Search for physics beyond the standard model with Atomic Clocks

Optical vs. microwave clocks



Applications of atomic clocks







GPS

Very Long Baseline Interferometry

Relativistic geodesy



Definition of the second



Quantum simulation



Search for physics beyond the Standard Model

Image Credits: NOAA, Science 281,1825; 346, 1467, University of Hannover, PTB

Search for physics beyond the standard model with atomic clocks

Atomic clocks can measure and compare frequencies to exceptional precisions!

If fundamental constants change (now) due to for various "new physics" effects atomic clock may be able to detect it.



Fermions: spin = 1/2 particles

Standard Model



According to the Standard Model



Our Universe can not exist !

We don't know what most of the Universe is!



ARE FUNDAMENTAL CONSTANTS CONSTANT???

Being able to compare and reproduce experiments is at the foundation of the scientific approach, which makes sense only if the laws of nature do not depend on time and space.

J.-P. Uzan, Rev. Mod. Phys. 75, 403 (2003)

FUNDAMENTAL CONSTANTS

Quantity	Symbol	Numerical Value
Speed of light (in vacuum)	С	$3.00 \times 10^8 \mathrm{m \ s^{-1}}$
Gravitational constant	G	$6.67 \times 10^{-11} \text{ N m}^2 \text{ kg}^{-2}$
Avogadro's number	N _A	6.02×10^{23} molecules mole ⁻¹
Universal gas constant	R	8.31 J K ⁻¹ mole ⁻¹
Boltzmann constant	k _B	$1.38 imes 10^{-23} { m J K^{-1}}$
		$8.62 \times 10^{-5} \mathrm{eV} \mathrm{K}^{-1}$
Stefan's constant	σ	$5.67 \times 10^{-8} \text{ W m}^{-2} \text{ K}^{-4}$
Atomic mass unit	u	1.66×10^{-27} kilograms
Coulomb constant	k	$9.00 \times 10^9 \text{ N m}^2 \text{ C}^{-2}$
	$\epsilon_0 = 1/4\pi k$	$8.85 \times 10^{-12} \mathrm{C}^2 \mathrm{N}^{-1} \mathrm{m}^2$
Biot-Savart constant	k'	10 ⁻⁷ T m A ⁻¹
Electron charge	-e	-1.60×10^{-19} coulombs
Electron mass	m _e	9.11×10^{-31} kilograms
Proton charge	e	1.60×10^{-19} coulombs
Proton mass	m_p	1.673×10^{-27} kilograms
Neutron mass	m _n	1.675×10^{-27} kilograms
Planck's constant	h	$6.63 \times 10^{-34} \text{ J s}$
		$4.14 \times 10^{-15} \mathrm{eV} \mathrm{s}$
	$\hbar = h/2\pi$	$1.055 \times 10^{-34} \text{ J s}$
		$6.58 \times 10^{-16} \mathrm{eV}\mathrm{s}$
Rydberg constant	R _H	$1.10 \times 10^7 \text{ metres}^{-1}$
Bohr radius	a ₀	5.29×10^{-11} metres
Bohr magneton	μ_B	$9.27 \times 10^{-24} \mathrm{J} \mathrm{T}^{-1}$

Note: we are still measuring them ...

2006 CODATA RECOMMENDED VALUES OF THE FUNDAMENTAL CONSTANTS OF PHYSICS AND CHEMISTRY NIST SP 959 (Aug/2008) /

Values from: P. J. Mohr, B. N. Taylor, and D. B. Newell, *Rev. Mod. Phys.* **80**, 633 (2008) and *J. Phys. Chem. Ref. Data* **37**, 1187 (2008). The number in parentheses is the one-sigma (1σ) uncertainty in the last two digits of the given value.

Quantity	Symbol	Numerical value	Unit
speed of light in vacuum	c, c_0	299792458 (exact)	$m s^{-1}$
magnetic constant	μ_0	$4\pi \times 10^{-7}$ (exact)	$N A^{-2}$
electric constant $1/\mu_0 c^2$	ϵ_0	$8.854187817 imes 10^{-12}$	$\rm F~m^{-1}$
Newtonian constant of gravitatio	n G	$6.67428(67) \times 10^{-11}$	$m^3 kg^{-1} s^{-2}$ 324/80
Planck constant	h	$6.62606896(33) \times 10^{-34}$	Js 957(29)
$h/2\pi$	\hbar	$1.054571\frac{628(53)}{2} \times 10^{-34}$	Js 726/47
elementary charge	e	$1.602176\frac{487(40)}{487(40)} \times 10^{-19}$	C 565/35
fine-structure constant $e^2/4\pi\epsilon_0\hbar c$	α	$7.2973525376(50) \times 10^{-3}$	18(215,00)
inverse fine-structure constant	α^{-1}	137.035999 679(94) 274 (4	T)
Rydberg constant $\alpha^2 m_{\rm e} c/2h$	R_{∞}	10973731.568527(73)	m^{-1} 39(33)
Bohr radius $\alpha/4\pi R_{\infty}$	a_0	$0.52917720859(36) \times 10^{-10}$	m 1092(17)
Bohr magneton $e\hbar/2m_{\rm e}$	$\mu_{ m B}$	$927.400915(23) \times 10^{-26}$	JT ⁻¹ 68(20)

From NIST Tech Beat: July 19, 2011

Which fundamental constants to consider?

