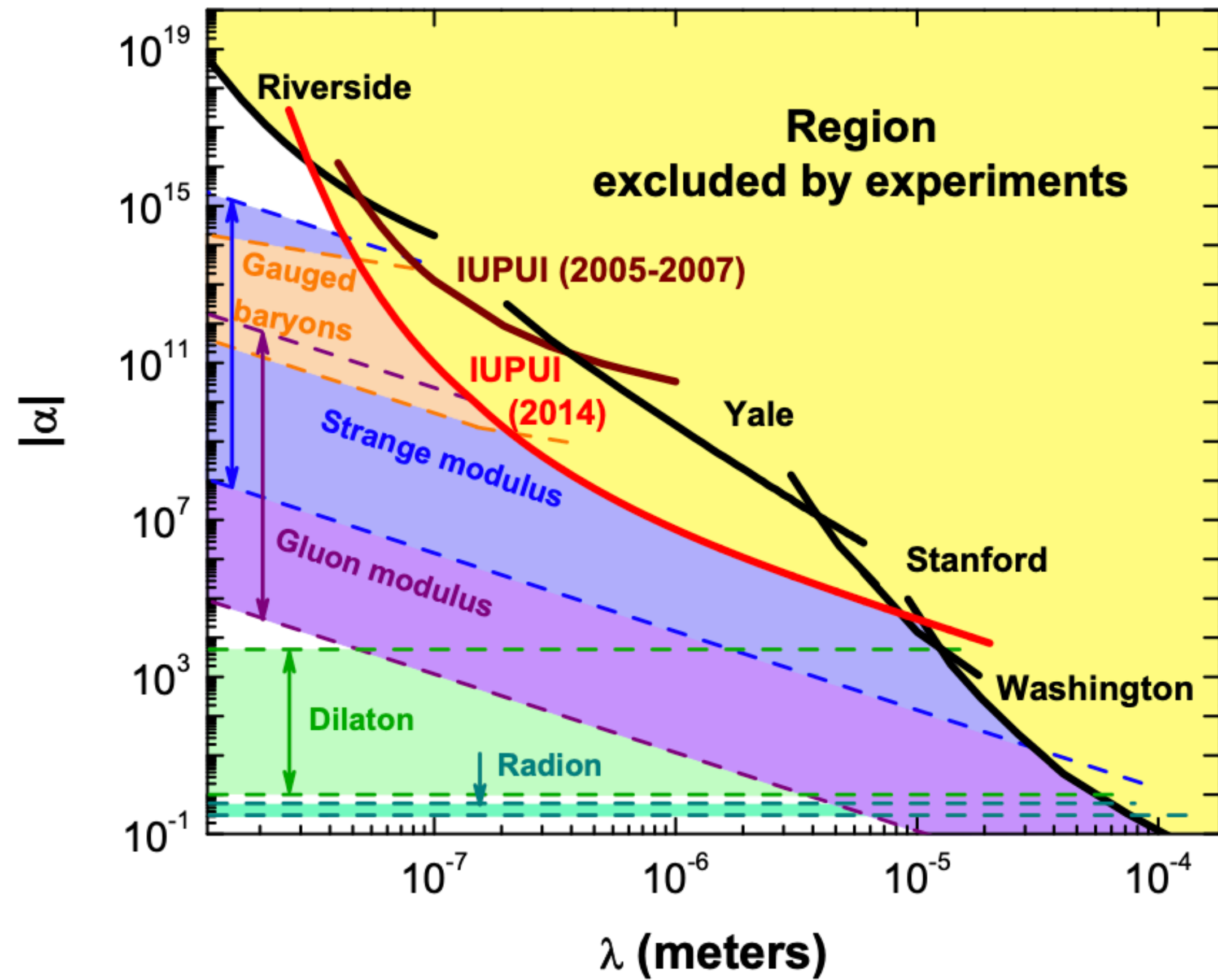


Very strong constraints at long ($> \mu\text{m}$) distances

Sensitivity rapidly drops at short ($< \mu\text{m}$) distances

$$F = \alpha \frac{G m_p^2}{r^2} e^{-\frac{r}{\lambda}}$$

Where are we?



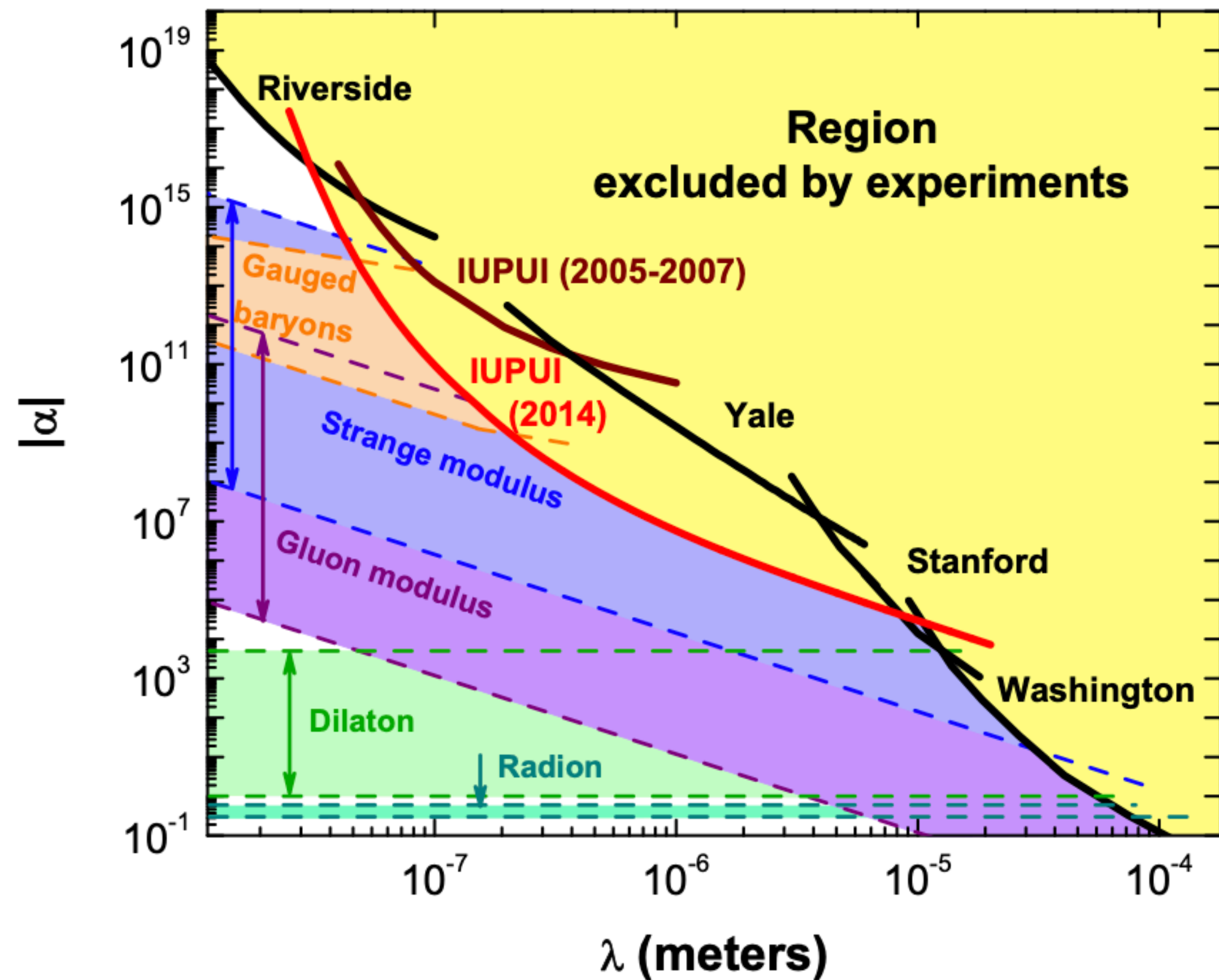
Very strong constraints at long ($> \mu\text{m}$) distances

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Why?

$$F = \alpha \frac{G m_p^2}{r^2} e^{-\frac{r}{\lambda}}$$

Where are we?



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Why?

Short Range \Rightarrow Objects need to be close

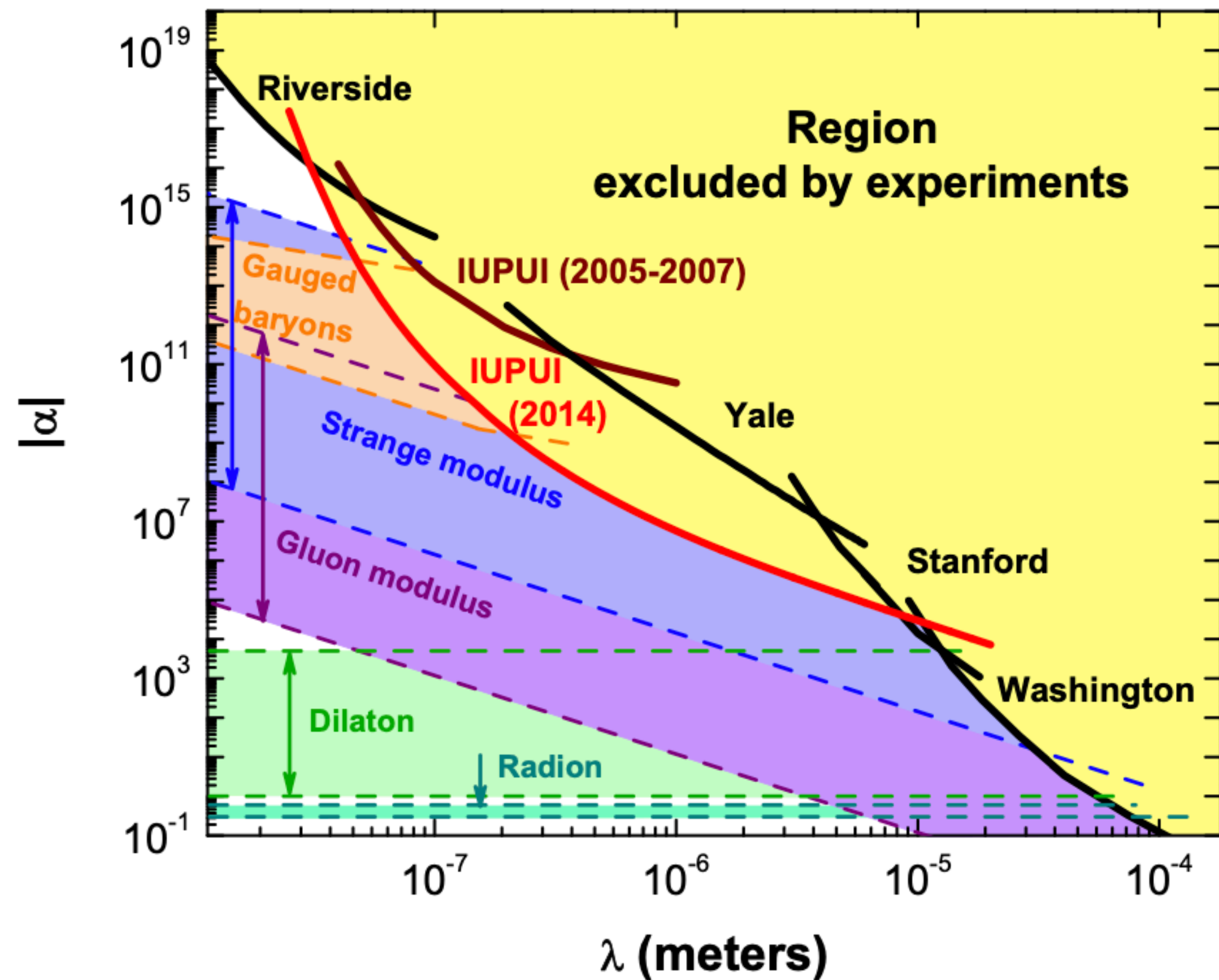
Electromagnetism \gg New Physics

Short Range \Rightarrow Only material within λ affected

Need to deal with thin objects with high precision

$$F = \alpha \frac{G m_p^2}{r^2} e^{-\frac{r}{\lambda}}$$

Where are we?



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Progress?

Outline

1. Mossbauer Effect

2. Setup

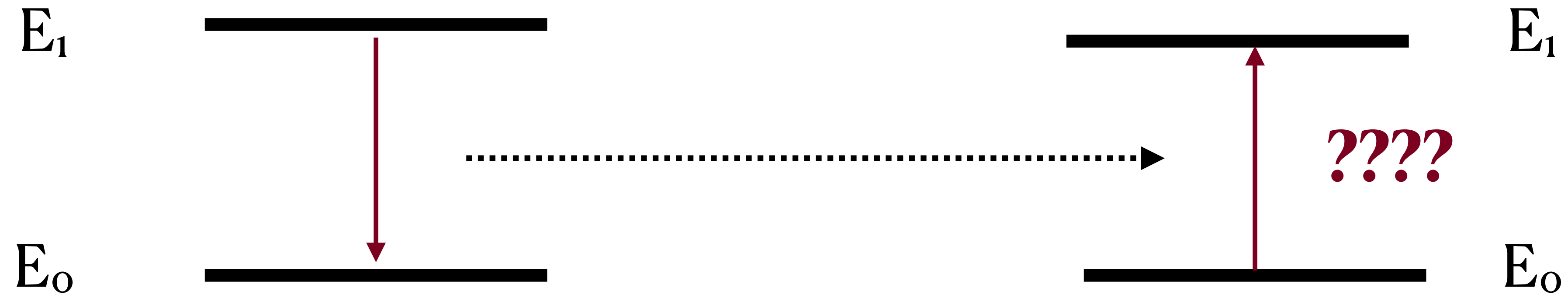
3. Backgrounds

4. Sensitivity

5. Synchrotron Light Sources?

6. Conclusions

Mossbauer Effect



Excited nuclear state decays via γ emission

Can the γ be reabsorbed?

Mossbauer Effect



Excited nuclear state decays via γ emission

Can the γ be reabsorbed?