A pragmatic approach: choose a theoretical framework so that the set of undetermined fixed parameters is fully known. Then, try to determine if these values are constant.

It only makes sense to **consider the variation of dimensionless ratios:**

Fine-structure constant
$$\alpha_{\rm EM} = \frac{e^2}{\hbar c} \sim 1/137.036$$

Electron or quark mass/QCD strong interaction scale $\frac{m_{e,q}}{\Lambda_{QCD}}$
The electron-proton mass ratio $\frac{m_e}{M_P}$

VARIATION OF FUNDAMENTAL CONSTANTS



The modern theories directed toward unifying gravitation with the three other fundamental interactions suggest variation of the fundamental constants in an expanding universe.

Variation of fundamental constants

Theories with varying dimensionless fundamental constants

String theories

J.-P. Uzan, Living Rev. Relativity 14, 2 (2011)

- Other theories with extra dimensions
- Loop quantum gravity
- Dark energy theories: chameleon and quintessence models
- ...many others



Life needs very specific fundamental constants!



If α is too big \rightarrow small nuclei can not exist Electric repulsion of the protons > strong nuclear binding force

 $\alpha \sim 1/137$





 $\alpha \sim 1/10$

will blow carbon apart

Carbon-12

Life needs very specific fundamental constants!



Nuclear reaction in stars are particularly sensitive to α . If α were different by 4%: **no carbon produced by stars**. No life.

Life needs very specific fundamental constants!



No carbon produced by stars: No life in the Universe



A. Derevianko, Conf. Ser. 723 (2016) 012043



How to test if α changed with time?



Atomic transition energies depend on α^2



Astrophysics searches for variation of α : looking for changes in quasar light





Astrophysical searches for α -variation

Alkali-doublet method



Murphy et al. (2001)
$$\frac{\Delta \alpha}{\alpha} = -0.5(1.3) \times 10^{-5}$$

Astrophysical searches for α -variation

Many-multiplet method: compare spectra of different atoms



Need atomic calculations to find corresponding factor q



Conflicting results

Murphy et al., 2003: Keck telescope, 143 systems, 23 lines, 0.2<z<4.2

Quast et al, 2004: VL telescope, 1 system, Fe II, 6 lines

Molaro et al., 2007

Z=1.84

Srianand et al, 2004: VL telescope, 23 systems, 12 lines, Fe II, Mg I, Si II, Al II, 0.4<*z*<*2.3*

Murphy et al., 2007

 $\Delta \alpha / \alpha = -0.54(12) \times 10^{-5}$

$$\Delta \alpha / \alpha = -0.4(1.9)(2.7) \times 10^{-6}$$

$$\Delta \alpha / \alpha = -0.12(1.8) \times 10^{-6}$$

 $\Delta \alpha / \alpha = 5.7(2.7) \times 10^{-6}$

 $\Delta \alpha / \alpha = -0.06(0.06) \times 10^{-5}$

 $\Delta \alpha / \alpha = -0.64(36) \times 10^{-5}$

V.V. Flambaum, Variation of Fundamental Constants

Astrophysical searches for variation of fundamental constants



Can we look for α -variation in a lab?

Different optical atomic clocks use transitions that have different contributions of the relativistic corrections to frequencies.

Therefore, comparison of these clocks can be used to search for α -variation.

Laboratory searches for variation of fundamental constants

 $\alpha = \frac{1}{4\pi\varepsilon_0} \frac{e^2}{\hbar c}$

 $\mu = \frac{m_p}{m_e}$

- 1. Frequency of optical transitions
- $\nu \simeq c R_{\infty} A F(\alpha)$ Depends only on α

2. Frequency of hyperfine transitions

$$\nu_{\rm hfs} \simeq cR_{\infty}A_{\rm hfs} \times g_i \times \frac{m_e}{m_p} \times \alpha^2 F_{\rm hfs}(\alpha)$$

Depends on α , μ , g-factors (quark masses to QCD scale)

2. Transitions in molecules: μ only, μ and α , or all three

$$E_{\rm el}: E_{\rm vib}: E_{\rm rot} \sim 1: \bar{\mu}^{1/2}: \bar{\mu}$$
 $\bar{\mu} = 1/\mu$