Issue: Small nuclear cross-sections

Efficient reabsorption only possible on resonance

Mossbauer Effect



Excited nuclear state decays via γ emission

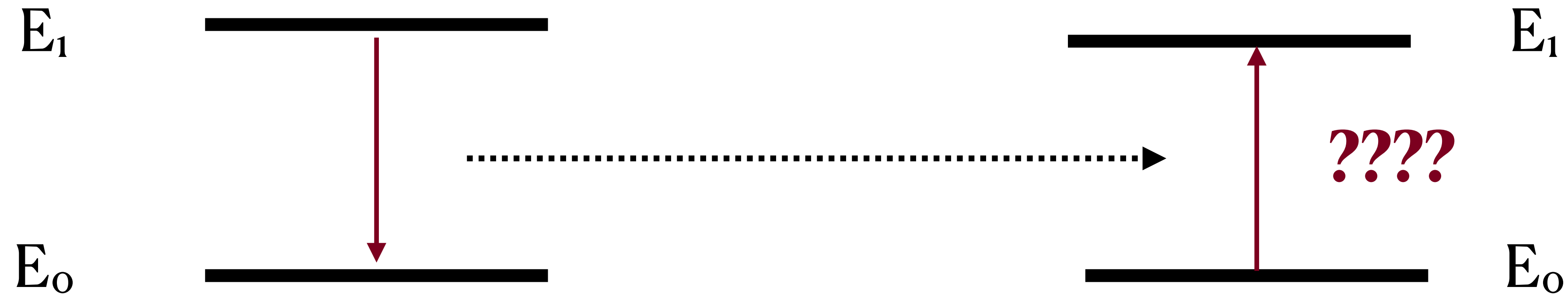
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Isn't emitted γ at transition energy? Automatically Resonant?

Mossbauer Effect



Excited nuclear state decays via γ emission

Can the γ be reabsorbed?

Issue: Small nuclear cross-sections

Efficient reabsorption only possible on resonance

Isn't emitted γ at transition energy? Automatically Resonant?

No : Recoiling nucleus takes energy, γ outside narrow width

Mossbauer Effect



Small enough E_γ , entire lattice recoils!

Negligible lattice kinetic energy - monochromatic E_γ

Resonant Reabsorption possible!

Mossbauer Effect



Small enough E_γ , entire lattice recoils!

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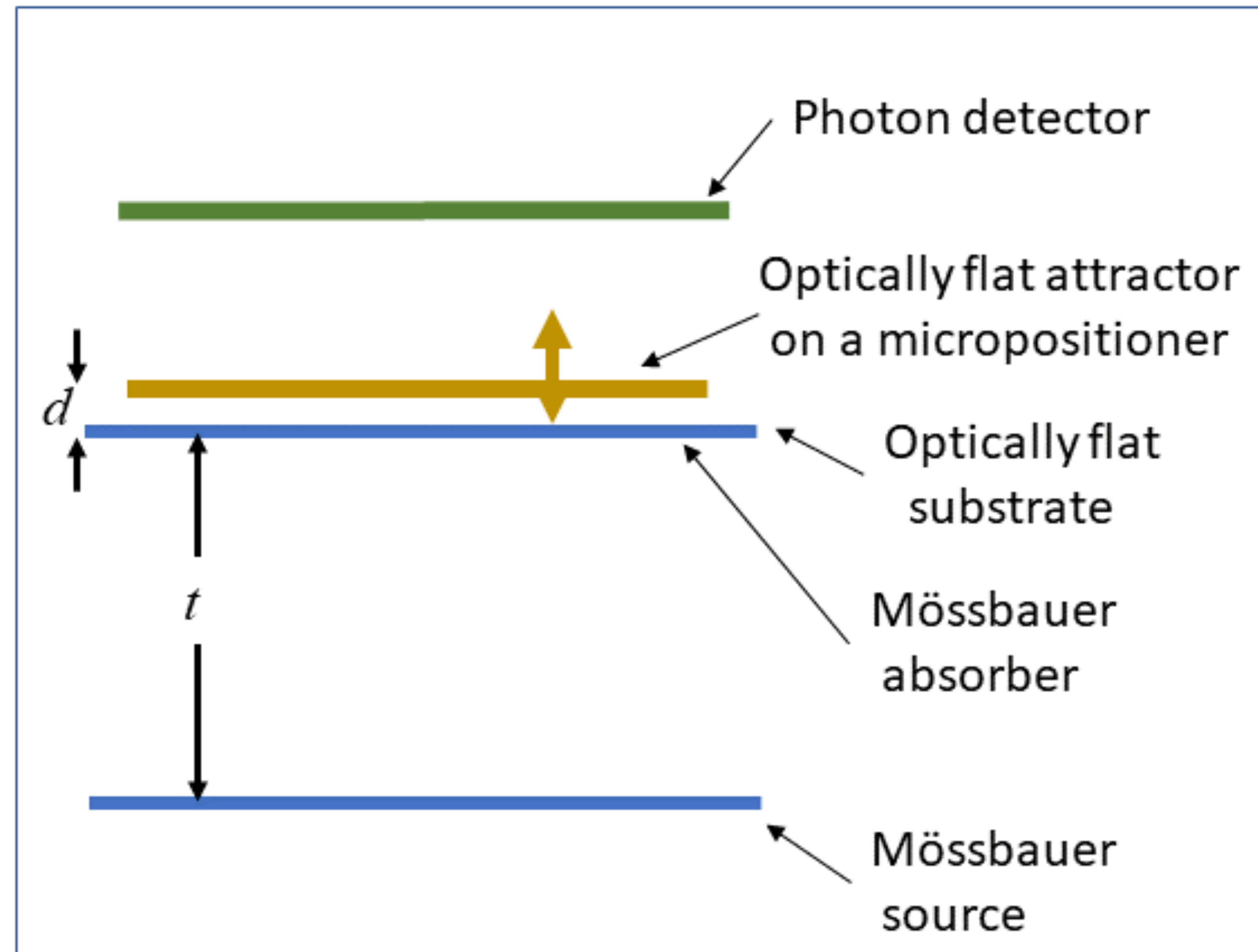
Resonant Reabsorption possible!

Narrow Nuclear Lines \Rightarrow High Sensitivity to energy shifts

Setup

$$\mathcal{L} \supset y_q \phi \bar{q} q + \frac{\phi}{f_\gamma} F_{\mu\nu}^2 + \frac{\phi}{f_g} G_{\mu\nu}^2 + \frac{\tilde{h}_{\mu\nu}}{f_T} F^\mu{}_\sigma F^{\nu\sigma} + g\phi h^2 + \frac{m_\phi^2}{2} \phi^2$$

**New interaction
shifts nuclear energy**

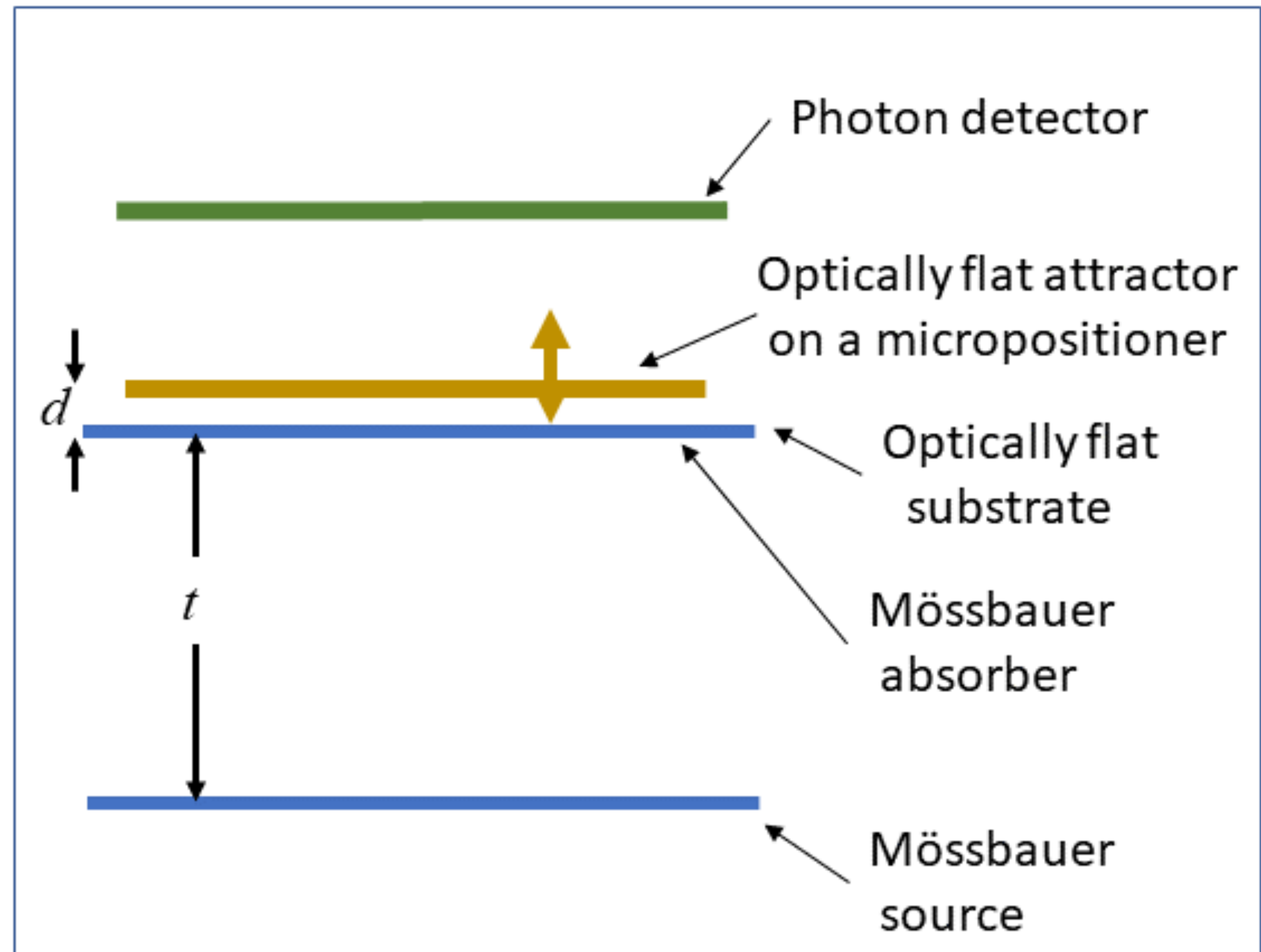


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**New interaction
shifts nuclear energy**

**Resonant
Reabsorption as a
function of d**



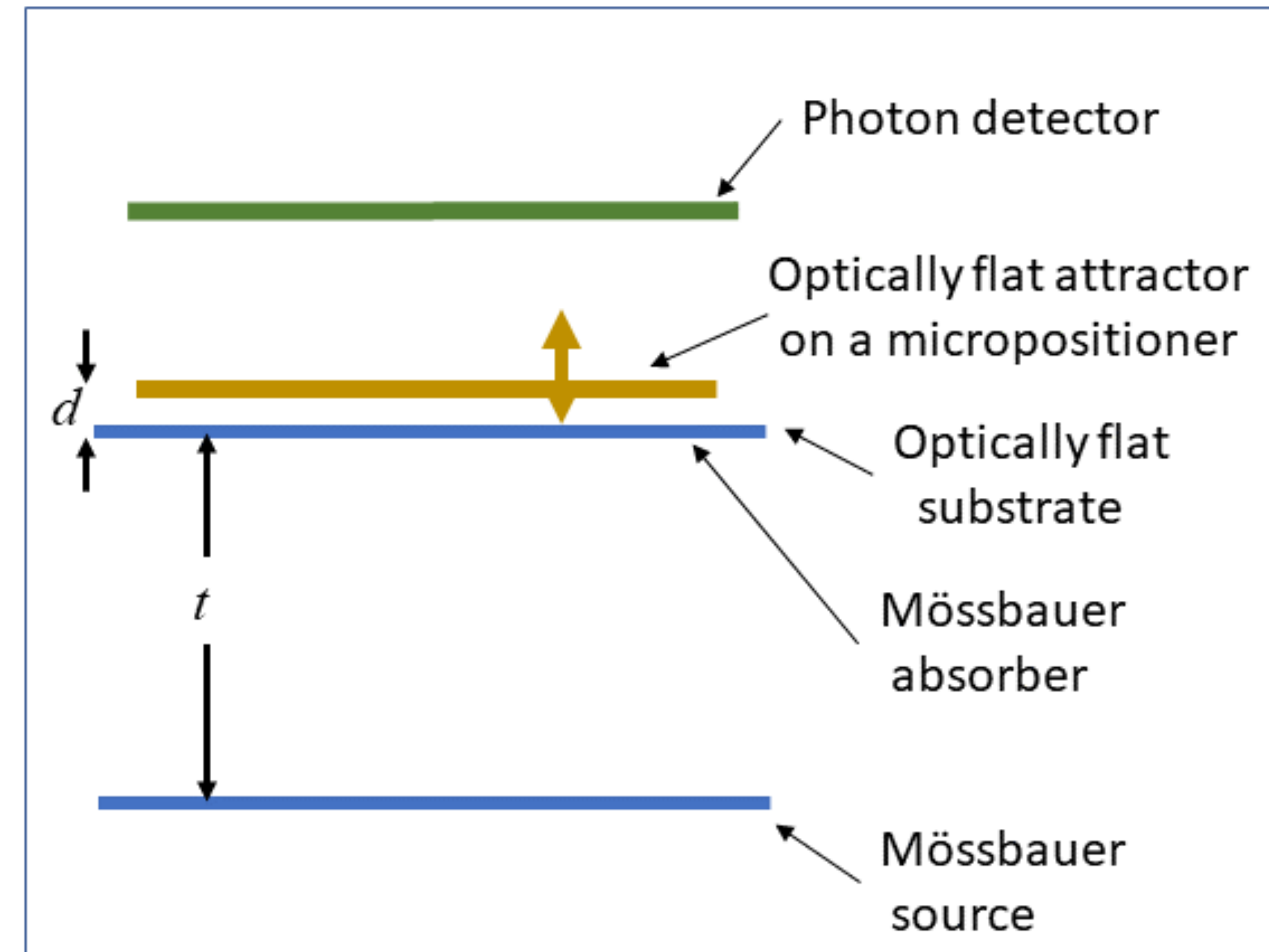
Backgrounds

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Electromagnetism?

Needs to change nuclear transition energies

Suppressed by small nuclear moments,
electron shielding



Backgrounds

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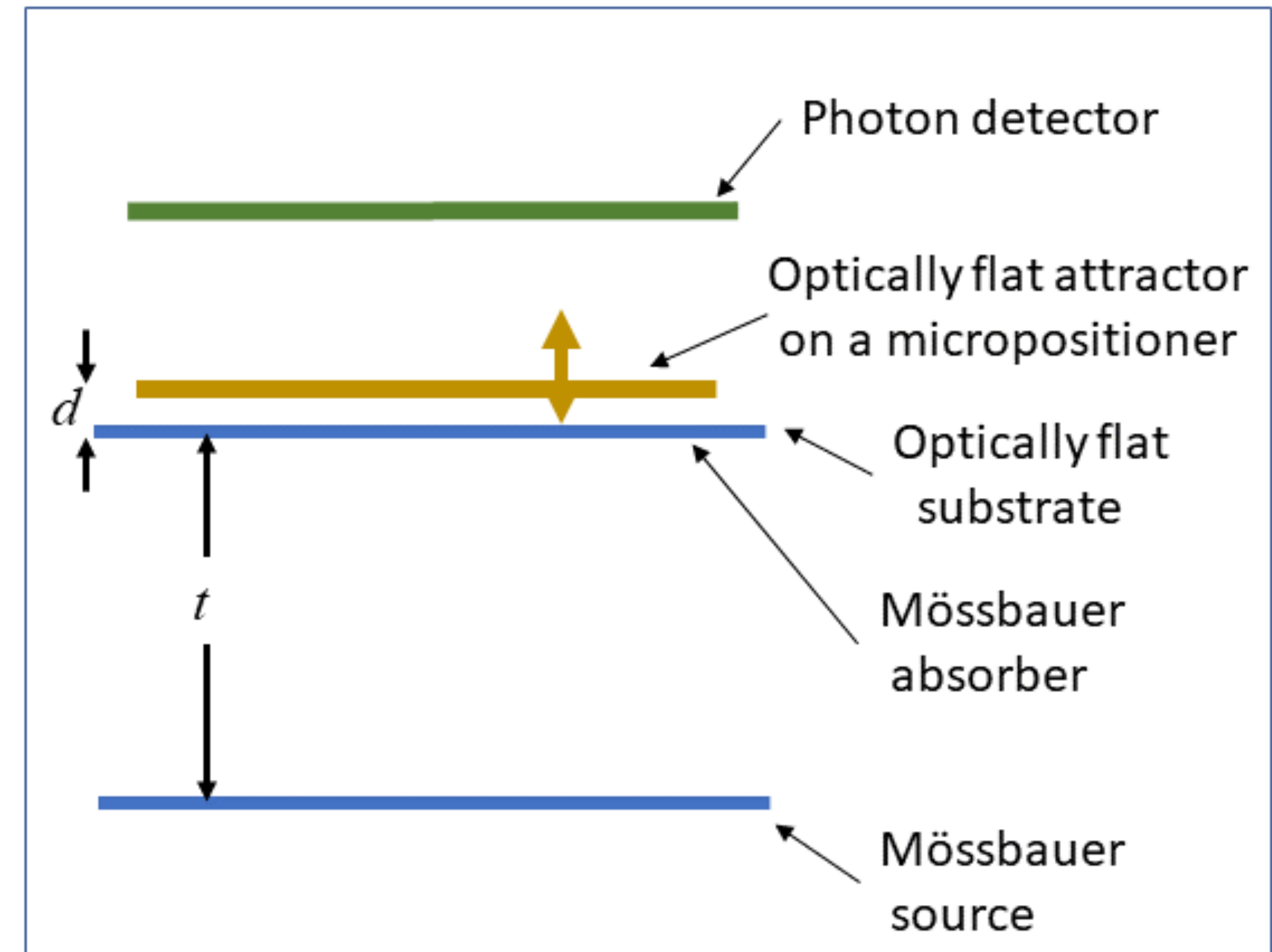
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Unpolarized nuclear spin => effects average
down

Signal from new scalar and tensor
interactions are not suppressed!



Backgrounds

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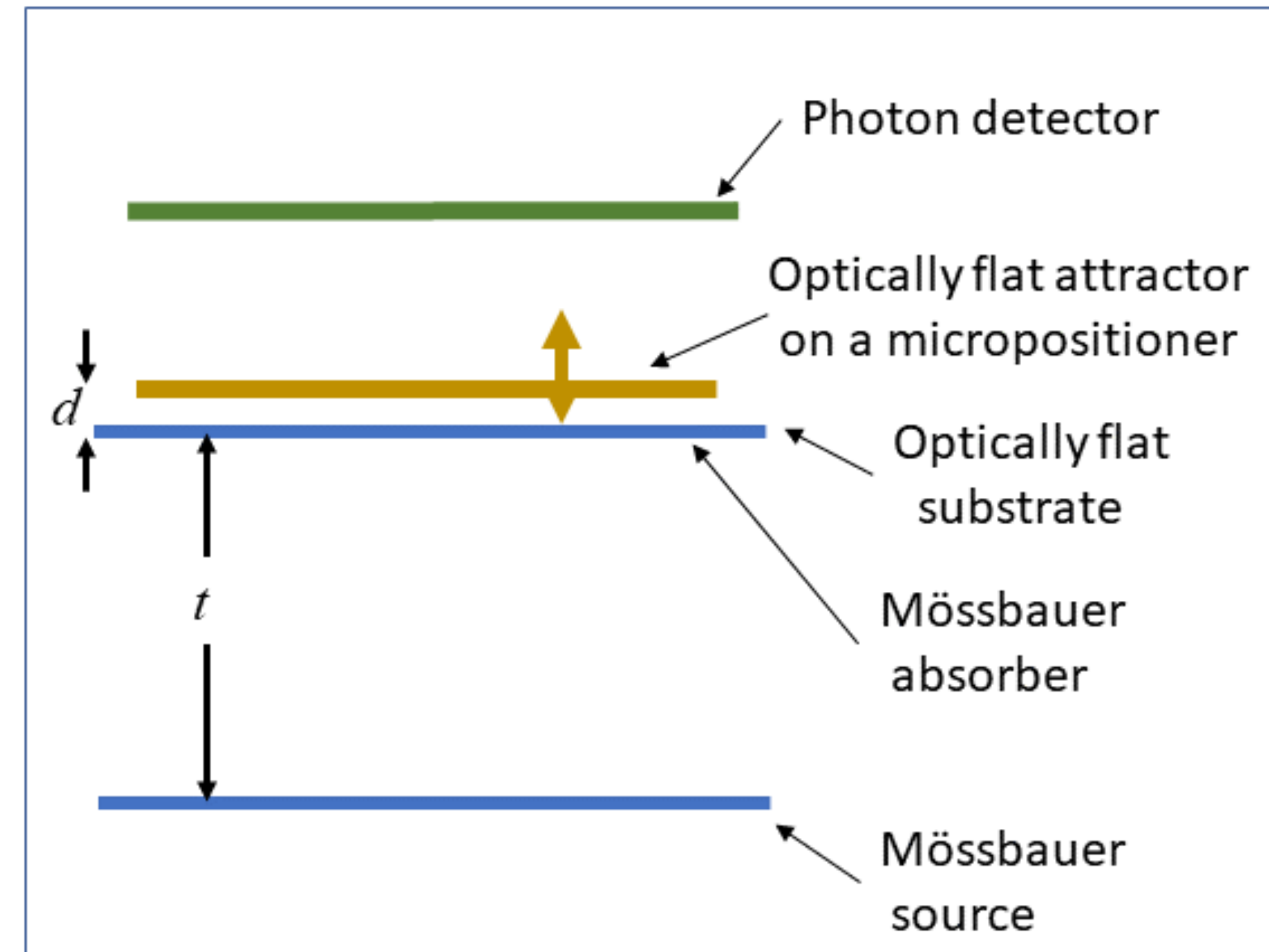
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Unpolarized nuclear spin => effects average
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First order effects irrelevant

Leading Background: Chemical Shift from Casimir

Sensitivity

$$\mathcal{L} \supset y_q \phi \bar{q} q + \frac{\phi}{f_\gamma} F_{\mu\nu}^2 + \frac{\phi}{f_g} G_{\mu\nu}^2 + \frac{\tilde{h}_{\mu\nu}}{f_T} F^\mu{}_\sigma F^{\nu\sigma} + g \phi h^2 + \frac{m_\phi^2}{2} \phi^2$$

For given coupling, compute energy shift ΔE

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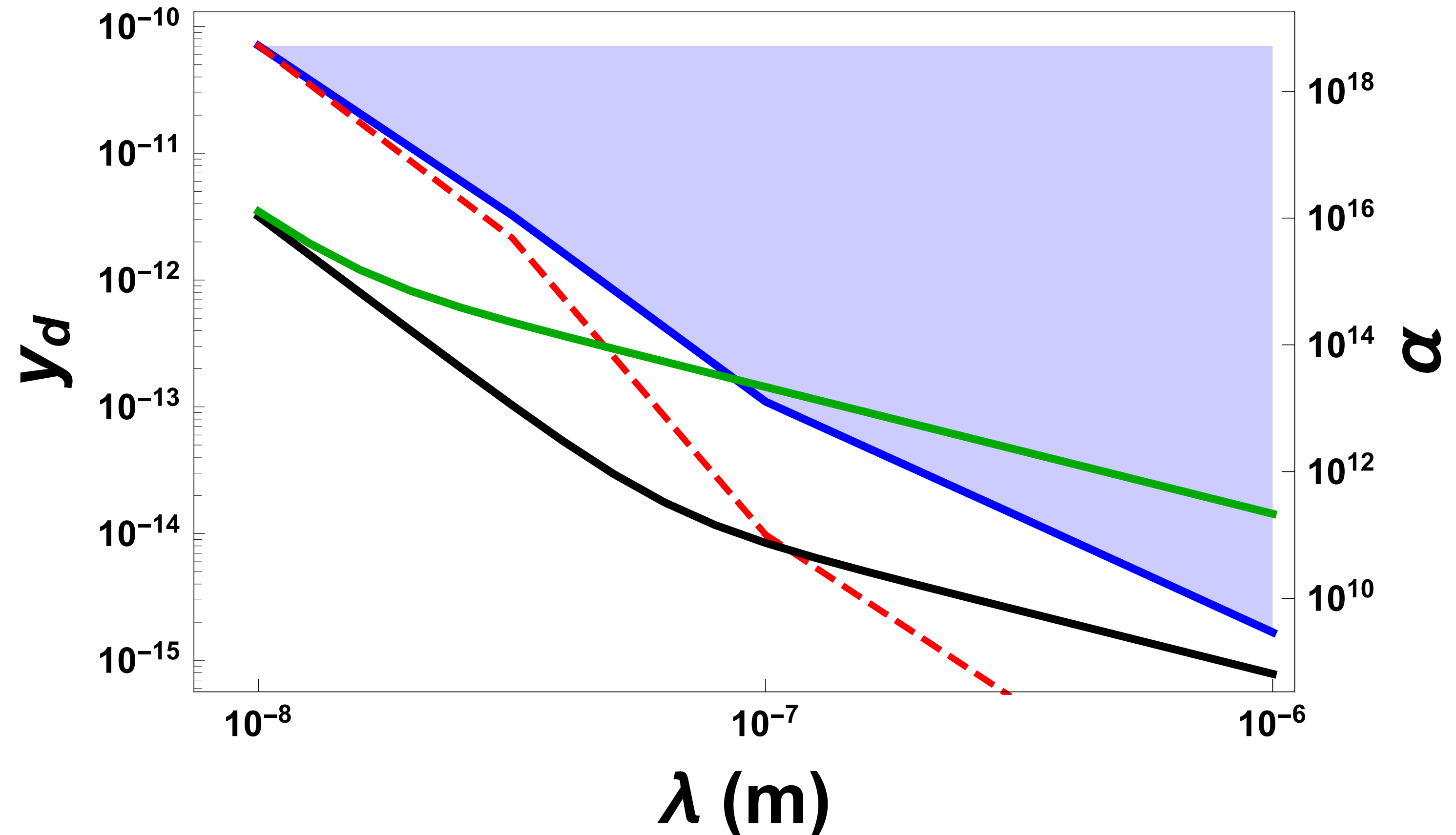
For given coupling, compute energy shift ΔE

$$\Delta E = \frac{\Gamma}{\sqrt{N_\gamma}}$$

$$^{57}\text{Fe} \quad \Delta E = 10^{-15} \text{ eV}$$

$$^{181}\text{Tl} \quad \Delta E = 10^{-17} \text{ eV}$$

$$N_\gamma = 3 \times 10^{14}$$



Sensitivity

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For given coupling, compute energy shift ΔE

$$\Delta E = \frac{\Gamma}{\sqrt{N_\gamma}}$$

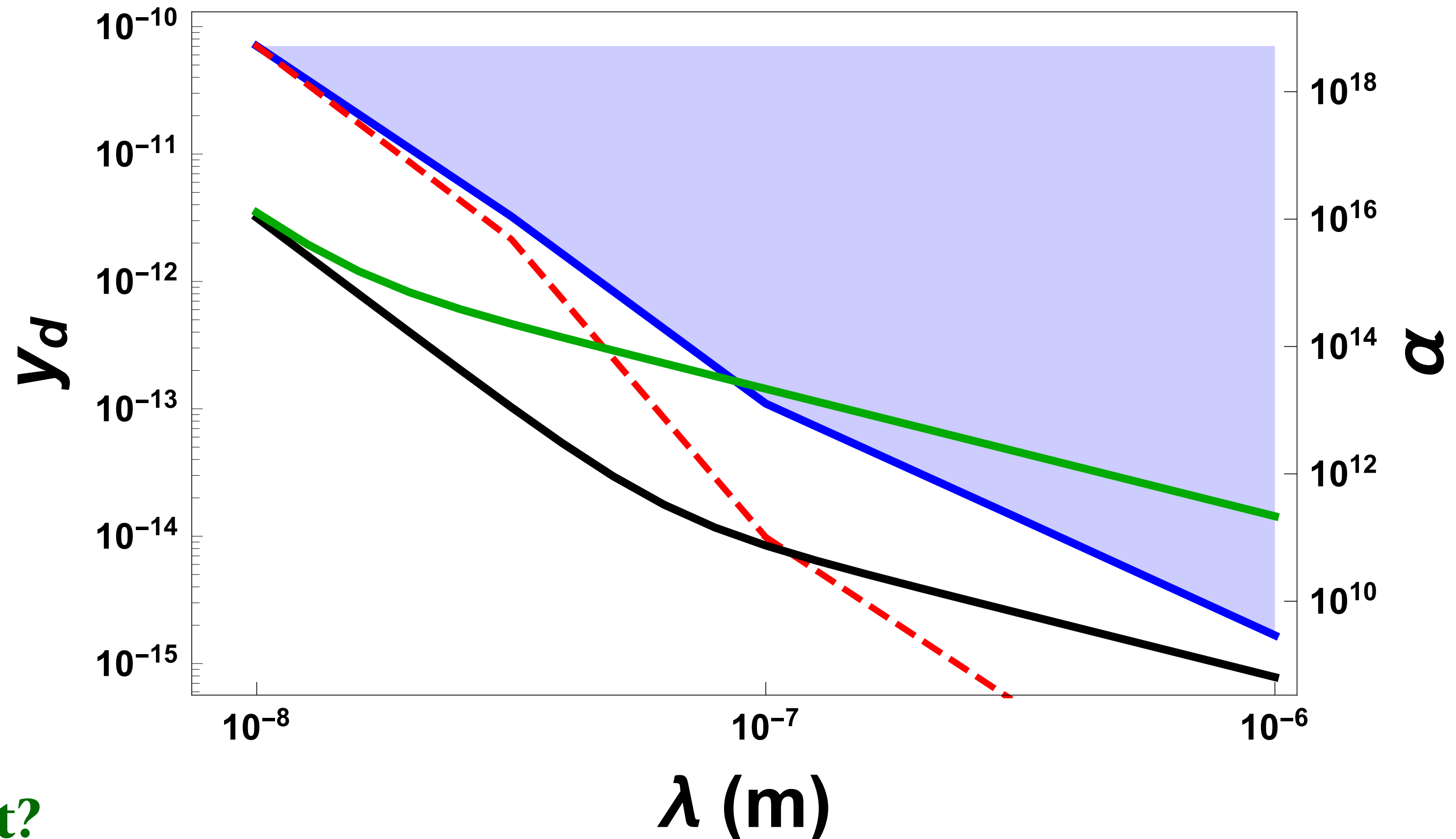
$$^{57}\text{Fe} \quad \Delta E = 10^{-15} \text{ eV}$$