Comparing different types of transitions probes different constants

(1) Measure the ratio R of optical to hyperfine (microwave Cs or Rb) clock frequencies: sensitive α , μ , g-factors (quark masses to QCD scale)

(2) Measure the ratio R of two optical clock frequencies: sensitive only to α -variation

$$E = E_0 + \boldsymbol{q} \left(\frac{\boldsymbol{\alpha}^2}{\boldsymbol{\alpha}_0^2} - 1 \right)$$

Calculate with good precision

Sensitivity of optical clocks to α -variation

$$E = E_0 + \boldsymbol{q} \left(\frac{\boldsymbol{\alpha}^2}{\boldsymbol{\alpha}_0^2} - 1 \right)$$

Enhancement factor



Need: large K for at least one for the clocks **Best case:** large K_2 and K_1 of opposite sign for clocks 1 and 2

$$\frac{\partial}{\partial t} \ln \frac{v_2}{v_1} = (K_2 - K_1) \frac{1}{\alpha} \frac{\partial \alpha}{\partial t}$$
Frequency ratio
accuracy 10⁻¹⁸ 100 10⁻²⁰
Easier to measure large effects!

Laboratory searches for α -variation

Different **atomic clocks** use transitions that have different contributions of the relativistic corrections to frequencies.

$$v(x) = v_0 + qx$$
 $x = (\alpha/\alpha_0)^2 - 1$

Therefore, comparison of different clocks can be used to search for α -variation.

NIST [Rosenband et al., Science 319, 1808 (2008)] Al⁺ / Hg⁺ atomic clocks

$$\dot{\alpha}/\alpha = -1.6(2.3) \times 10^{-17} \text{ y}^{-1}$$

199Hg⁺ Energy Levels



- Atomic line $Q \approx 5 \times 10^{14}$
- State detection by electron shelving.

from Jim Bergquist' talk NIST

Trapped ions in an rf trap

- No static E or B fields; Trap acts on total charge of ion, not internal structure
- Trap ion at trap center where trapping fields approach zero



• Can operate in tight-confinement (Lamb-Dicke) regime \Rightarrow First-order doppler free. 2nd-order doppler shift (time dilation) due to micromotion will limit accuracy $\frac{\Delta f}{f_0} \sim 10^{-18}$

from Jim Bergquist' talk NIST



Some facts about Al⁺

- 8 mHz linewidth clock transition
- Small quadratic Zeeman shift (6x10⁻¹⁶ /Gauss²)
- Negligible electric-quadrupole shift (J=0)
- Smallest known blackbody shift (8x10⁻¹⁸ at 300K)
- Linear Zeeman shift 4 kHz/Gauss (easily compensated)
- Light mass (2nd order Doppler shifts)
- No accessible strong transition for cooling & state detection

167 nm

 S_0

267 nm

1121 THz

I = 5/2

Clock state transfer to Be⁺ (simplified)

1. Cool to motional quantum ground state with Be+

- 2. Depending on clock state, add vibrational energy via Al+
- 3. Detect vibrational energy via Be⁺

PHYSICAL REVIEW A

VOLUME 42, NUMBER 5

1 SEPTEMBER 1990

Quantum-limited cooling and detection of radio-frequency oscillations by laser-cooled ions

D. J. Heinzen and D. J. Wineland

Time and Frequency Division, National Institute of Standards and Technology, Boulder, Colorado 80303 (Received 13 March 1990)

A single trapped ion, laser cooled into its quantum ground state of motion, may be used as a very-low-temperature detector of radio-frequency signals applied to the trap end caps. If the signal

from Jim Bergquist' talk

Using two ions



from Jim Bergquist' talk

from Jim Bergquist' talk

Al+/Hg+ Comparison fs-comb locked to Hg+ measure beat with Al+



from Jim Bergquist' talk

Al+/Hg+ Comparison





Constraints on temporal variations of α and μ from comparisons of atomic transition frequencies. Huntemann et al., PRL 113, 210802 (2014)

Variation of fundamental constants and dark matter



Rotation curves







The dark matter



dark matter



Image credit: Simon Knapen,

Ultralight dark matter

Bosonic dark matter (DM) with mass $m_{\phi} < 1 \text{eV}$ Dark matter density in our Galaxy > λ_{dR}^{-3}

 λ_{dB} is the de Broglie wavelength of the particle.

Then, the scalar dark matter exhibits coherence and behaves like a wave.

$$\phi(t) = \phi_0 \cos\left(m_\phi t + \bar{k}_\psi \times \bar{x} + \dots\right)$$

Dilatons: appears in theories with extra dimensions when the volume of the compactified dimensions is allowed to vary.