$$^{181}\text{Tl} \quad \Delta E = 10^{-17} \text{ eV}$$

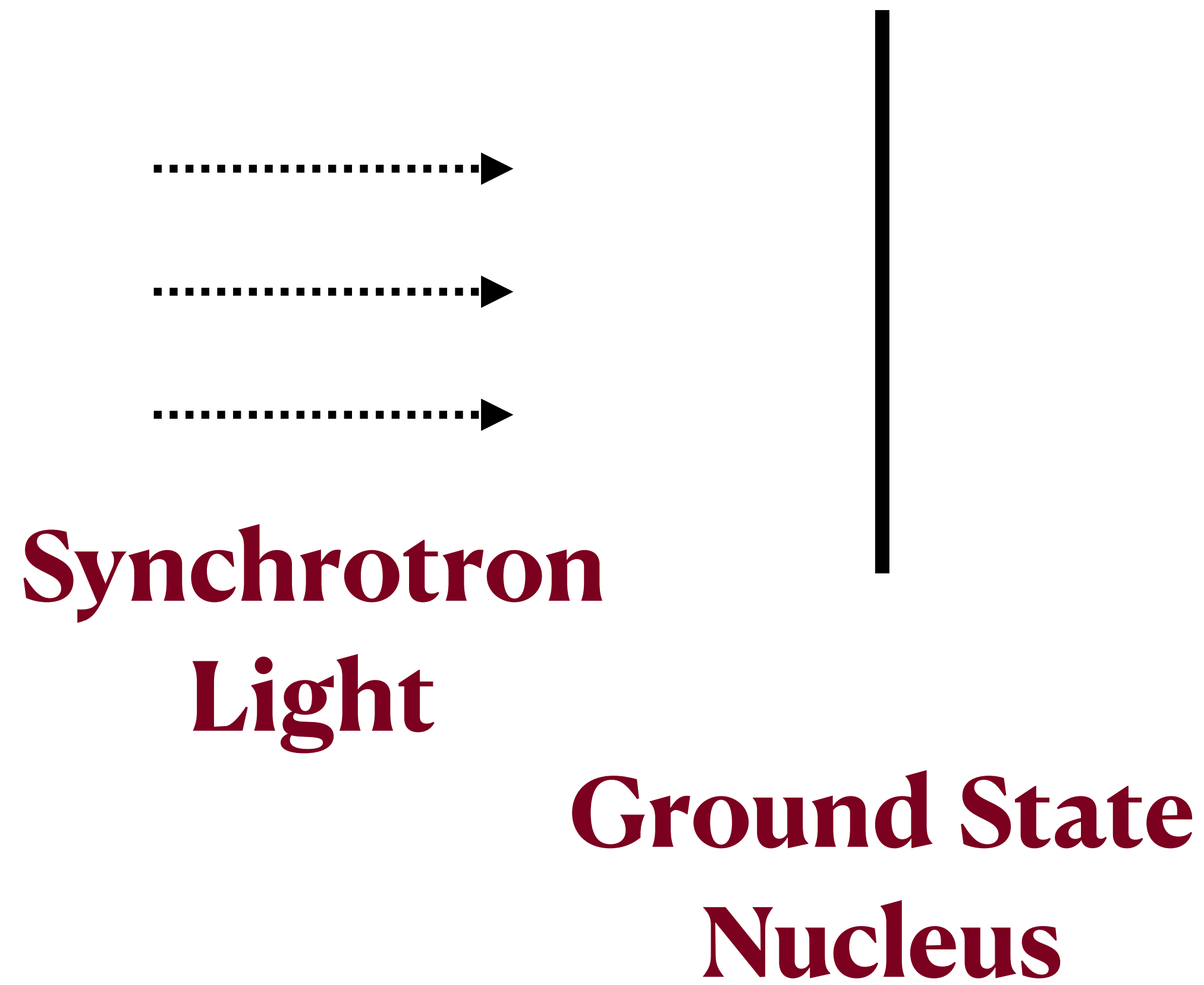
$$N_\gamma = 3 \times 10^{14}$$

Second Order Casimir Background at
shortest distances

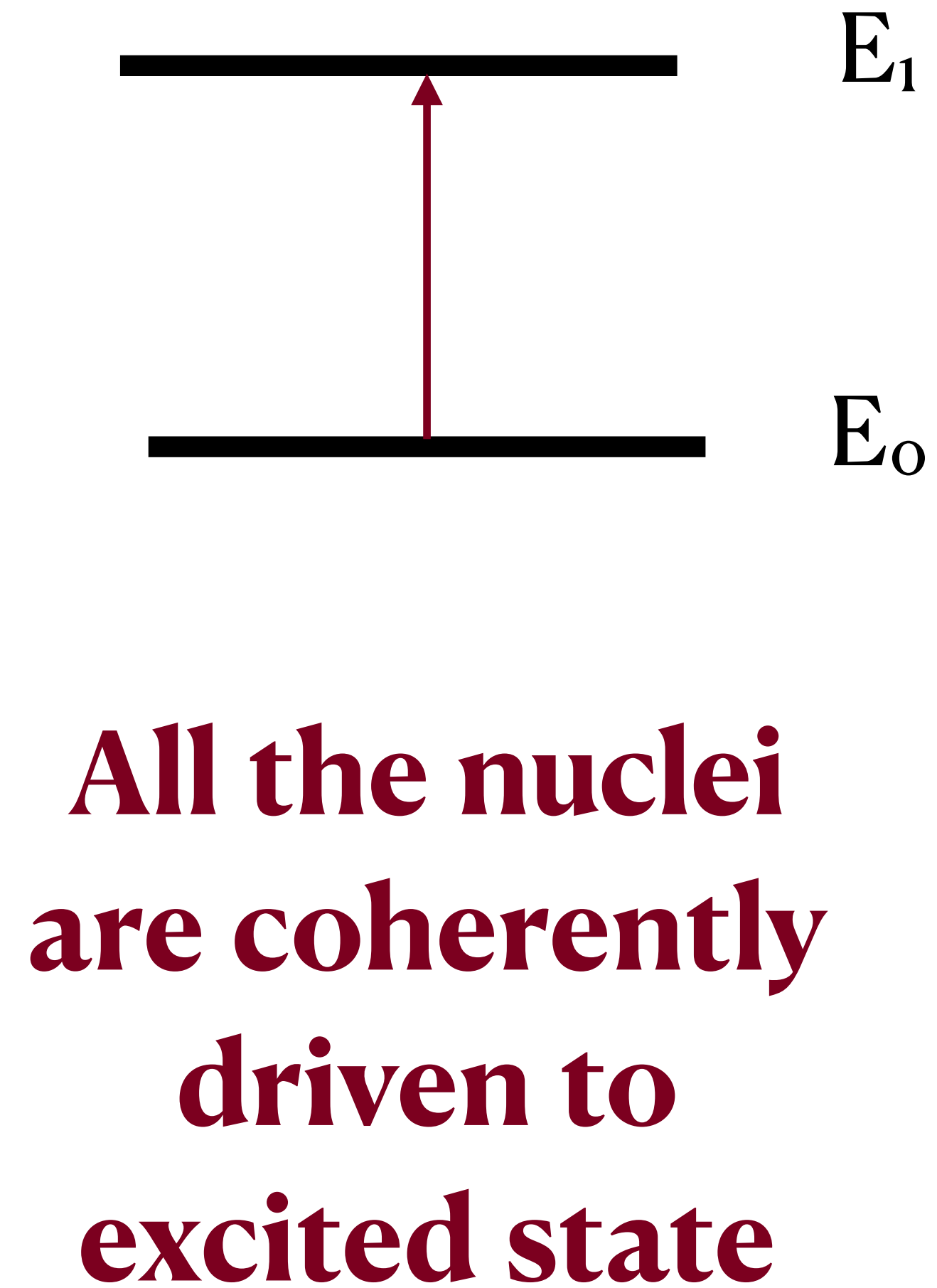
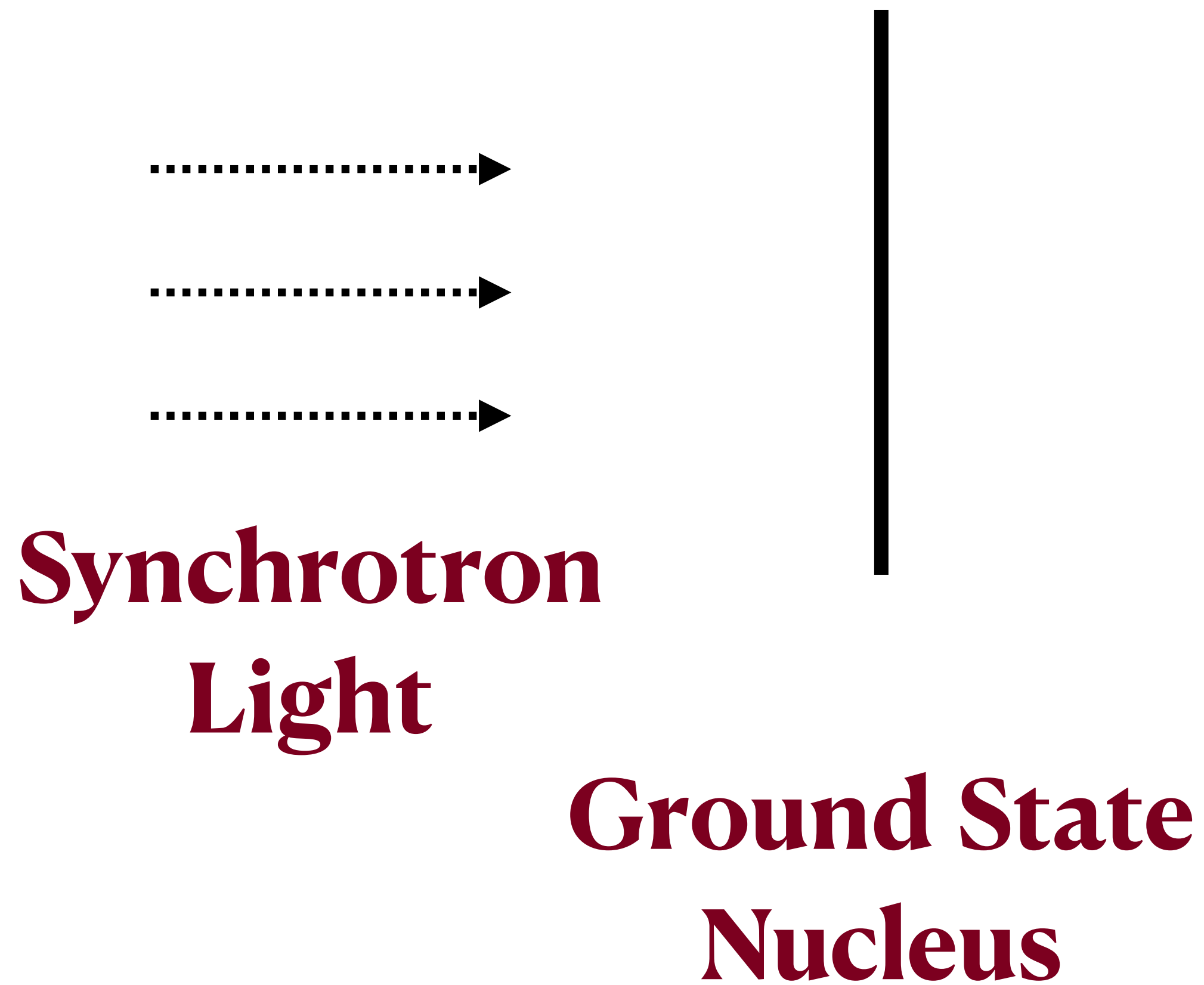
Mitigate using differential measurement?



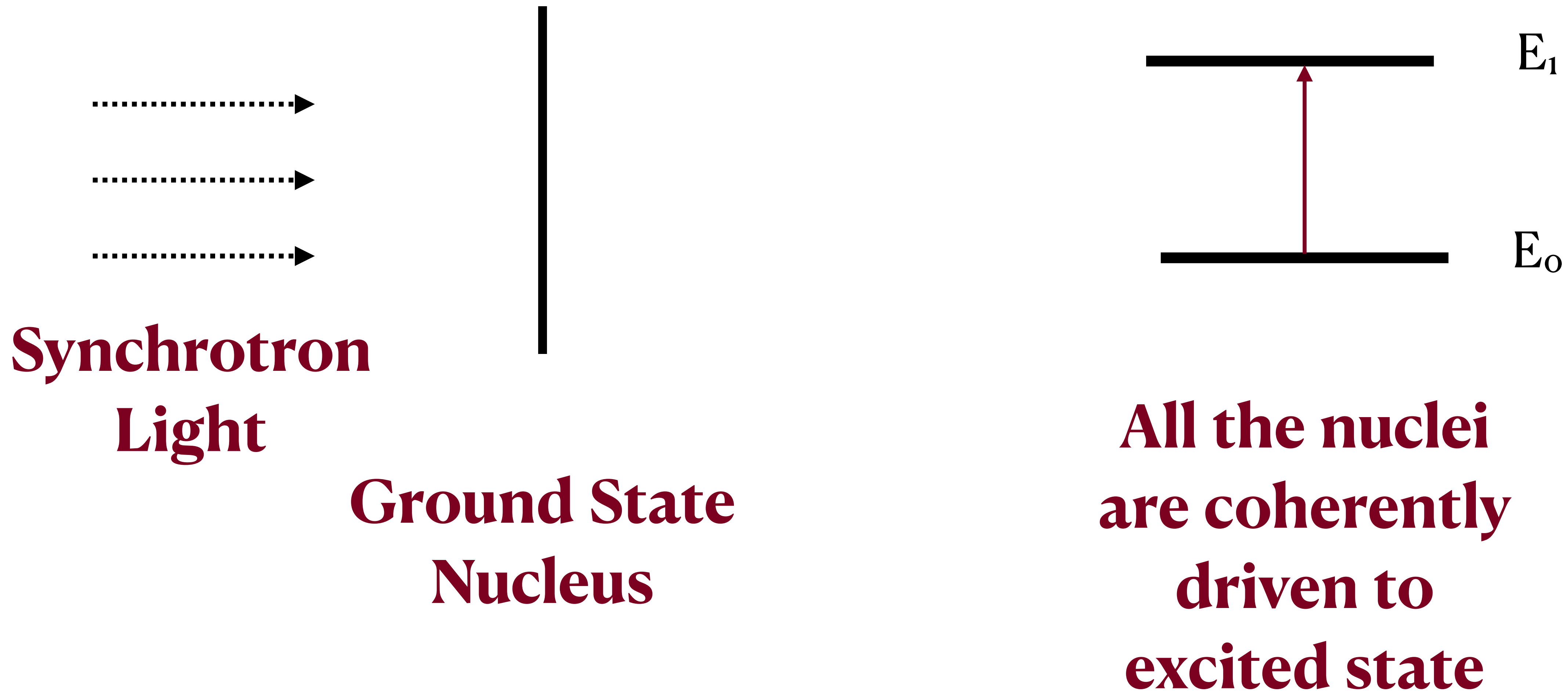
Mossbauer at Synchrotron Source



Mossbauer at Synchrotron Source



Mossbauer at Synchrotron Source



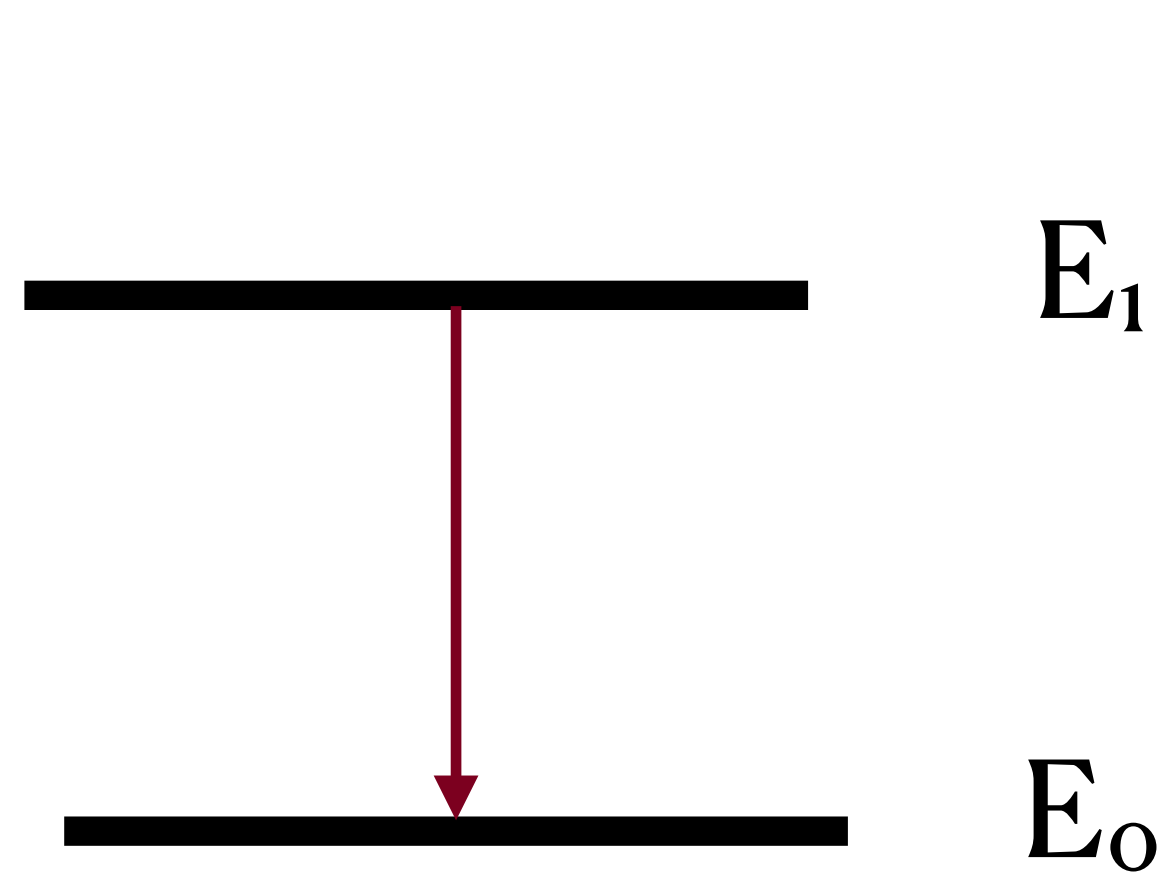
Short Pulse \ll Lifetime of State

Mossbauer at Synchrotron Source

Well after pulse, state starts to decay

Coherent initial excitation \Rightarrow decays in phase

Decay along forward direction amplified by in phase addition



In phase decay

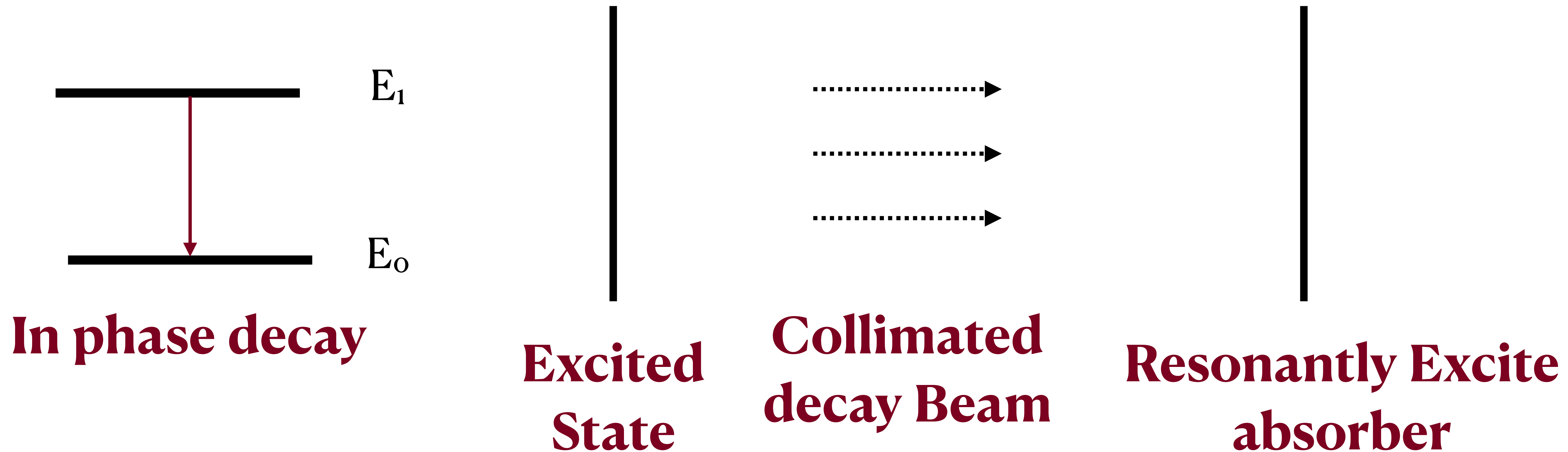
**Excited
State**

Mossbauer at Synchrotron Source

Well after pulse, state starts to decay

Coherent initial excitation => decays in phase

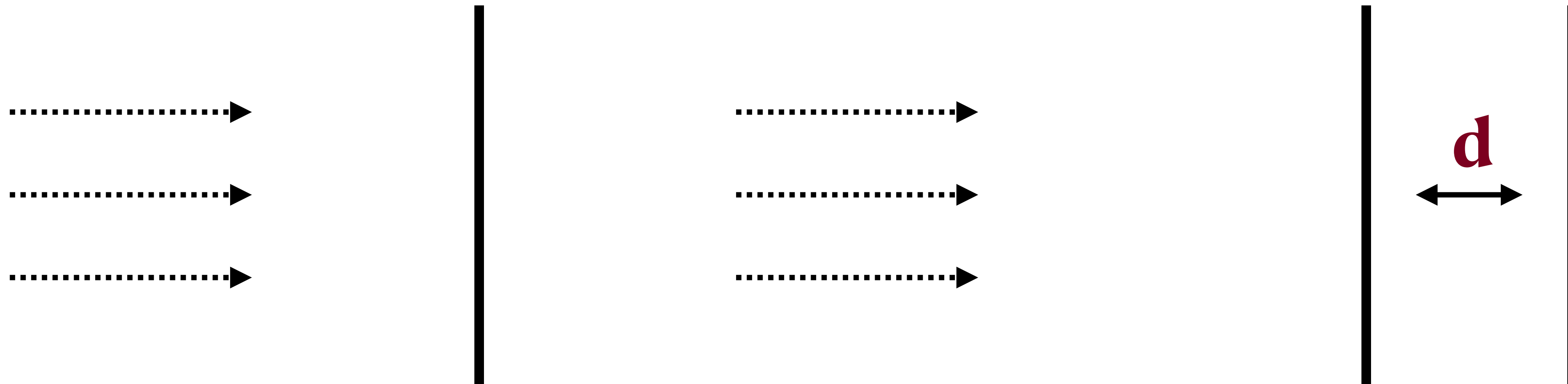
Decay along forward direction amplified by in phase addition



Mossbauer at Synchrotron Source



Mossbauer at Synchrotron Source

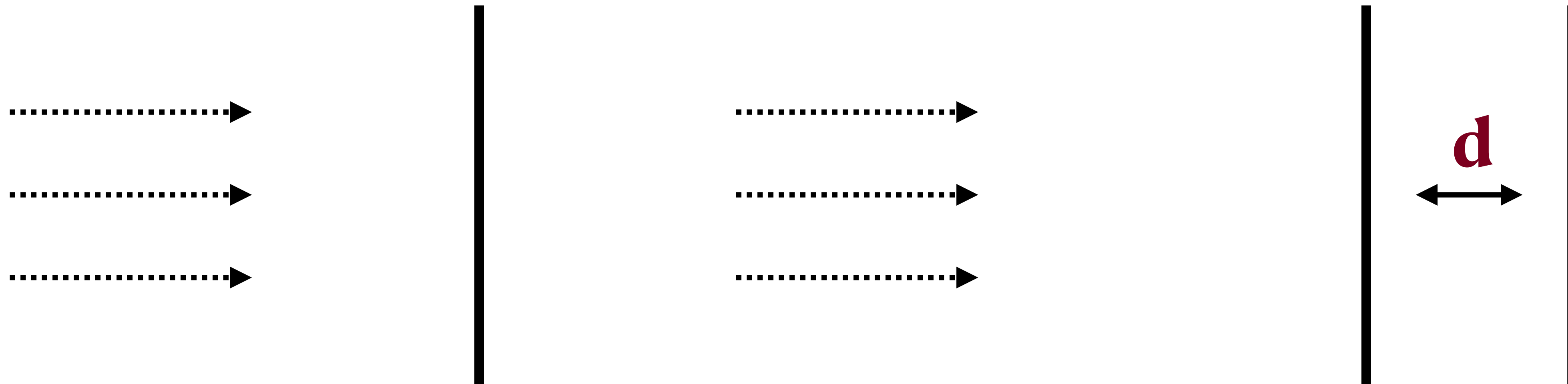


Send Synchrotron Pulse

Well after pulse, collimated emission

Measure resonant reabsorption as a function of d

Mossbauer at Synchrotron Source



Send Synchrotron Pulse

Well after pulse, collimated emission

Measure resonant reabsorption as a function of d

Why?

Clean excitation unlike radioactive decay

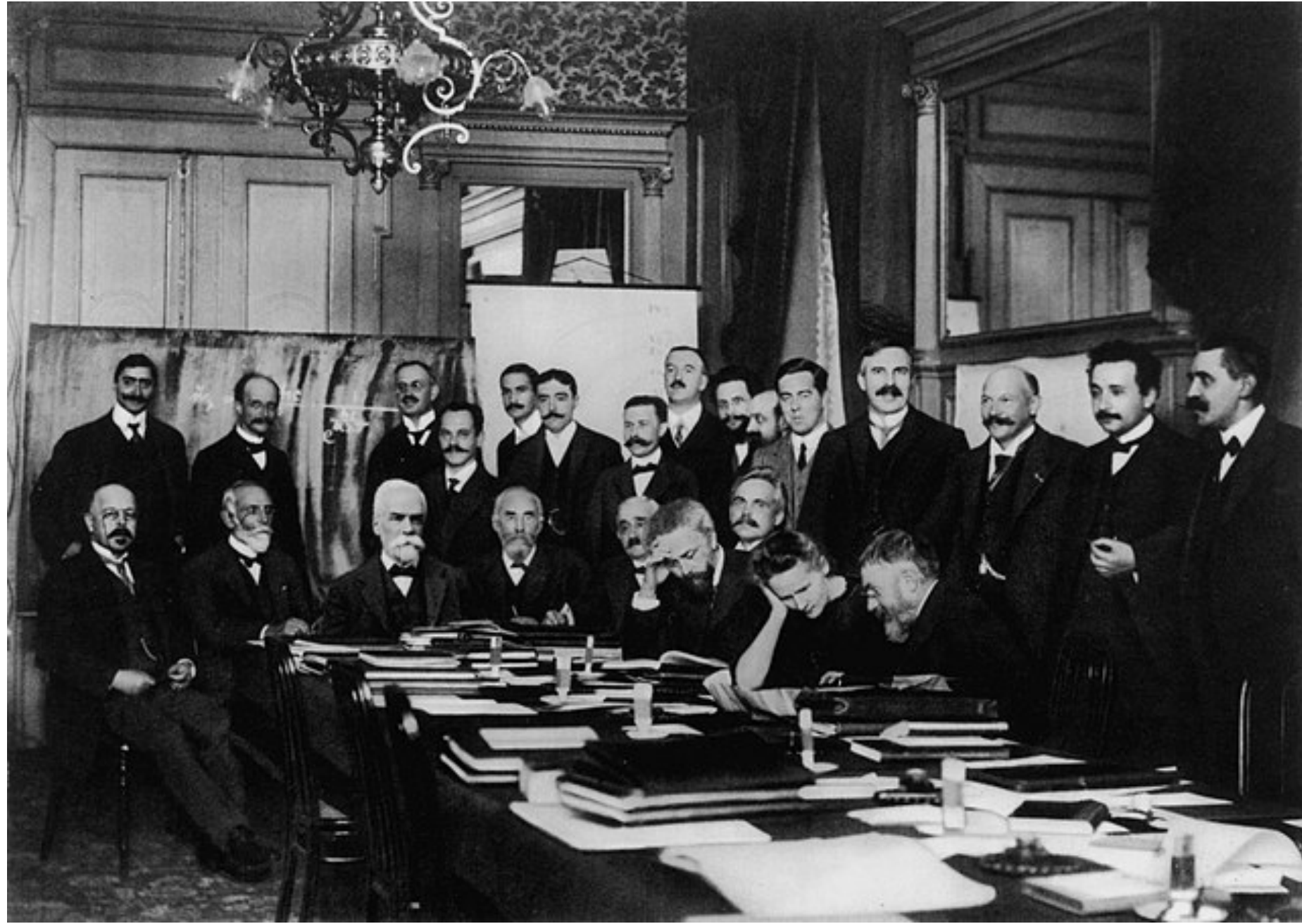
May enable new class of ultra narrow Mossbauer