Dilatonic couplings to the Standard Model



A. Arvanitaki et al., PRD 91, 015015 (2015)

Ultralight dark matter

$$\phi(t) = \phi_0 \cos\left(m_\phi t + \bar{k}_\phi \times \bar{x} + \dots\right)$$

DM virial velocities ~ 300 km/s

Dark matter parameters

	$m_{\phi} [{ m eV}]$	$f = 2\pi/m_{\phi} \; [\text{Hz}]$	$ au~[{ m s}]$
	4×10^{-9}	$1 \mathrm{~MHz}$	10^{-6}
	4×10^{-12}	$1 \mathrm{~kHz}$	10^{-3}
One oscillation per second	4×10^{-15}	1	1
	4×10^{-18}	$1 \mathrm{~mHz}$	1000
One oscillation per 11 days	4×10^{-21}	10^{-6}	10^{6}

Clock measurement protocols for the dark matter detection

Single clock ratio measurement: averaging over time τ_1

Make N such measurements, preferably regularly spaced



Detection signal:

A peak with monochromatic frequency $f = 2\pi/m_{\phi}$ in the discrete Fourier transform of this time series.

A. Arvanitaki et al., PRD 91, 015015 (2015)

Measuring ratios of optical clock frequencies for dark matter detection

$$\frac{\delta(\nu_2/\nu_1)}{(\nu_2/\nu_1)} \simeq d_e(K_2 - K_1)\kappa\phi(t)$$

Need:

- Best short-term stability σ_1 at $\Delta \tau$
- Long total measurement time to improve sensitivity

$$\sigma_N = \sigma_1 / \sqrt{N}$$

But: only until you reach the DM coherence time

$$au_{
m coh} \simeq 2\pi (m_{\phi}v^2)^{-1}$$
 $v \approx 10^{-3}$
(Not an issue for 10⁶ s)

- Lowest systematic uncertainty
- Largest possible enhancement factor combination (K₂-K₁)

Experimental constraints for dilaton dark matter



nature physics

Transient variations

STATION 3

STATION &

STATION 4

Hunting for topological dark matter with atomic clocks

A. Derevianko^{1*} and M. Pospelov^{2,3}

Dark matter clumps: point-like monopoles, onedimensional strings or two-dimensional sheets (domair. walls).

Topological dark matter may have formed when the early Universe cooled down after the Big Bang, similar to the domains formed in a ferromagnet below its Curie temperature.

If they are large (size of the Earth) and frequent enough we can detect this with atomic clocks.

Yang at al., Scientific Reports 5, 11469 (2015)



ETTERS

Nature Communications 8, 1195 (2017)



Topological dark matter may be detected by measuring changes in the synchronicity of a global network of atomic clocks, such as the Global Positioning System, as the Earth passes through the domain wall.

Rana Adhikari, Paul Hamiton & Holger Müller, Nature Physics 10, 906 (2014)

Search for the violation of the Lorentz invariance

Lorentz invariance: the laws of physics that govern a physical system are unchanged for different system orientations or velocities.



Search for the violation of the Lorentz invariance

Lorentz invariance: the laws of physics that govern a physical system are unchanged for different system orientations or velocities.



Animation credit: http://www.physics.indiana.edu/~kostelec/mov.html

The basic idea of atomic physics tests of Lorentz invariance:

Atomic energy levels are affected differently by Lorentz violation: transition frequency will change when experimental set up rotates or moves



Experimental strategy:

- (1) Pick two energy levels that should shift differently due to Lorentz violation.
- (2) Turn on magnetic field to define a quantization axis.
- (3) Keep measuring the transition frequency between these two levels while Earth rotates and and moves around the Sun.
 [DO NOT ROTATE EXPERIMENT YOURSELF].

Animation credit: http://www.physics.indiana.edu/~kostelec/mov.html

Measure frequency difference of two Yb+ PTB clocks to probe Lorentz violation



The Future: Atomic Clocks in the Next Quantum Revolution



Measurements beyond the quantum limitEntangled clocksOrders of magnitude improvements with current clocks

Image credits: NIST, Innsbruck group, MIT Vuletic group, Ye JILA group

The Future: Atomic Clocks in the Next Quantum Revolution



a) ++++++ b) x 5 + * 0.1 mm

Nuclear clock

Clocks with ultracold highly charged ions

Science 347, 1233 (2015)

Th ion nuclear clock

Th nuclear clock:

Nuclear isomer transition in 229 Thorium has been suggested as an etalon transition in a new type of optical frequency standard.



Possible orders of magnitude enhancement to the

variation of α and $\frac{m_q}{\Lambda_{QCD}}$ but orders of magnitude uncertainty in the enhancement factors.

Atomic Clocks of the Next Quantum Revolution: Quantum-science Enabled Discoveries



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Are fundamental constants constant?







Quantum gravity and violation of Lorentz invariance?

Does equivalence principle breaks down?