Conclusions

- 1. Mossbauer Effect seems well suited to probe short distance forces**
- 2. Natural electromagnetic background suppression**
- 3. Ideal for scalar and tensor forces**
- 4. Synchrotron light sources may enable new Mossbauer sources**

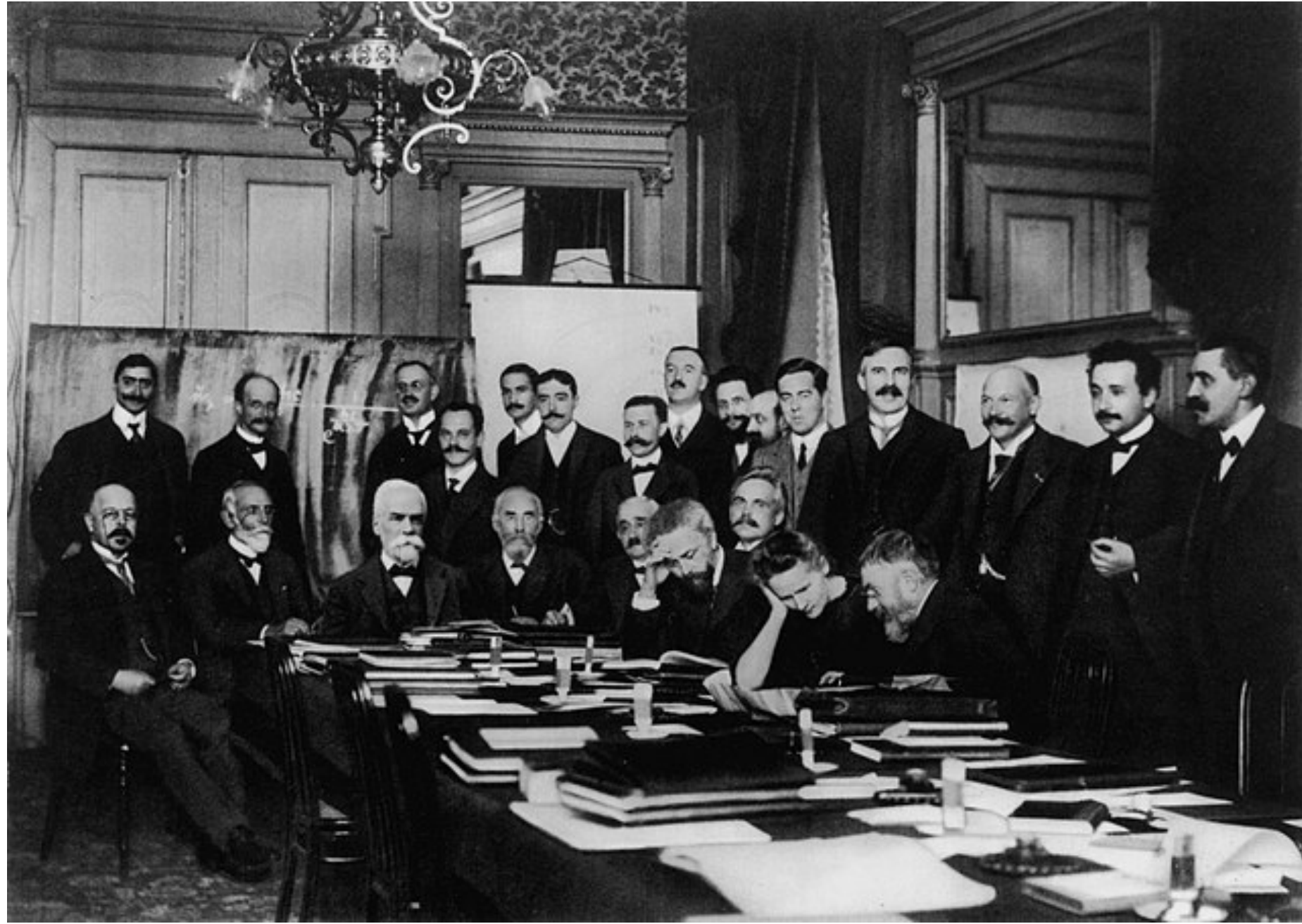
Non Linear Quantum Mechanics

Non Linear Quantum Mechanics?



**Theory built on observations in the 1900s
Why should it be “the absolute truth”?**

Non Linear Quantum Mechanics?



Theory built on observations in the 1900s
Why should it be “the absolute truth”?

What?

Two Postulates of Quantum Mechanics

Probability

Linearity

Which?

Probability

Probability

Finite system has a finite set of energies

Continuous observables and symmetries

Probability

Finite system has a finite set of energies
Continuous observables and symmetries } **Deterministic
Observables?**

Probability

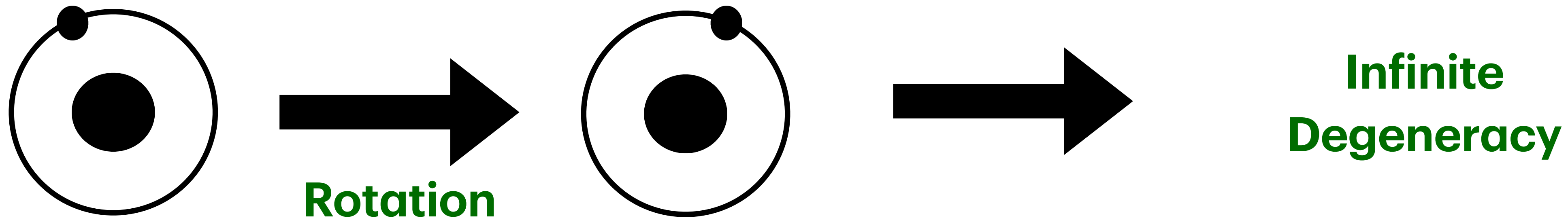
Finite system has a finite set of energies
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Could an electron in an atom have a well defined position?

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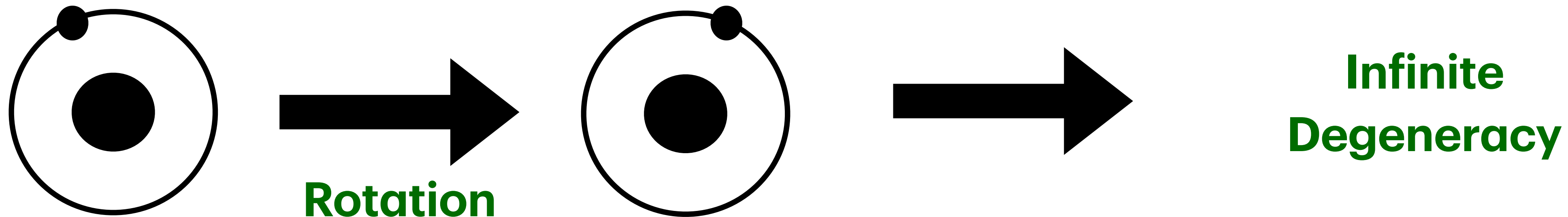
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Probability

Finite system has a finite set of energies
Continuous observables and symmetries } Deterministic Observables?

Could an electron in an atom have a well defined position?



Quantum Mechanics

Sacrifice Determinism.

Preserve finite set of energy states, continuous symmetries and observables

Bell Inequalities, Kochen-Specker, SSC Theorems

Causality and Entanglement

Trial Non-Linear Term

$$i \frac{\partial \Psi}{\partial t} = H_L \Psi + \epsilon \left(\Psi^2 + \Psi^{*2} \right) \Psi$$

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Entanglement is fundamental to quantum mechanics

$$\Psi (x, y; t) = \sum_{i,j} c_{ij} (t) \alpha_i (x) \beta_j (y)$$

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$$\Psi(x, y; t) = \sum_{i,j} c_{ij}(t) \alpha_i(x) \beta_j(y)$$

Apply some local operation on x: $\alpha_i(x) \rightarrow U \alpha_i(x)$

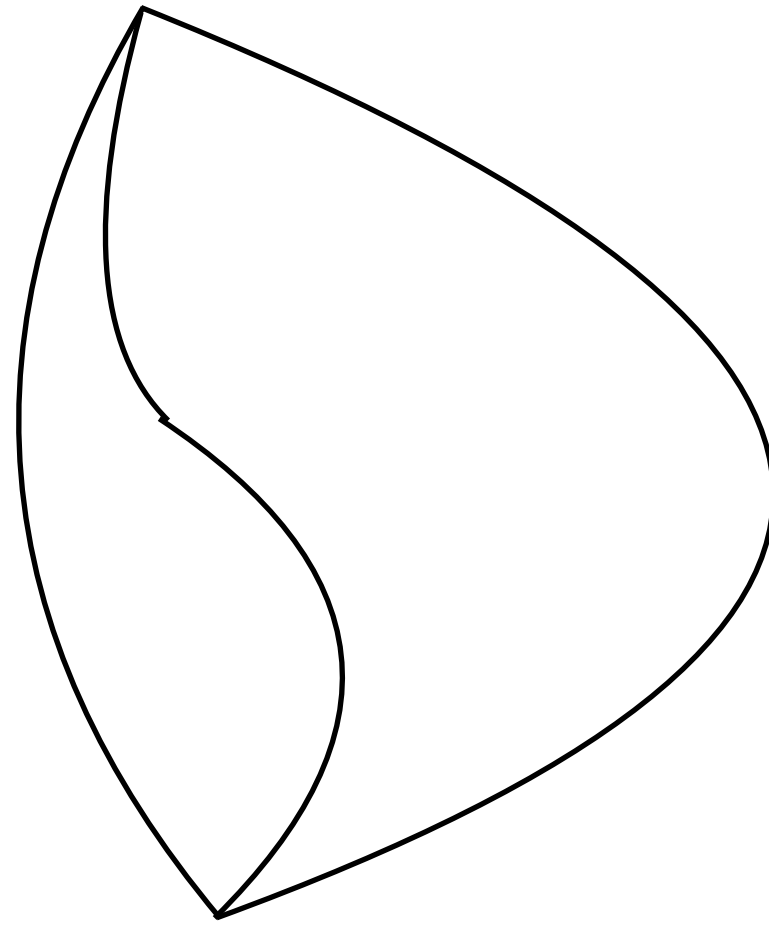
Does it instantly change the time evolution of y?

YES

Not causal

Linearity

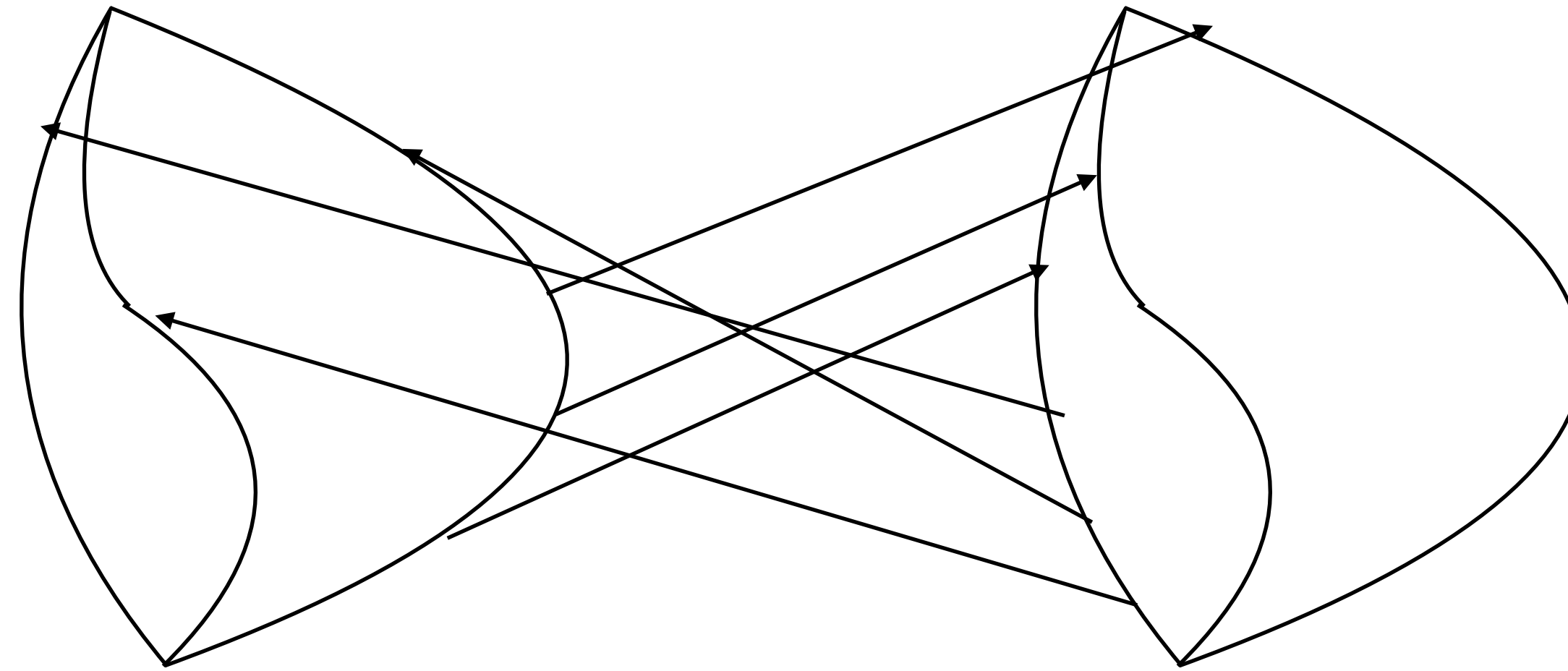
Electron Coupled to Electromagnetism



**Electron paths do not
interact via
electromagnetism**

Linearity

Electron Coupled to Electromagnetism

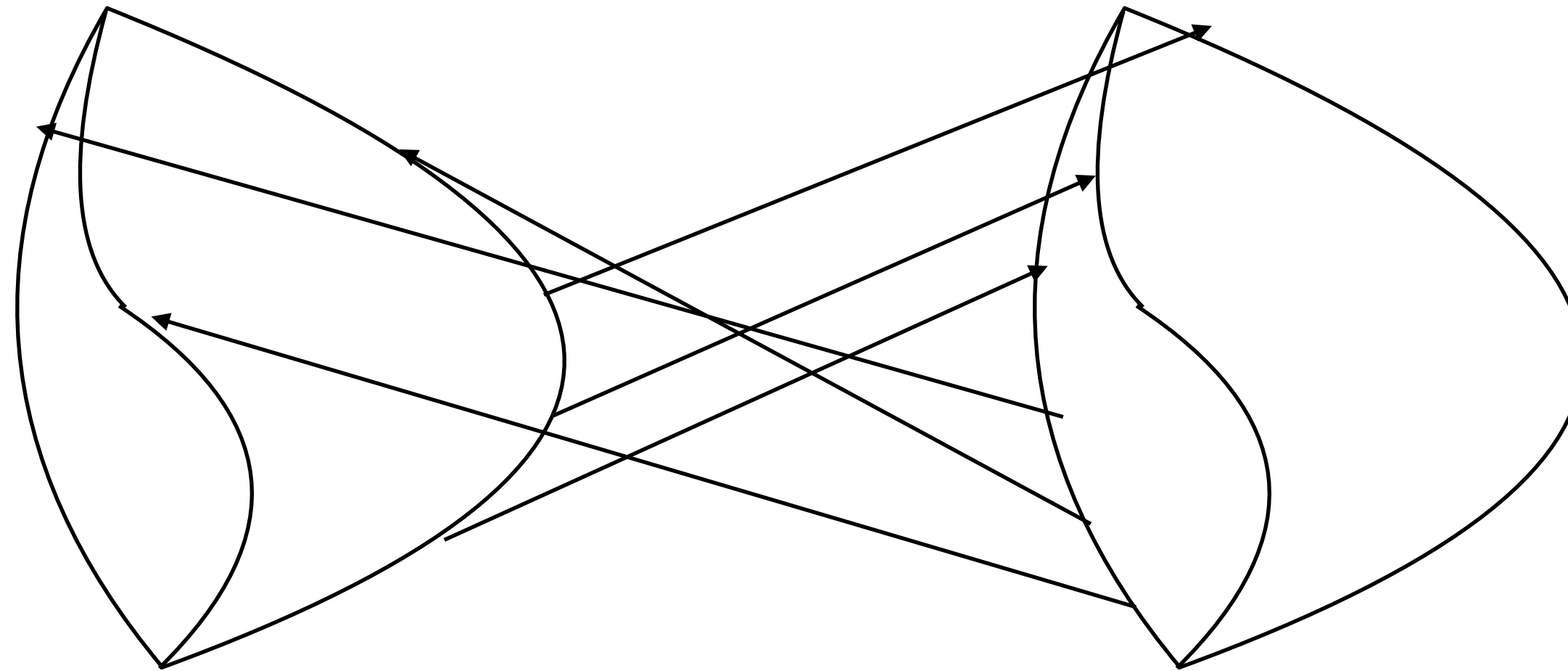


**Electron paths do not
interact via
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**Paths of two electrons
interact causally (QFT)**

Linearity

Electron Coupled to Electromagnetism



**Electron paths do not
interact via
electromagnetism**

**Paths of two electrons
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Why can't path talk to itself?

**Natural Language:
Quantum Field Theory**

The Framework

The Schrodinger Picture of Quantum Field Theory

$|\chi(t)\rangle$

Quantum State of Fields
(e.g. in Fock states)

$\phi(x)$

Time Independent
Operators

$$H = \int d^3x \mathcal{H}(\phi(x), \pi(x))$$

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Action

$$S = \int dt (i \langle \chi | \dot{\chi} \rangle - \langle \chi | H | \chi \rangle)$$

The Framework

Yukawa $H \supset \int d^3x \, y \, \phi(x) \, \bar{\Psi}(x) \, \Psi(x)$

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Non-linearities in the operators but not in the state

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Linear QFT: $S \supset \left(\int d^3x \, y \, \langle \chi(t) | \phi(x) \, \bar{\Psi}(x) \, \Psi(x) | \chi(t) \rangle \right)$

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Higher order in states - leads to state dependent quantum evolution

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Obeys all the rules

Higher order in states - leads to state dependent quantum evolution

Analyze non-linearity perturbatively

Perturbation Theory

$$\mathcal{H} \supset y\Phi\bar{\Psi}\Psi = (y\phi + \epsilon\langle\chi|\phi|\chi\rangle)\bar{\Psi}\Psi$$

$$i\frac{\partial|\chi\rangle}{\partial t} = H|\chi\rangle$$

At zeroth order, this is just standard QFT

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Causality: Non-linearity enters via expectation value. At lowest order, causal from QFT.

Causal background field for all higher orders

Single Particle

$$\mathcal{L} \supset y \Phi \bar{\Psi} \Psi = y (\phi + \tilde{\epsilon} \langle \chi | \phi | \chi \rangle) \bar{\Psi} \Psi$$

Suppose we have a ψ particle - how does its wave-function evolve?

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Straightforward Computation of Expectation Value

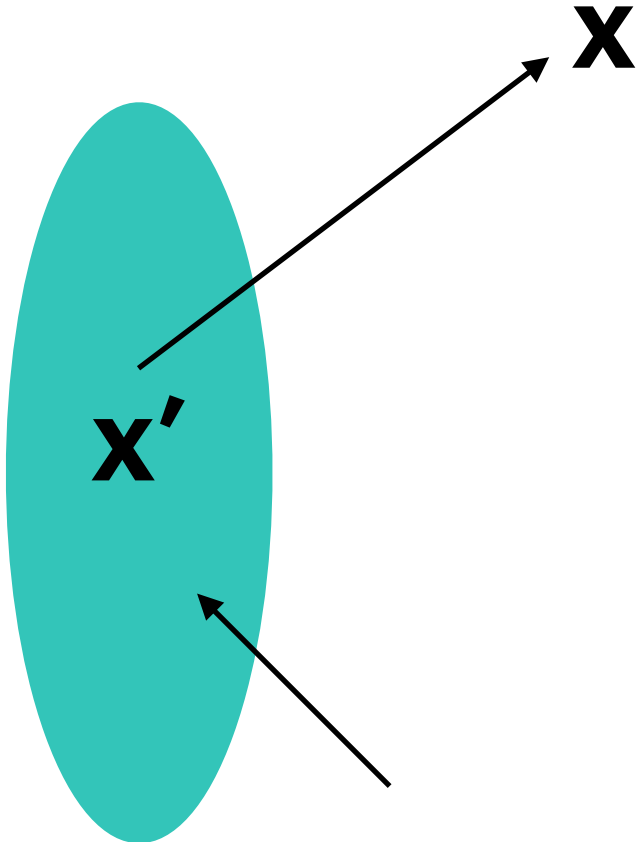
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Charge Density of ψ

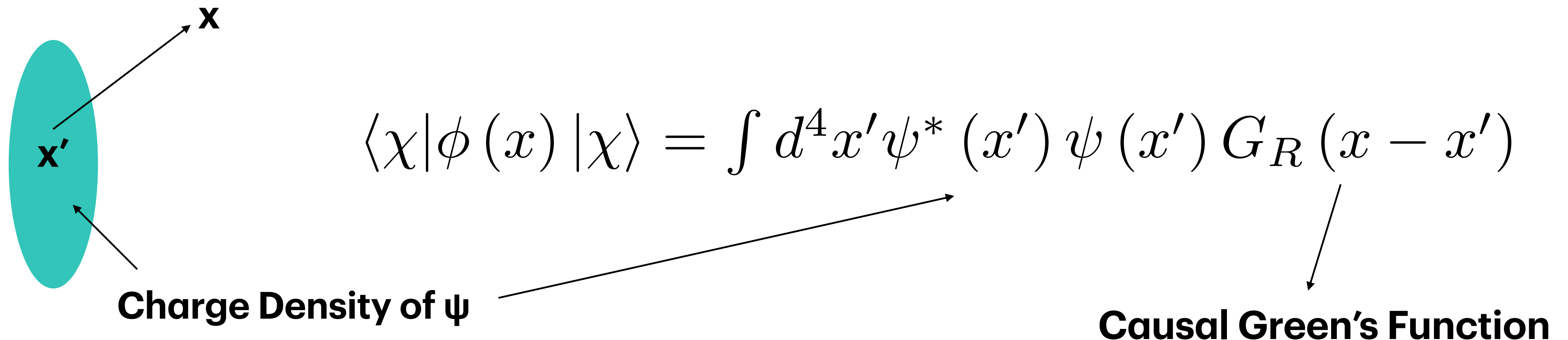
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Causal Green's Function

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Single particle equation derived from field theory

Equation depends upon theory (Yukawa, Φ^4 etc)

$$i\frac{\partial\Psi(t,\mathbf{x})}{\partial t} = \left(H + \tilde{\epsilon}y \int d^4x' \Psi^*(x) \Psi(x') G_R(x;x')\right) \Psi(t,\mathbf{x})$$

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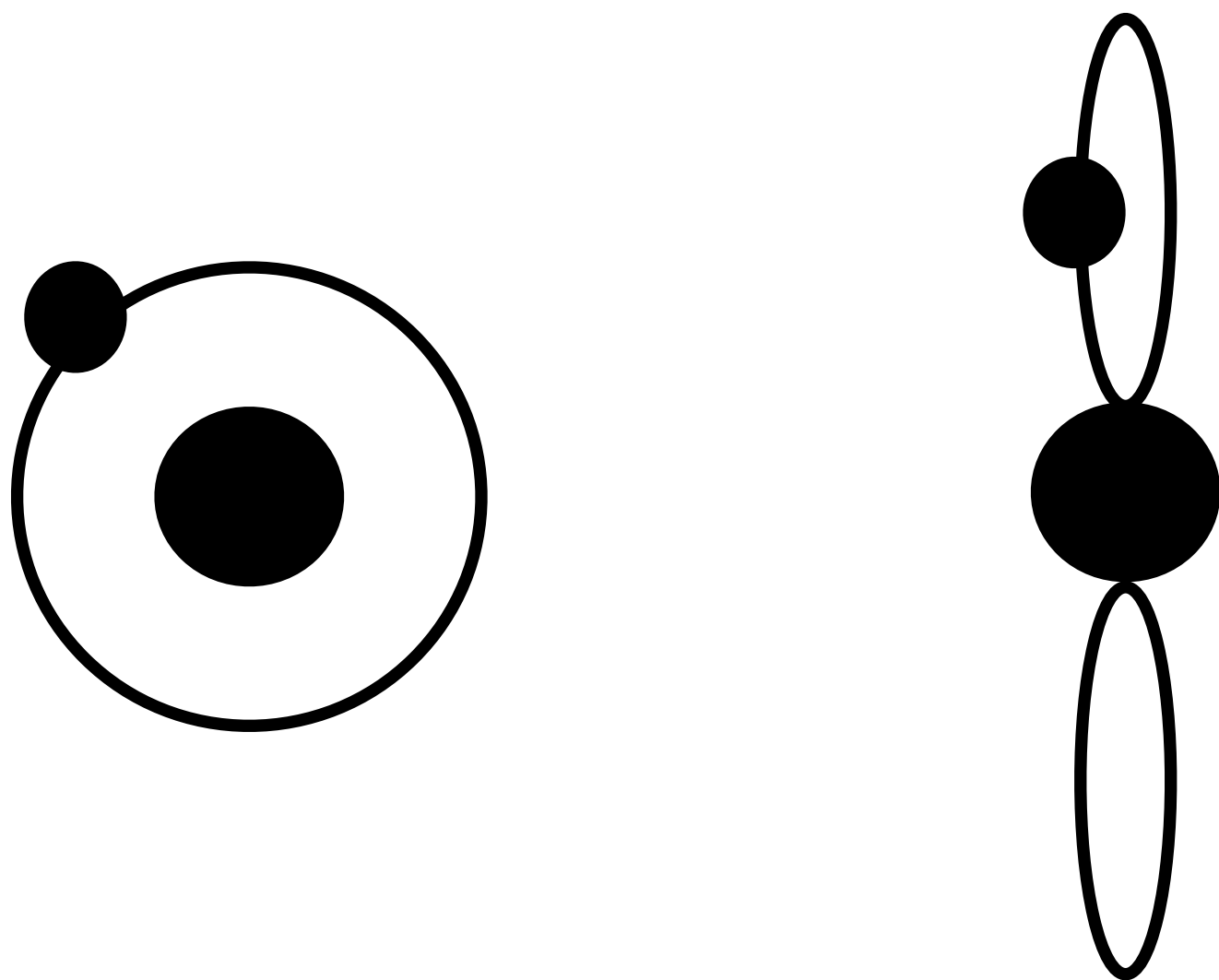
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Fixed Central particle

Self interaction of wave-function breaks degeneracy of levels



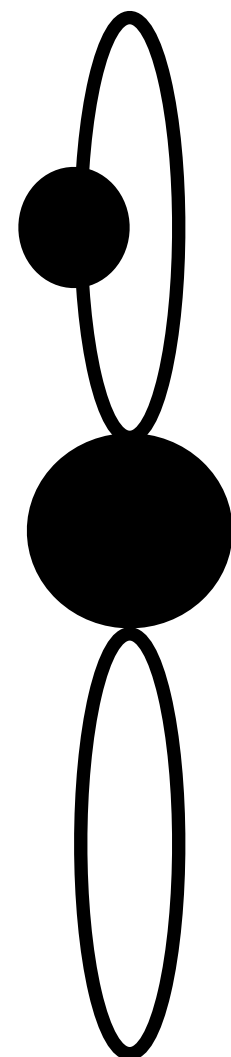
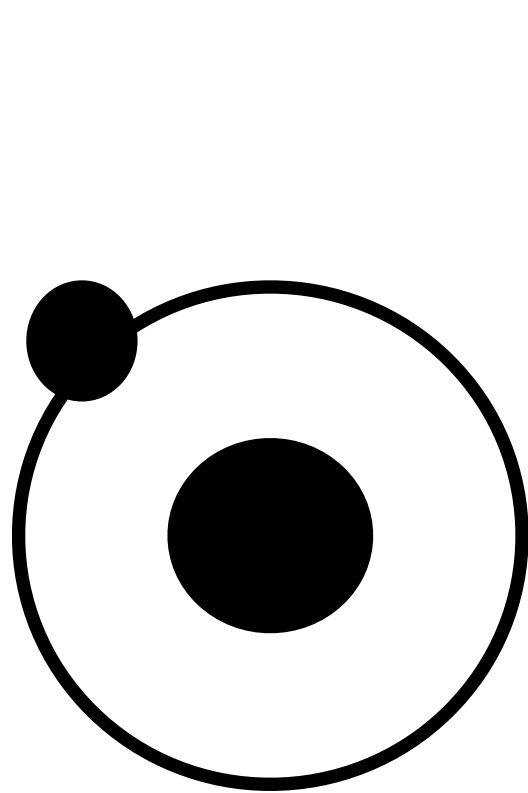
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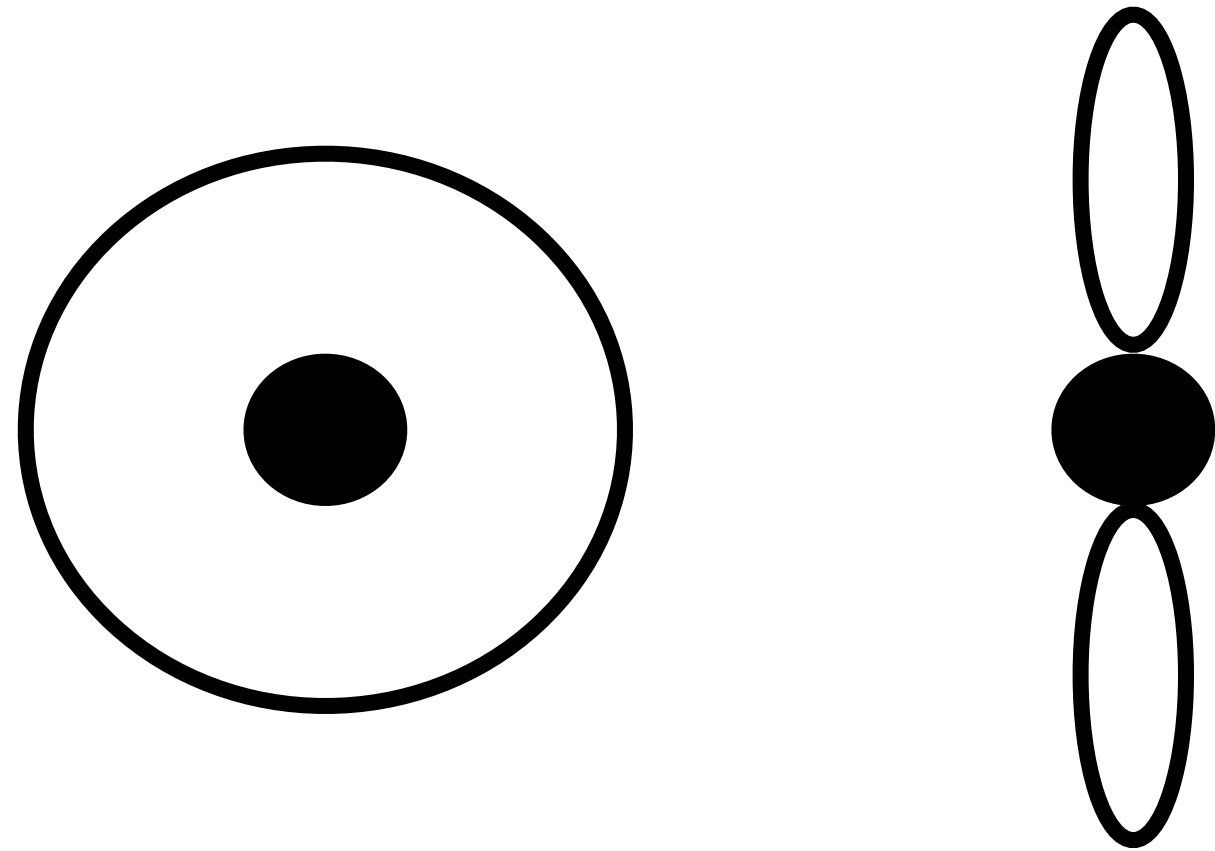
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**Hermitean Form of Hamiltonian implies conserved
norm**

Constraints

What does this do to the Lamb Shift?



Proton at Fixed Location

2S and 2P electron have different charge distribution

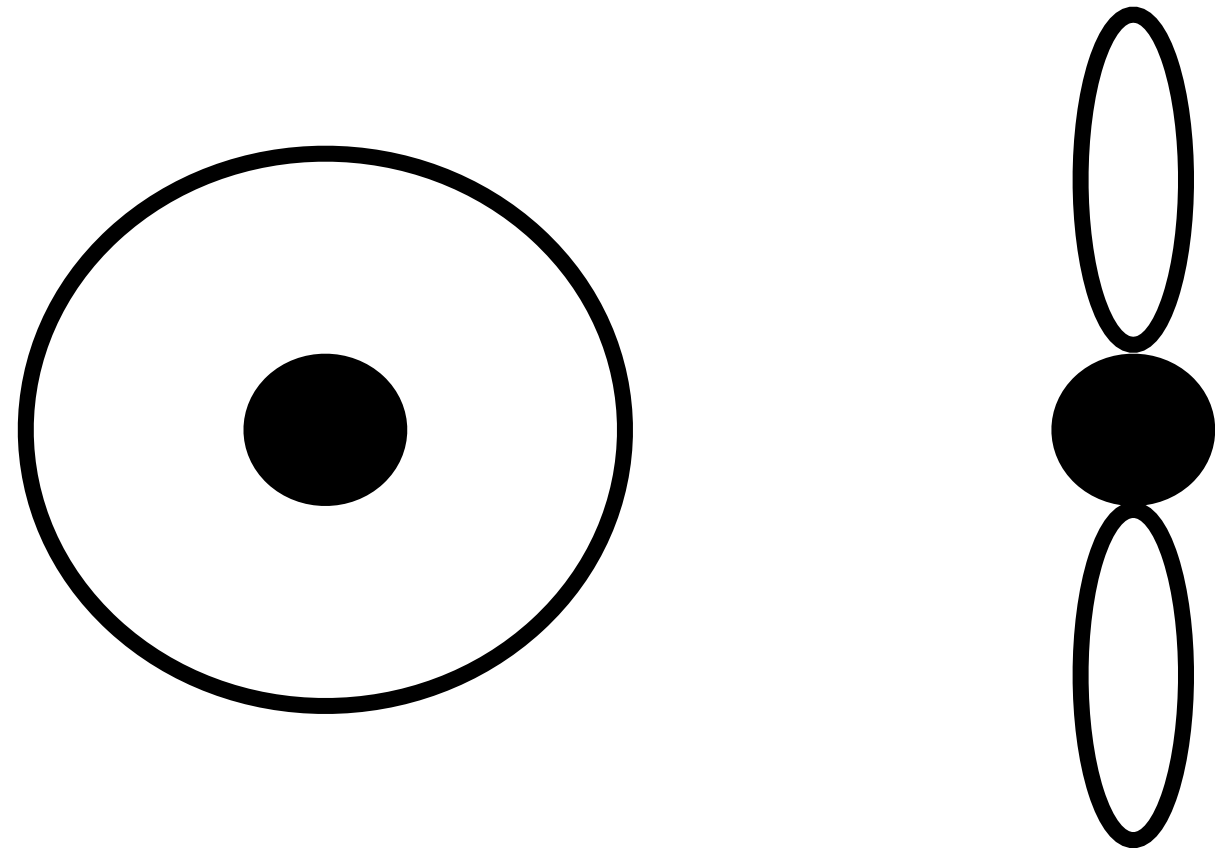
Different expectation value of electromagnetic field

Level Splitting!

$$\langle \chi | A_\mu | \chi \rangle J^\mu$$

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Level Splitting!

$$\langle \chi | A_\mu | \chi \rangle J^\mu$$

BUT: Cannot decouple center of mass and relative co-ordinates

Proton wave-function spread over some region (e.g. trap size ~ 100 nm)

Expectation value of electromagnetic field diluted

In neutral atom - heavily suppressed, except at edges!

$$\varepsilon < 10^{-2}$$

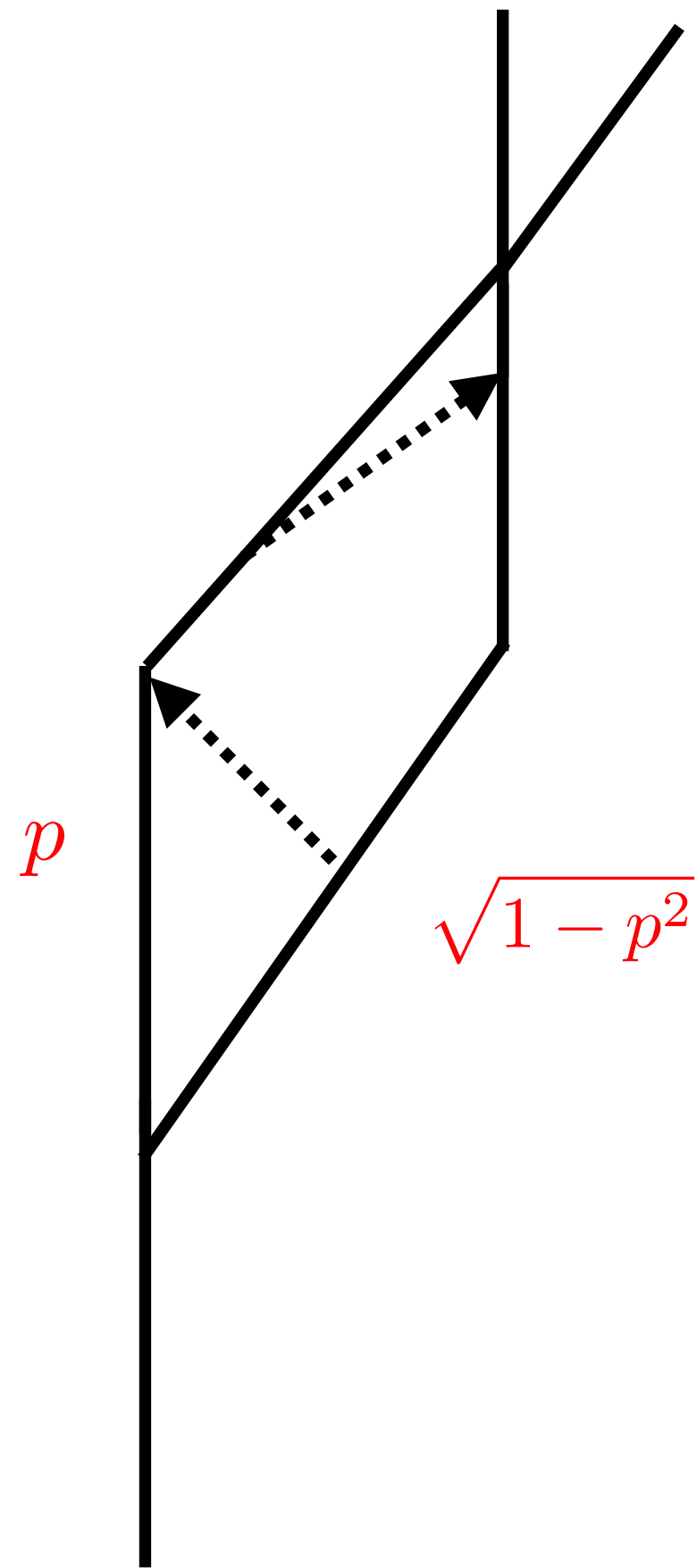
Similarly, kills possible bounds on QCD and gravity

Experimental Tests

Interferometry - interaction between paths

Take an ion - split its wave-function

Coulomb Field of one path interacts with the other path



Experimental Tests

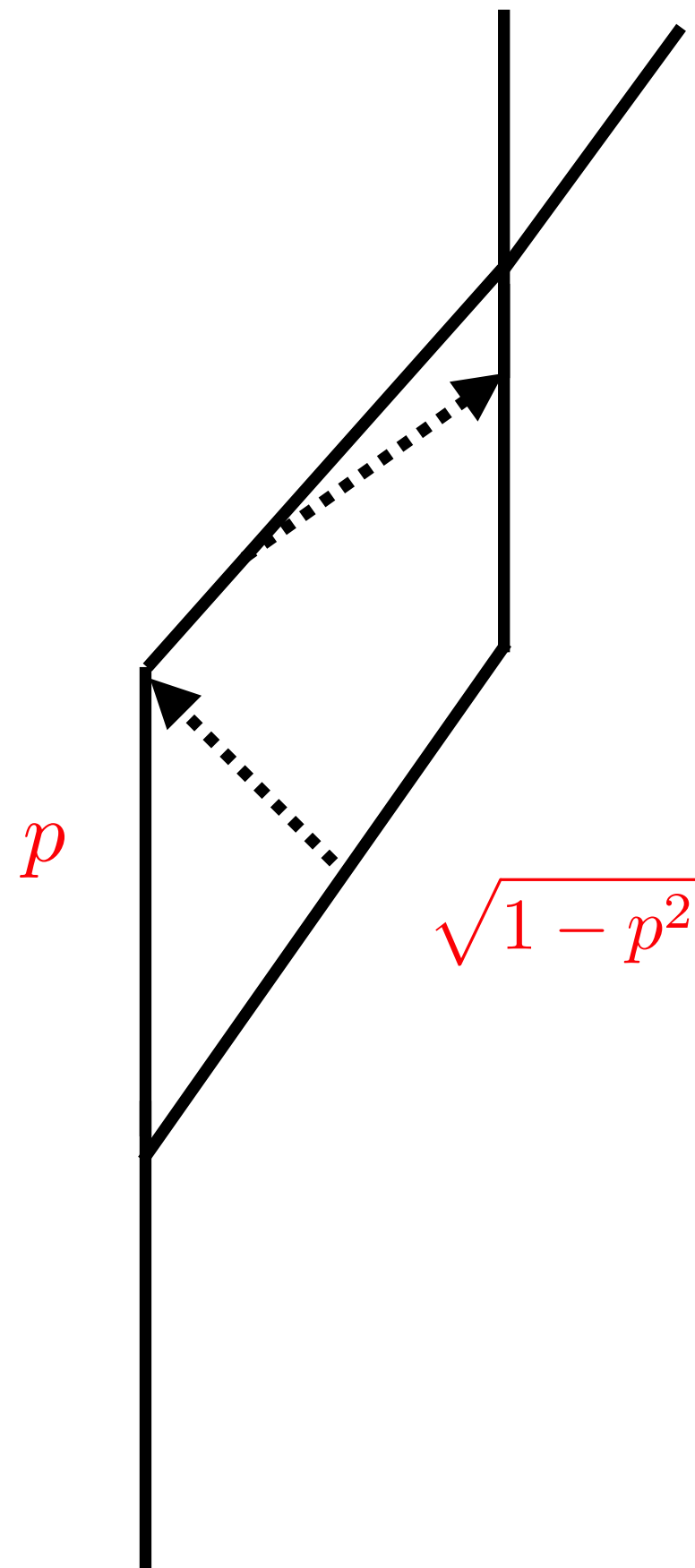
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Gives rise to phase shift that depends on the intensity p of the split

Use intensity dependence to combat systematics



Conclusions

- 1. Quantum Field Theory can be generalized to include non-linear, state dependent time evolution**
- 2. Conventional tests of quantum mechanics in atomic and nuclear systems do NOT probe causal non-linear quantum mechanics**
- 3. Straightforward set of experimental tests possible to probe non-linear quantum mechanics**
- 4. Motivation to test other extensions as well - e.g. Lindblad Decoherence**